

# ISM models with SN driven turbulence

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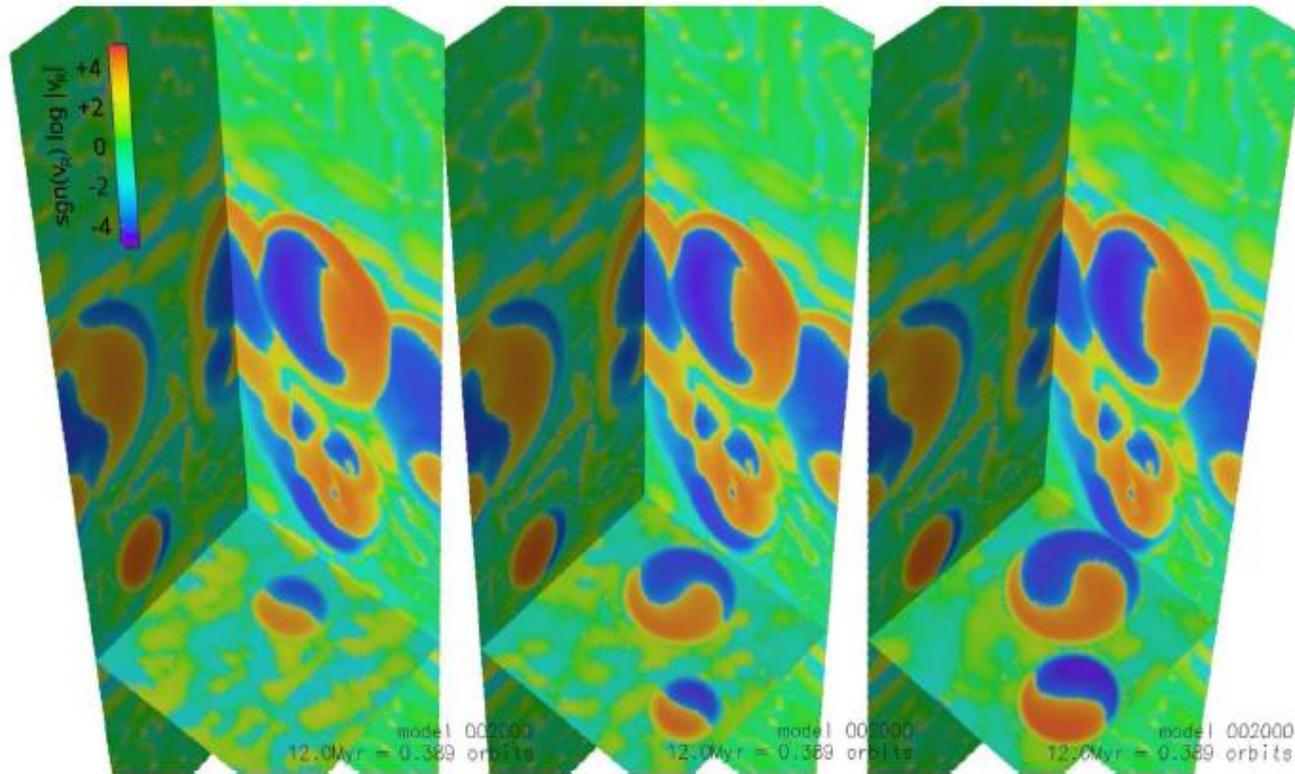
# Outline

- ISM simulations with SNE
- Dynamo models until saturation
- SNR dependence
- Explaining the saturation

# Introduction

- ISM consists of
  - i) Many gaseous phases
  - ii) Magnetic field (few  $\mu\text{G}$ )
  - iii) Cosmic rays
- Heating and cooling
- ISM turbulence = mainly contributed by SN II (isolated and clustered)
- Energy of SN II explosion :  $1.14 \times 10^{51}$  erg

# Simulation box



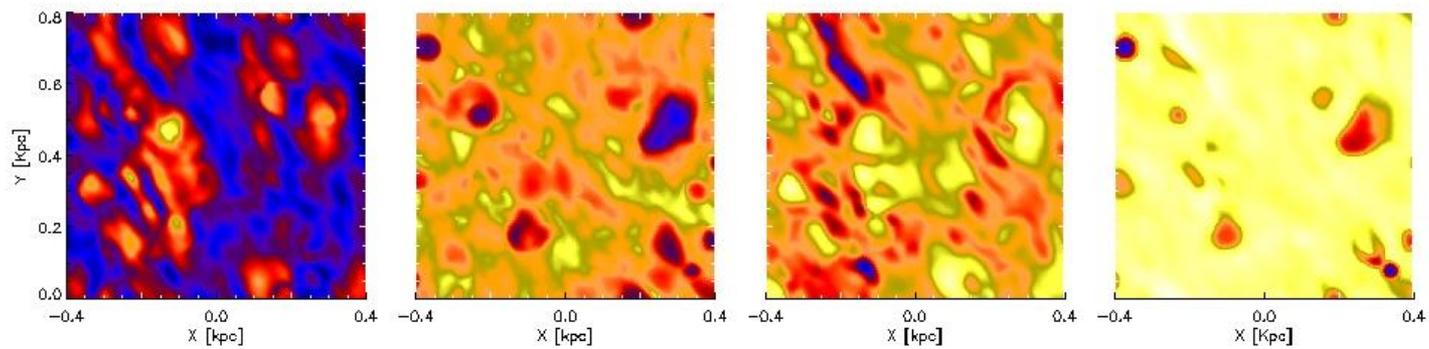
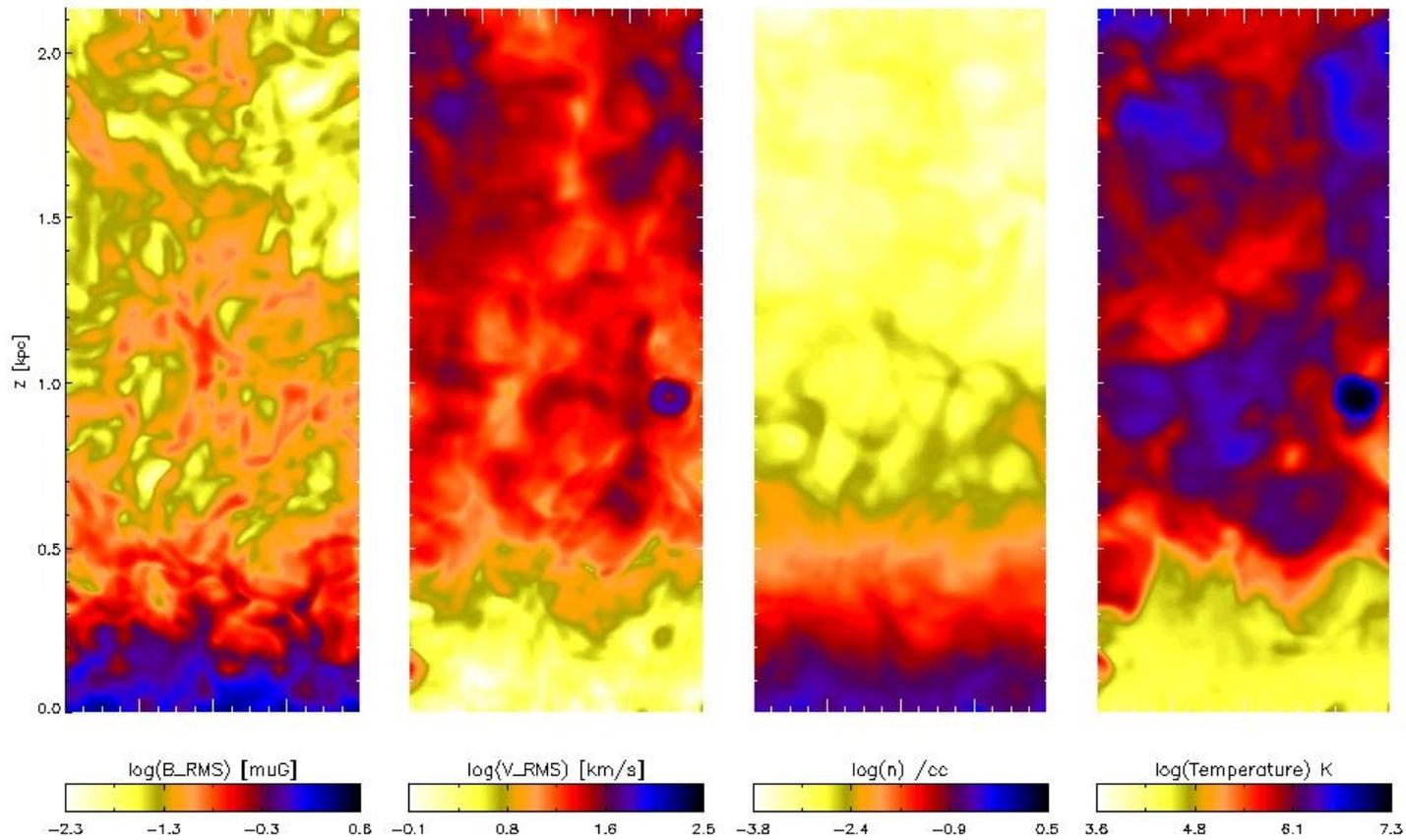
$xy$  plane: 0.8 kpc x 0.8 kpc (96x96 grid cells)

Vertical size  $z$ : -2 kpc to 2 kpc (512 grid cells)

# Modeling

- Density stratification
- Rotational shear is included (shearing periodic boundary conditions)
- Radiative cooling function
- SN explosions as localized injections of thermal energy
- SN explosions frequency,

$$\sigma_{\text{SN-I}} = 4 \text{ myr}^{-1} \text{ kpc}^{-2}, \quad \sigma_{\text{SN-II}} = 30 \text{ myr}^{-1} \text{ kpc}^{-2}$$



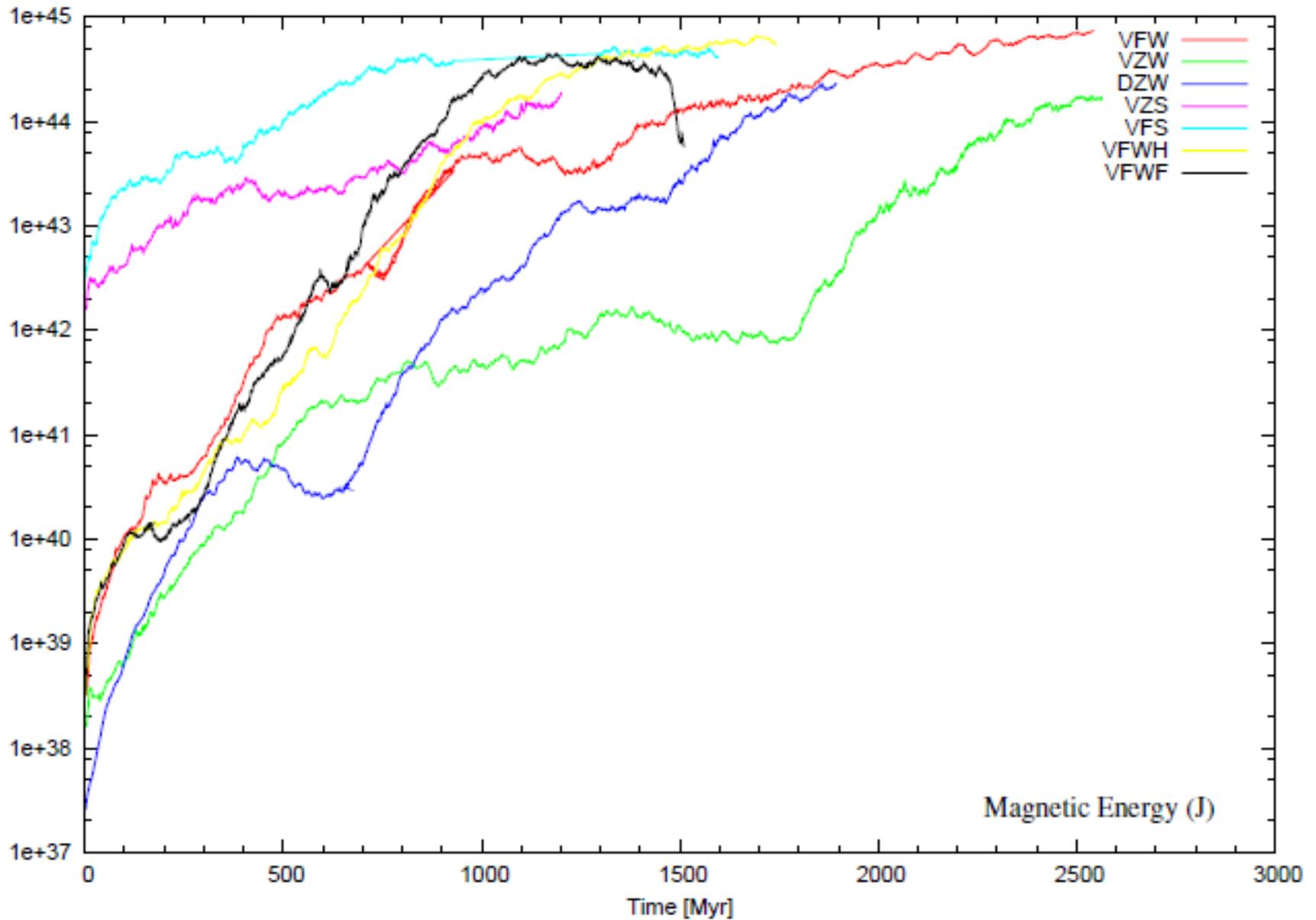
# Dependence on the seed fields

- Azimuthal seed fields (Gressel O.)
- Vertical seed fields
  - i) With and without flux
  - ii) Strong and weak initial field strength

IGM magnetic fields  $< 1e^{-8}G$

- Vertical flux passing through the disk of galaxy
- May influence MRI
- May influence dynamo action

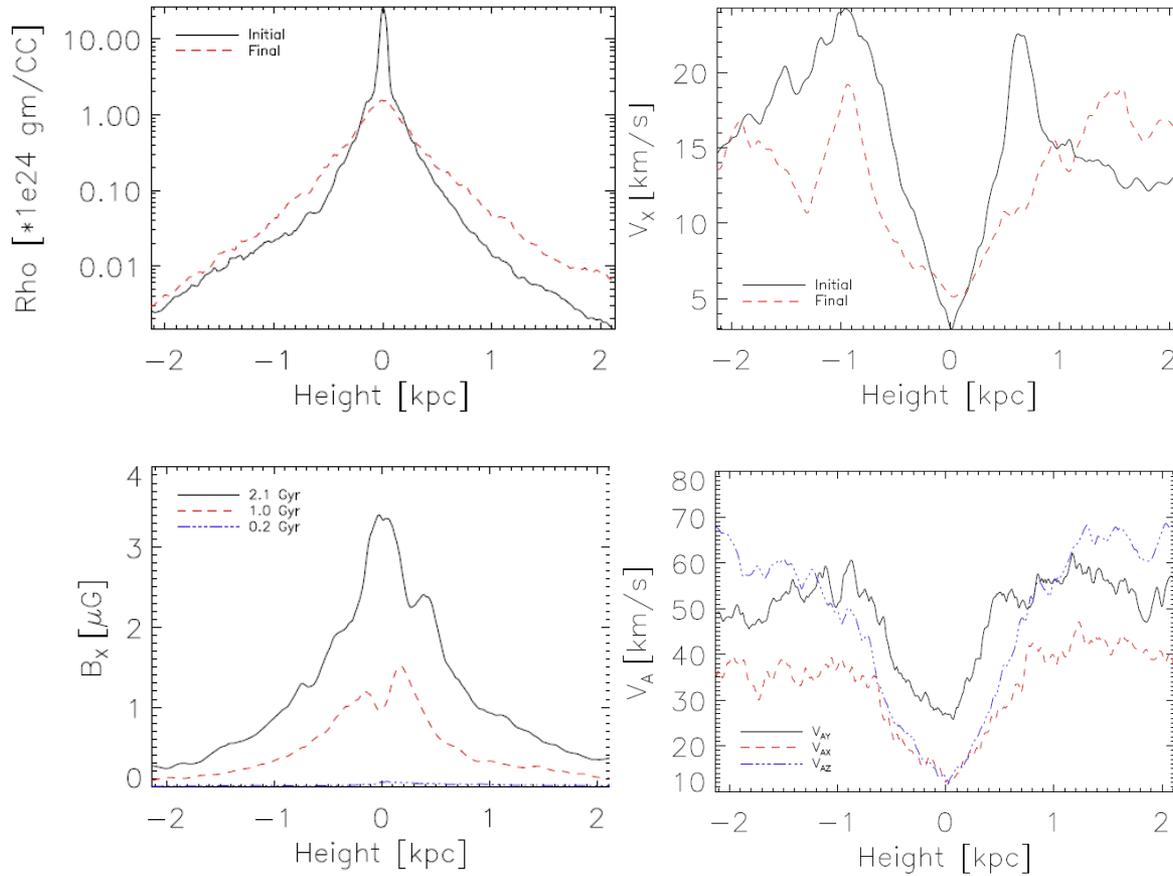
# Evolution of magnetic energy



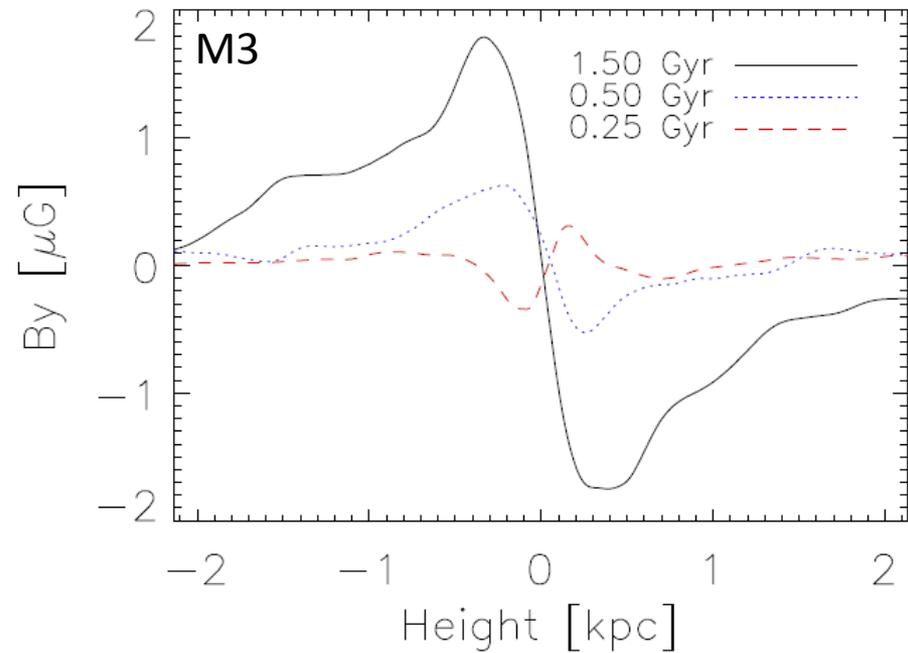
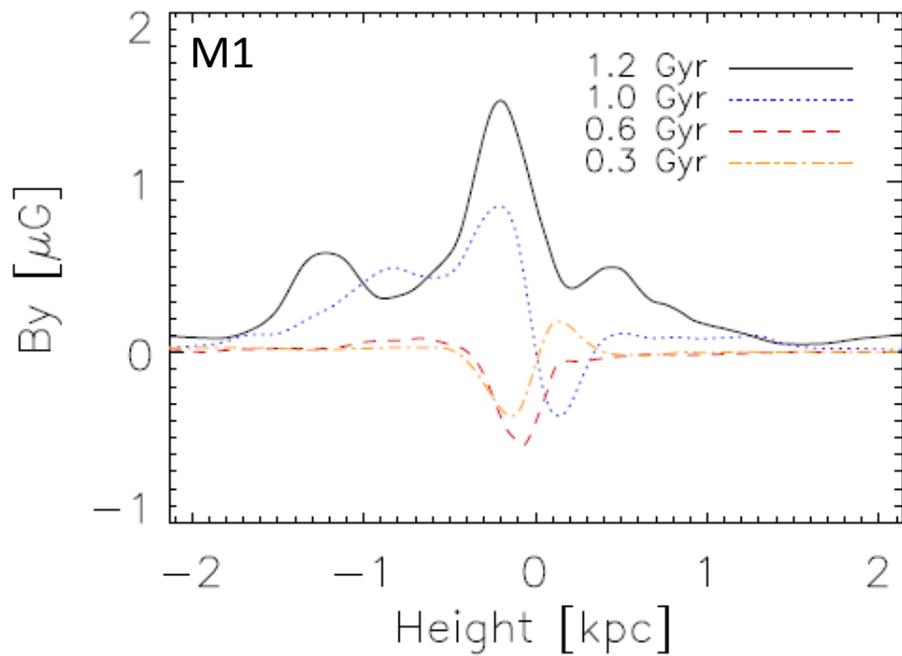
# Dependence on the seed fields

- Growth times of  $E_{\text{MAG}}$  for all models are same in the initial growing region (except the strong field models)
- Strong field vertical flux model is in the over equipartition regime and  $E_{\text{MAG}}$  Saturated
- Strong field zero flux model still in the growth phase

# Properties of the ISM



# B Azimuthal evolution

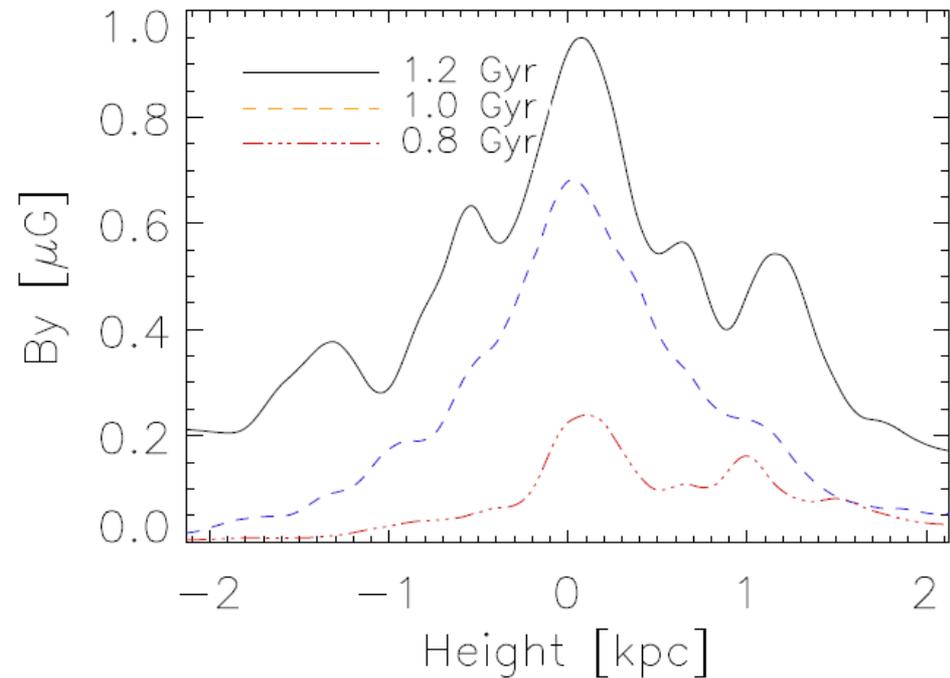
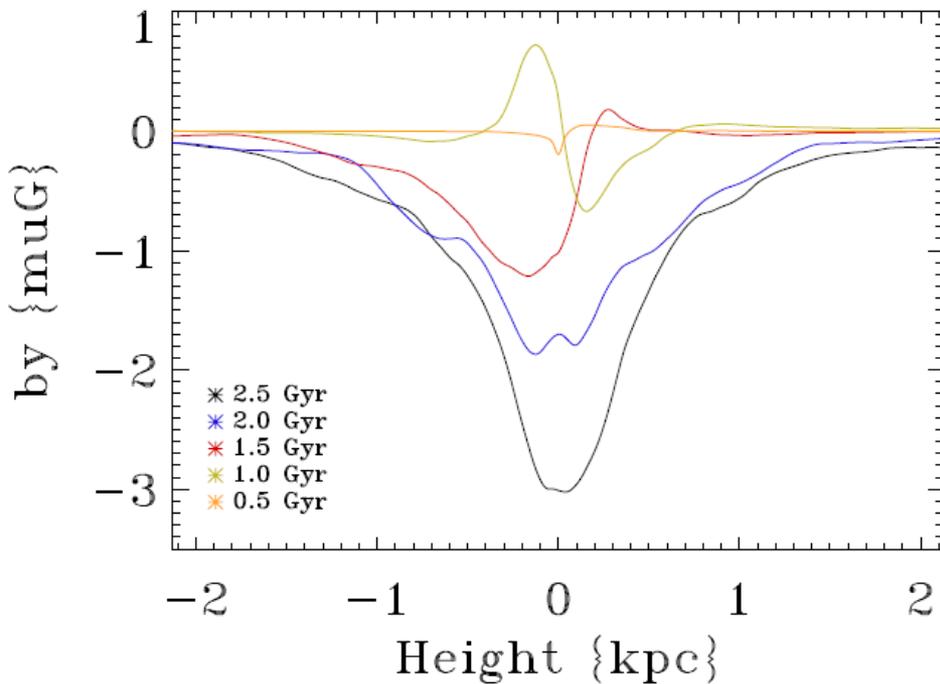


# SNR Dependence

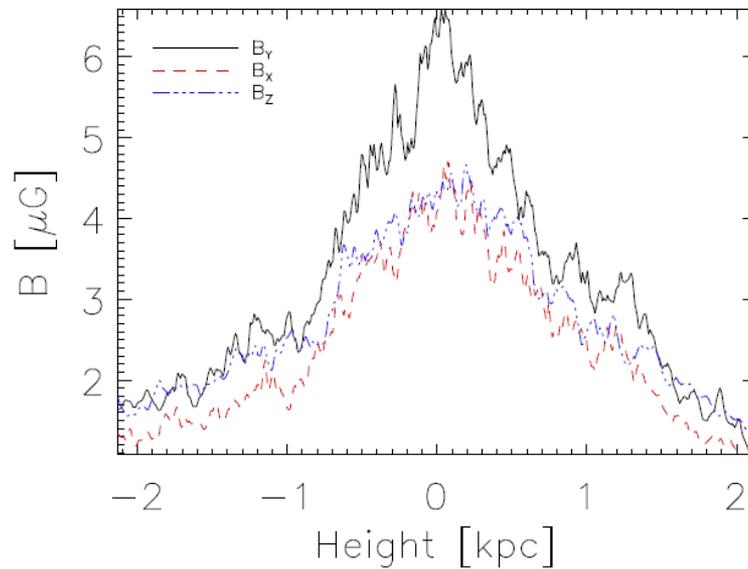
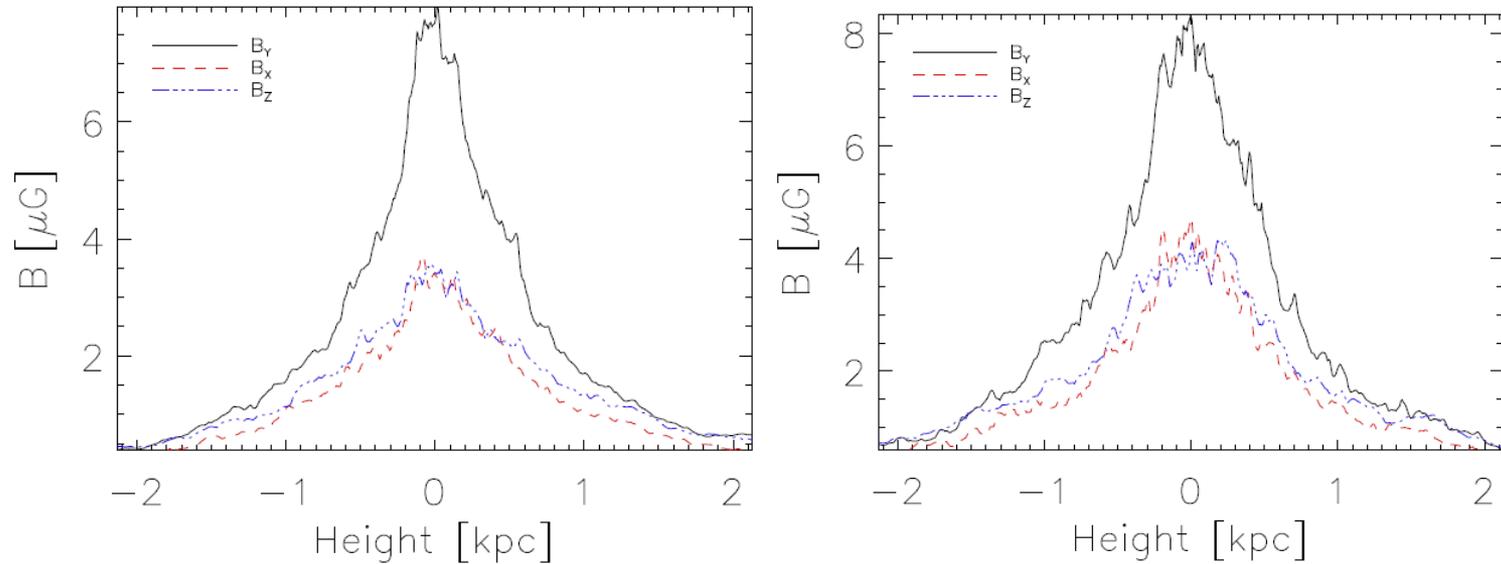
- For weak field strength with vertical flux
- Three models with different SNR (calculated until  $E_{\text{MAG.}}$  saturates)
- Seed field :  $B_z = 1 \times 10^{-3} \mu\text{G}$ , (Flux  $10^{39} \mu\text{Gcm}^2$ )
- Growth times of  $E_{\text{MAG.}}$  are same in the initial growing phase (200 Myr)
- Total magnetic energy is  $4 \times 10^{51}$  erg for all SNR

	VWQ	VWH	VWF
SNR (%) (times $\sigma$ )	25%	50%	100%
$E_{\text{mag}} : E_{\text{kin}}$ (final)	2.3	1.1	0.4

# B azimuthal Vs. Z, Profiles evolution

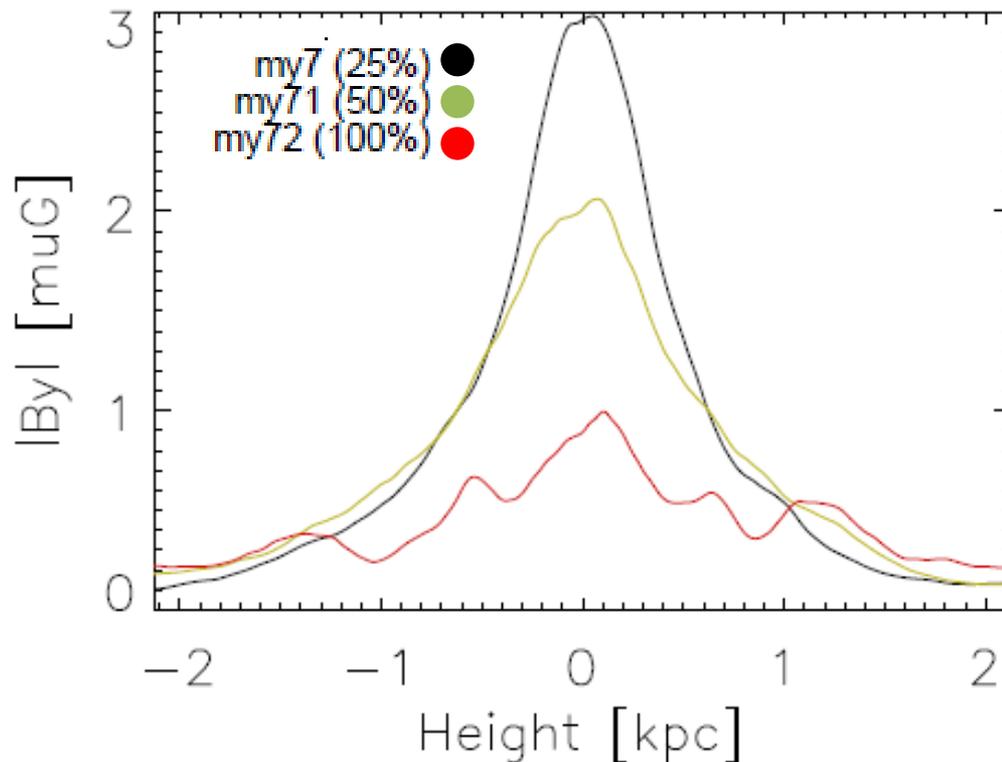


# Maximum $|B|$



# $B_y$ Profiles (SNR dependence)

Final field strengths are inversely proportional to SNR



# Saturation

Alpha quenching or/and Wind quenching

$$\partial_t \bar{B} = \nabla \times (\bar{\mathbf{v}} \times \bar{B}) + \nabla \times \boldsymbol{\varepsilon} + \eta \nabla^2 \bar{B} \quad \boldsymbol{\varepsilon} = \overline{\mathbf{v}' \times \mathbf{b}}$$

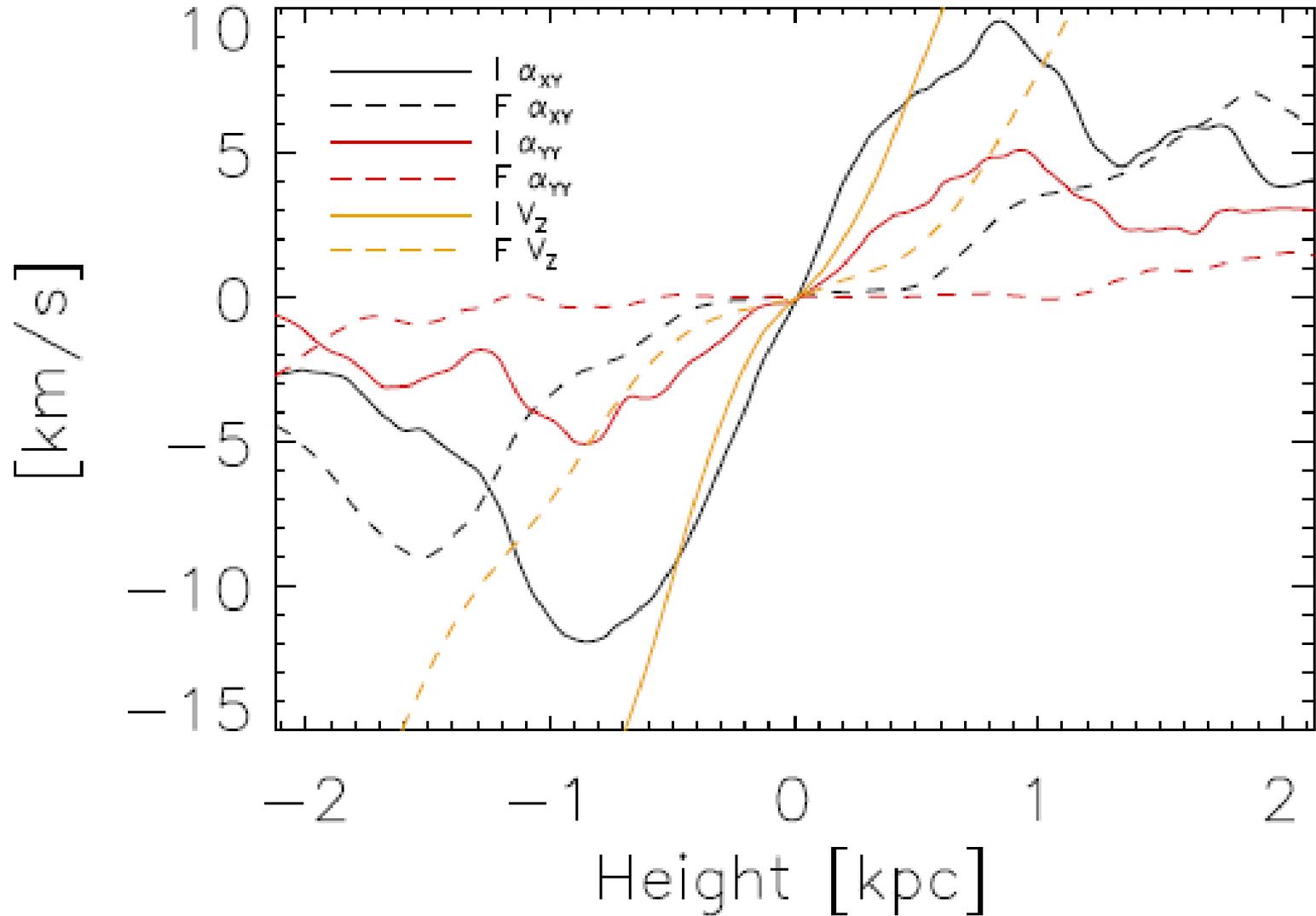
$$\varepsilon_i = \alpha_{ij} \bar{B}_j + \eta_{ijk} \partial_k \bar{B}_j$$

*Correlation*

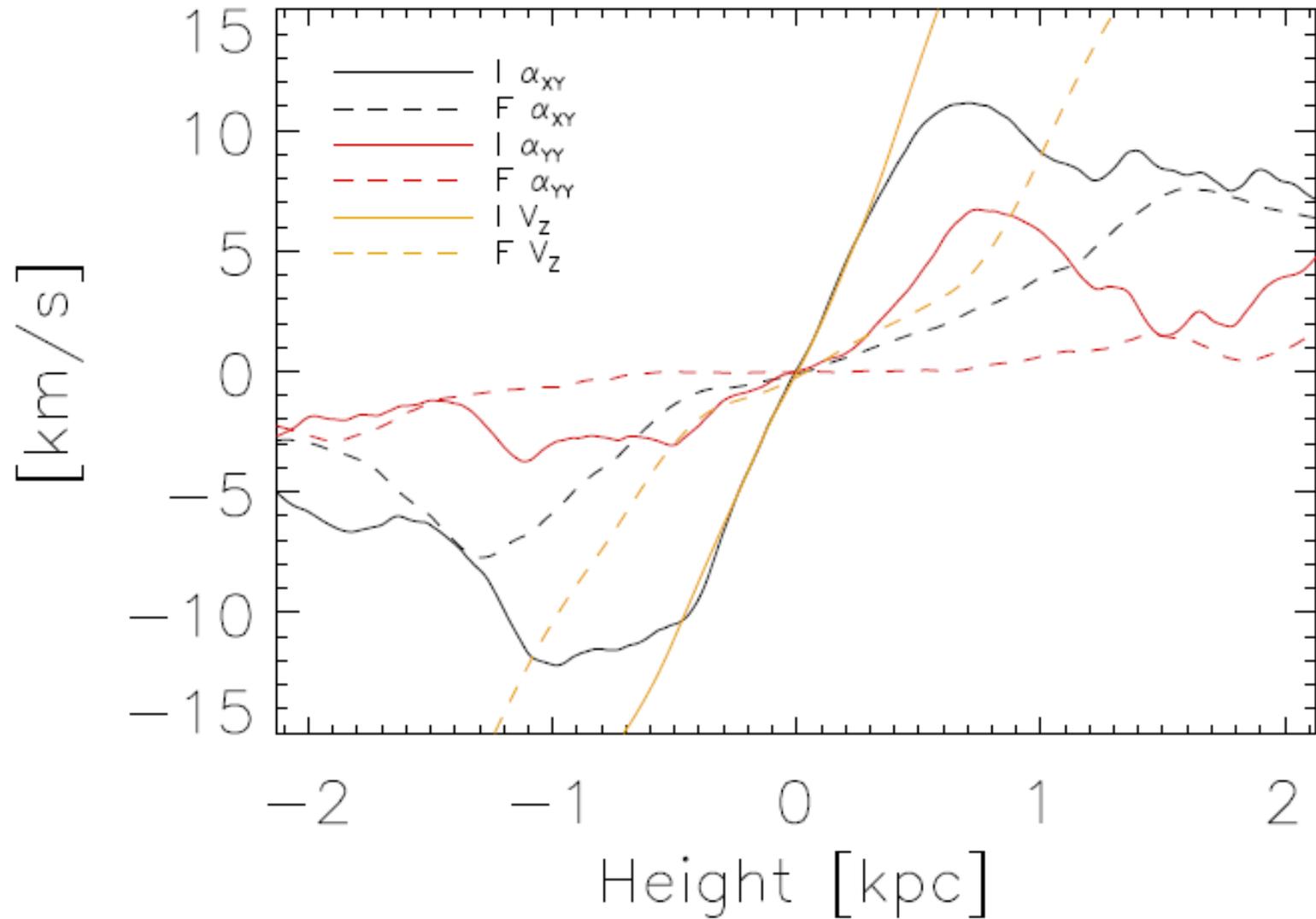
$$\begin{aligned} \partial_t \bar{B}_R &= \partial_z \left[ \begin{array}{l} \text{Wind} \\ \text{Pumping} \end{array} \right] (\bar{u}_z + \gamma) \bar{B}_R - \alpha_\varphi \bar{B}_\varphi + (\eta_\varphi + \eta) \partial_z \bar{B}_R + \delta \partial_z \bar{B}_\varphi \\ \partial_t \bar{B}_\varphi &= \partial_z \left[ \begin{array}{l} \text{Alpha} \\ \text{Turbulent diffusivity} \end{array} \right] \bar{B}_R - (\bar{u}_z + \gamma) \bar{B}_\varphi - \delta \partial_z \bar{B}_R + (\eta_R + \eta) \partial_z \bar{B}_\varphi - q \Omega \bar{B}_R \end{aligned}$$

*Differential rotation*

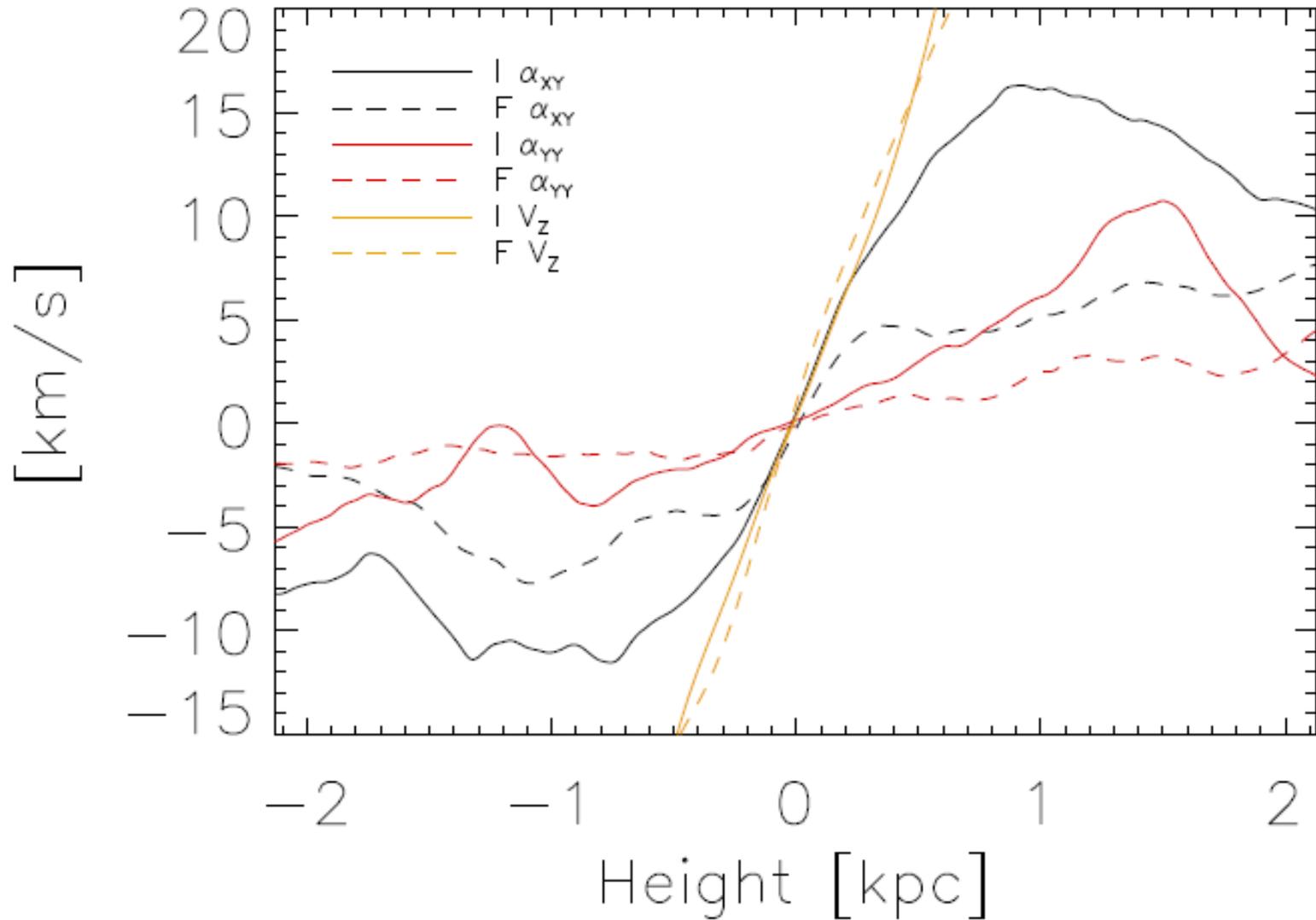
# Alpha Quenching (25% SNR)



# Alpha Quenching (50% SNR)



# Alpha Quenching (100% SNR)

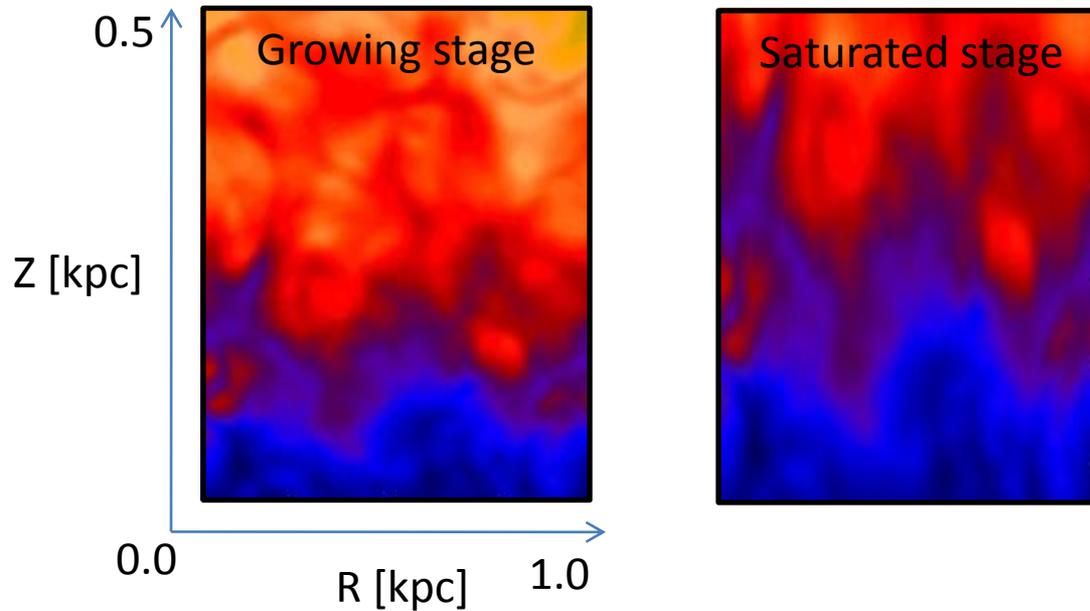


# Saturation Process

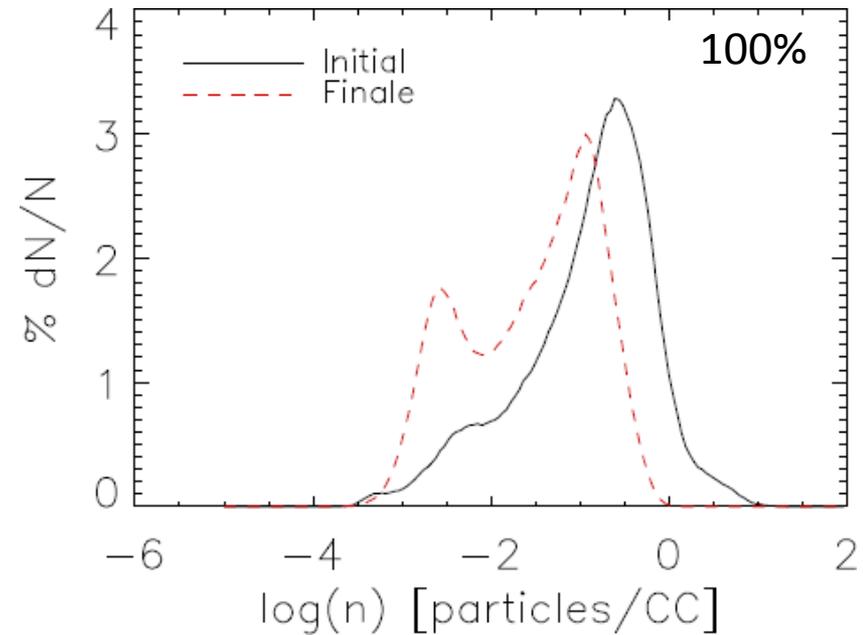
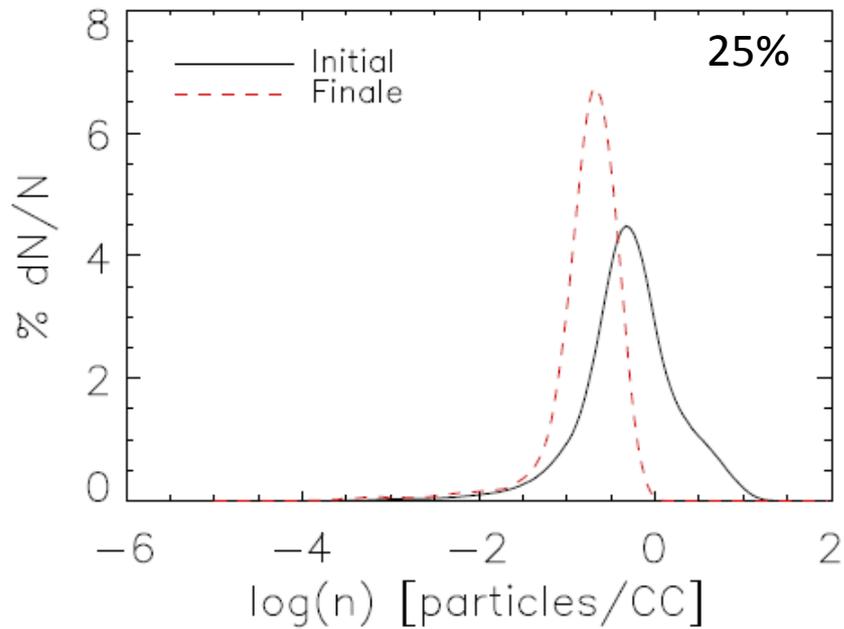
- Magnetic energy gets saturated at equal magnitude for all SNE rates, ( $4 \times 10^{51}$  erg)
- Quenching of wind and alpha profile -> dynamo stops
- In the last model with 100% SNR wind did not quench

Why?

# Mass transport Vs. density distribution



# Volume filling fraction (25% & 100% SNR)



May be due to the relatively broader density distribution

# Summary

- Magnetic energy saturates at the equal magnitude, irrespective of the seed fields and SNE frequency
- Growth rates are almost equal for all SNR (in the initial growing phase )
- Absolute value of the mean magnetic field decrease with increasing SNR
- $\alpha_{R\varphi}$  and  $\alpha_{\varphi\varphi}$  profiles are quenched in the saturated region
- $\alpha_{R\varphi}/\alpha_{\varphi\varphi} = 4$  for all SNR (Elstner D. , Gressel O.)
- Except the 100% SNR model, wind is also quenched in all other models
- Wind may have quenched due to the relatively narrow distributions of density in the mid planes

# Outlook

- Models with higher resolution
- Larger box sizes
- Cosmic Rays

# Saturated energies

Model	Final energy (erg)	Kinetic energy (erg)
Strong field vertical flux (25%)	$4 \times 10^{51}$	$2 \times 10^{51}$
Strong field zero flux (25%)	Growing phase	$2 \times 10^{51}$
Weak field vertical flux (25%)	$4 \times 10^{51}$	$2 \times 10^{51}$
Weak field zero flux (25%)	Growing phase	$2 \times 10^{51}$
Weak field vertical flux (50%)	$4 \times 10^{51}$	$4 \times 10^{51}$
Weak field vertical flux (100%)	$4 \times 10^{51}$	$10 \times 10^{51}$
Azimuthal and radial fields (25%)	Growing phase	$2 \times 10^{51}$

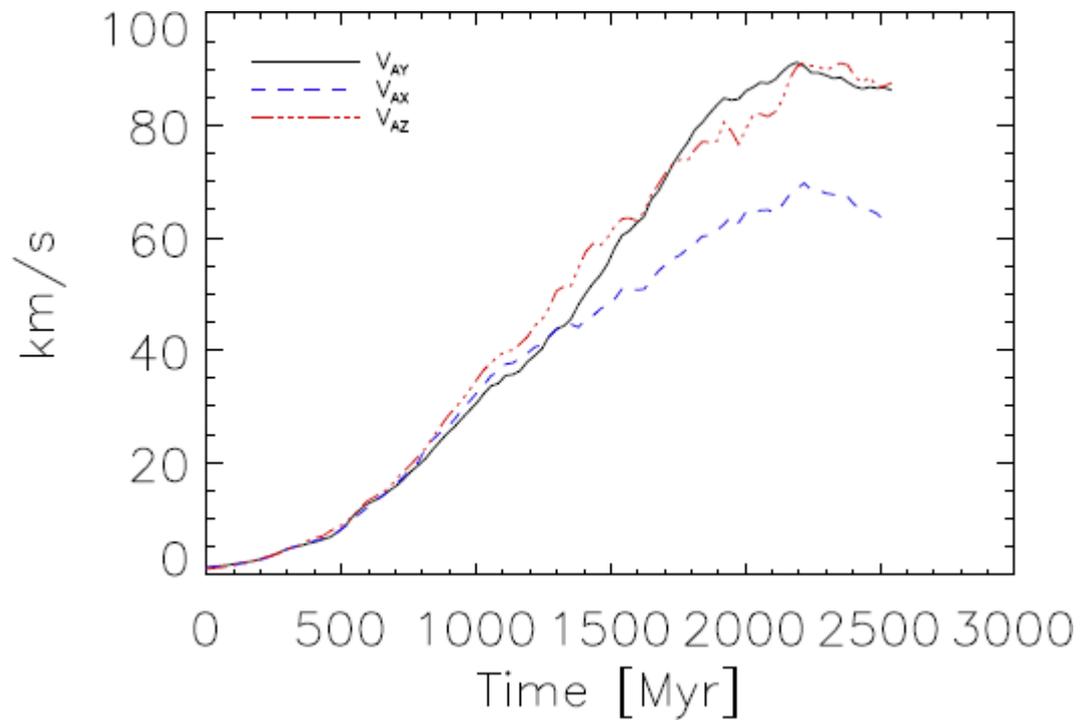
# Radiative cooling function

- Modeling of SNE via thermal energy injections
- Necessary to include radiative cooling function
- Adopted as piecewise power law

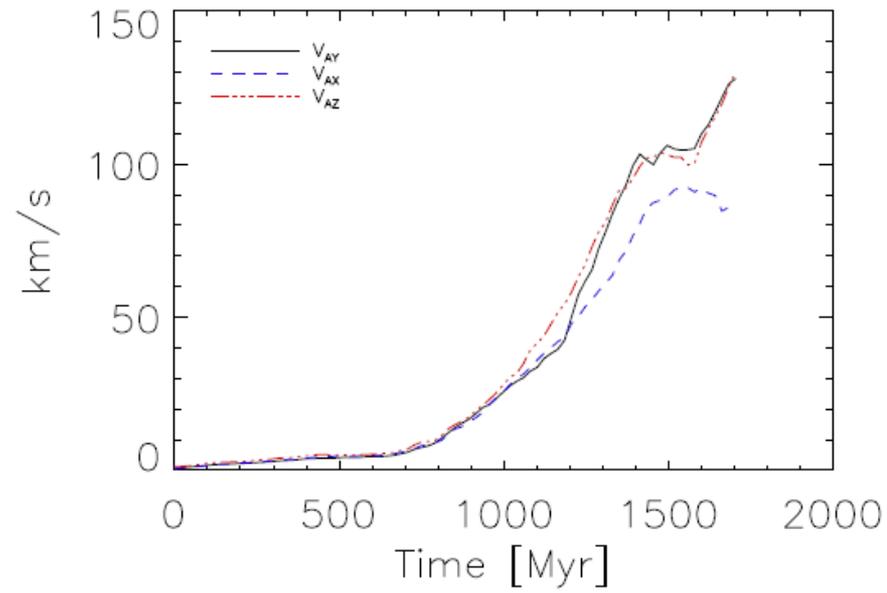
$$\Lambda(T) = \Lambda_i T^{\beta_i}, \quad \text{for } T_i \leq T < T_{i+1}.$$

- Thermally unstable range 141K to 6102K

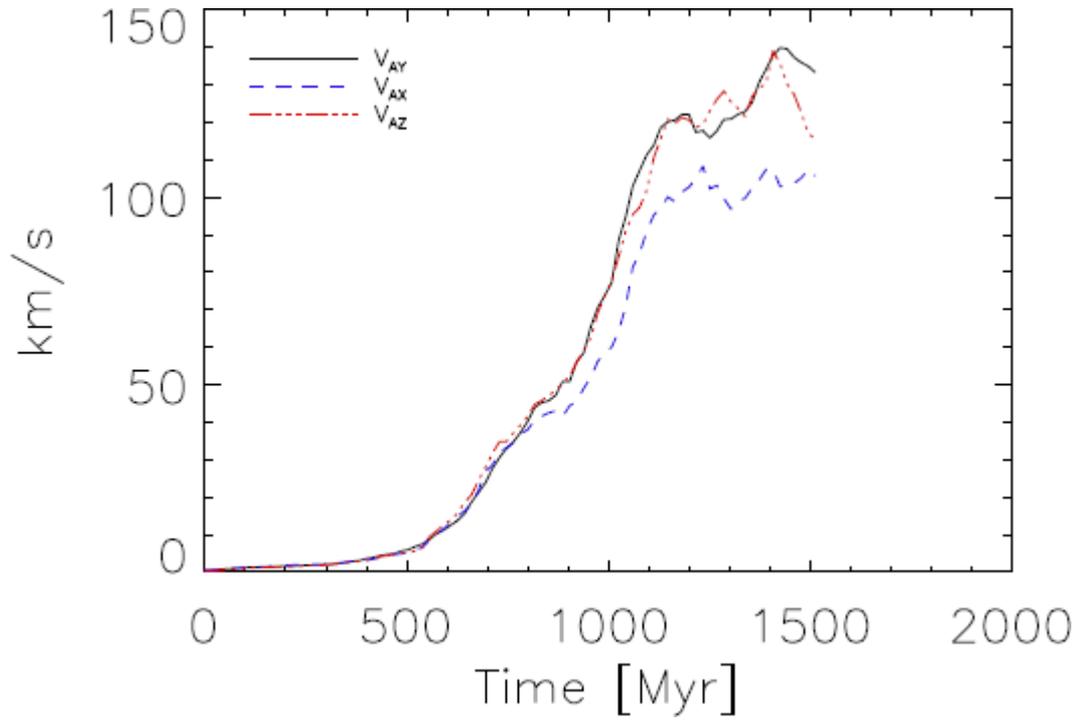
# 25% SNR



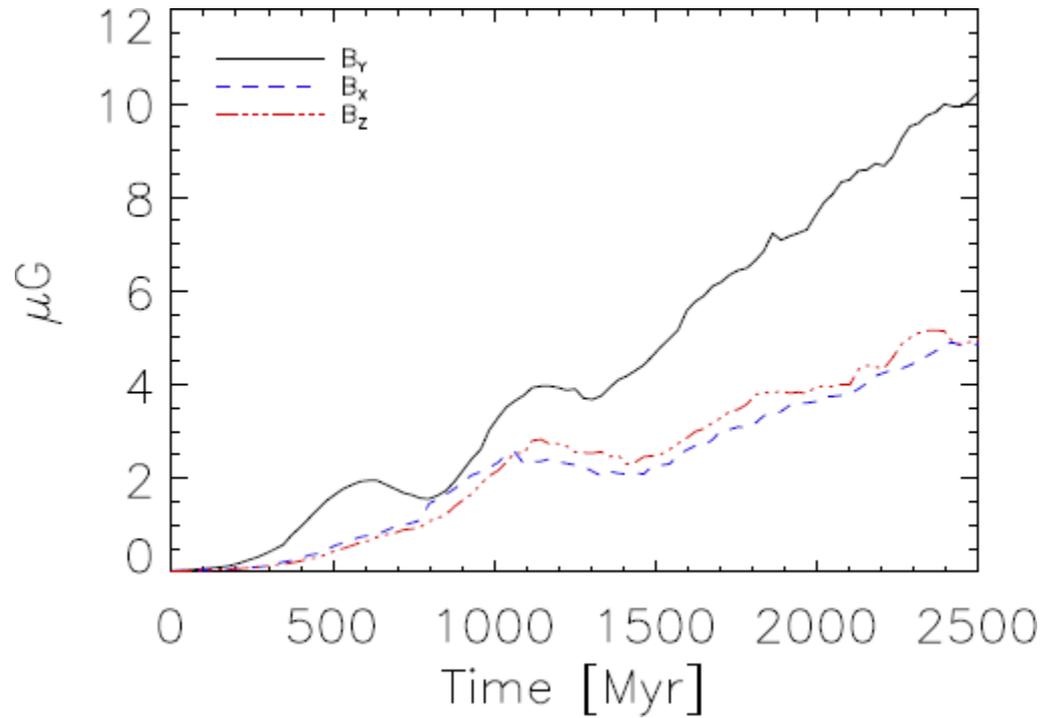
# 50% SNR



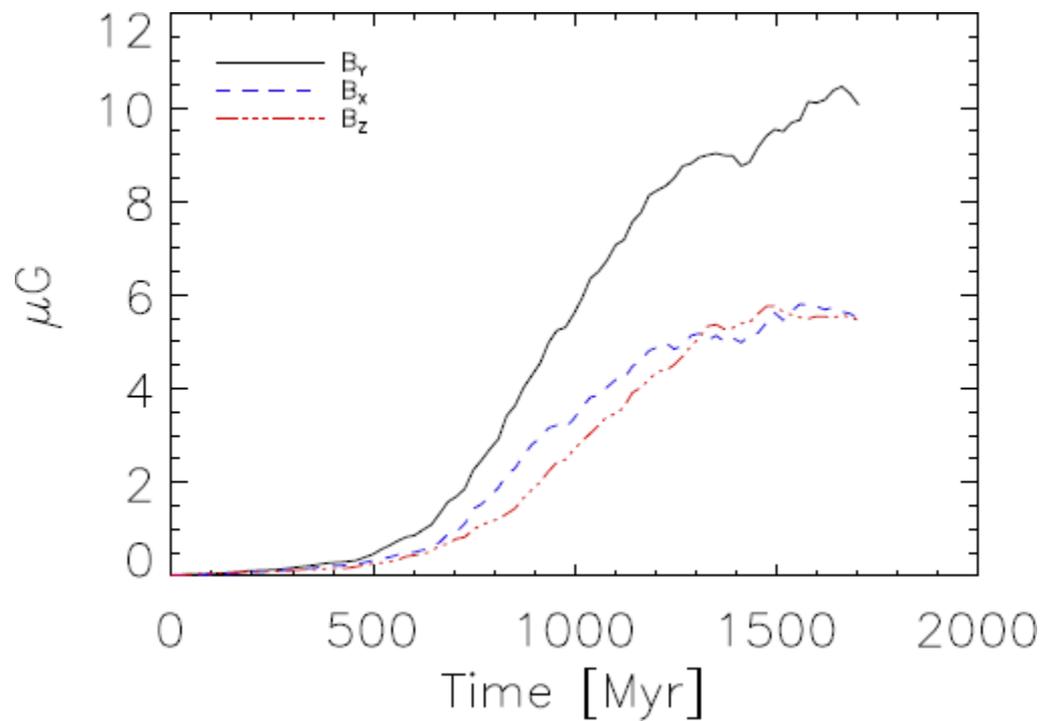
# 100% SNR



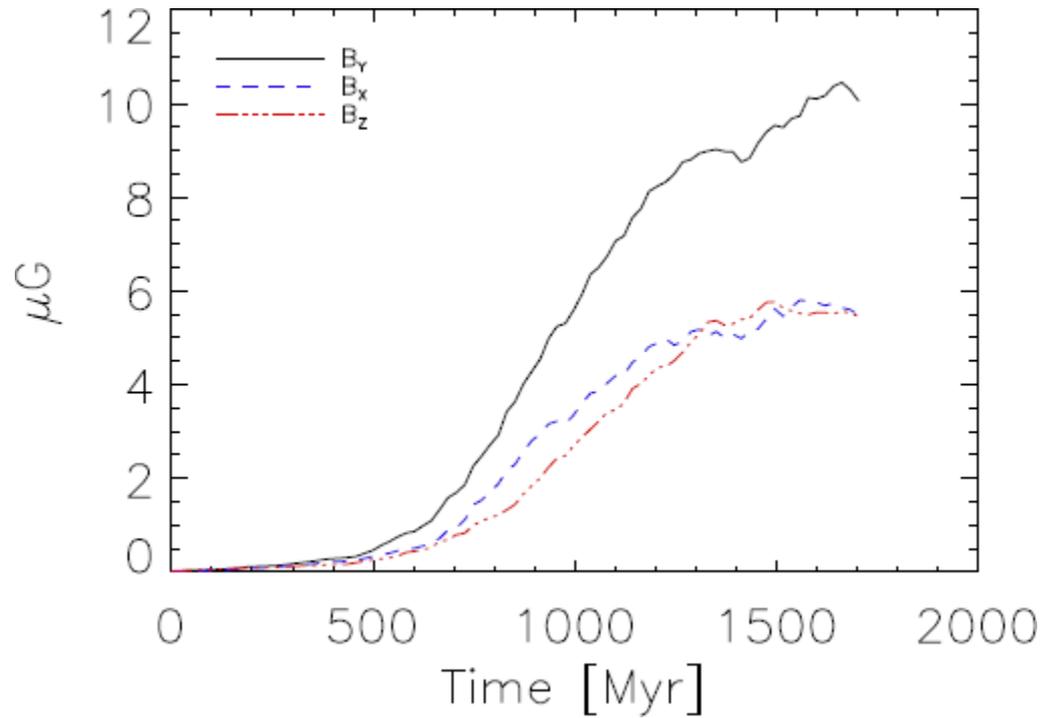
# 25% SNR



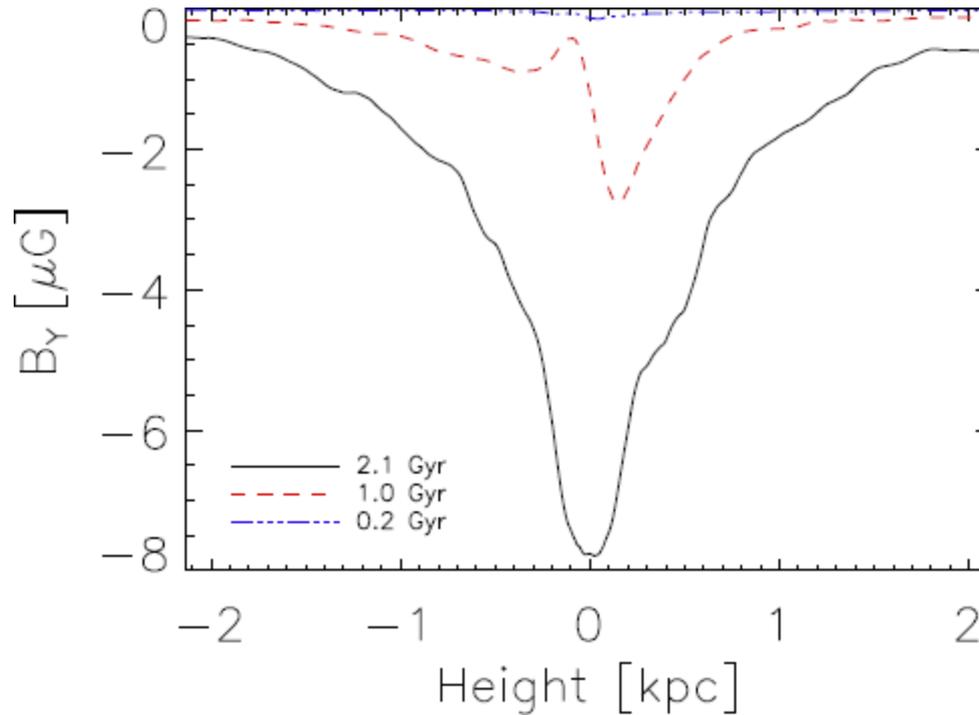
# 50% SNR



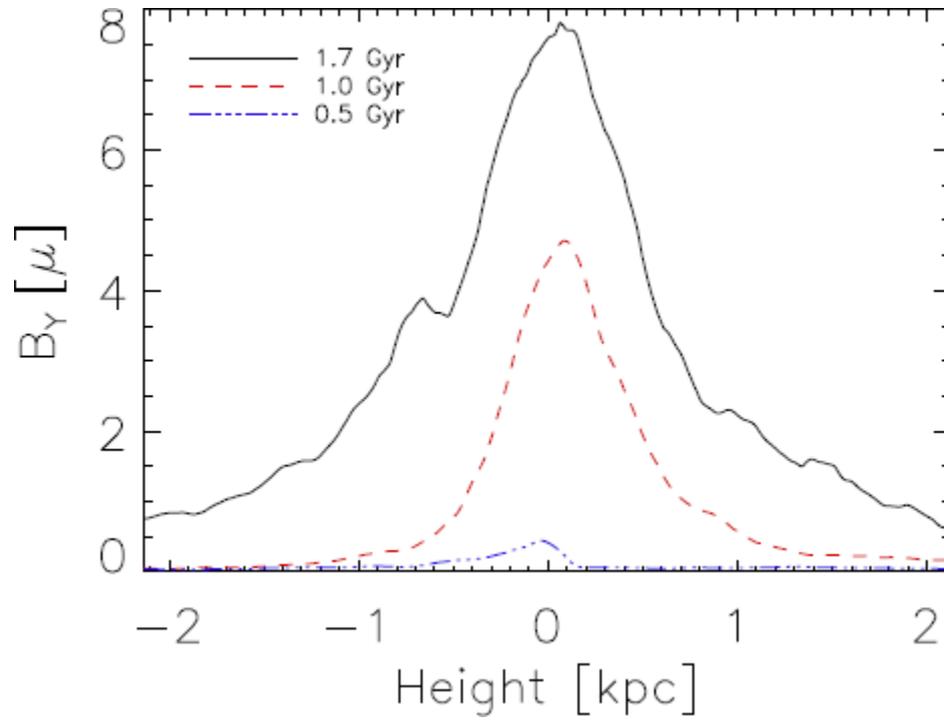
# 100% SNR



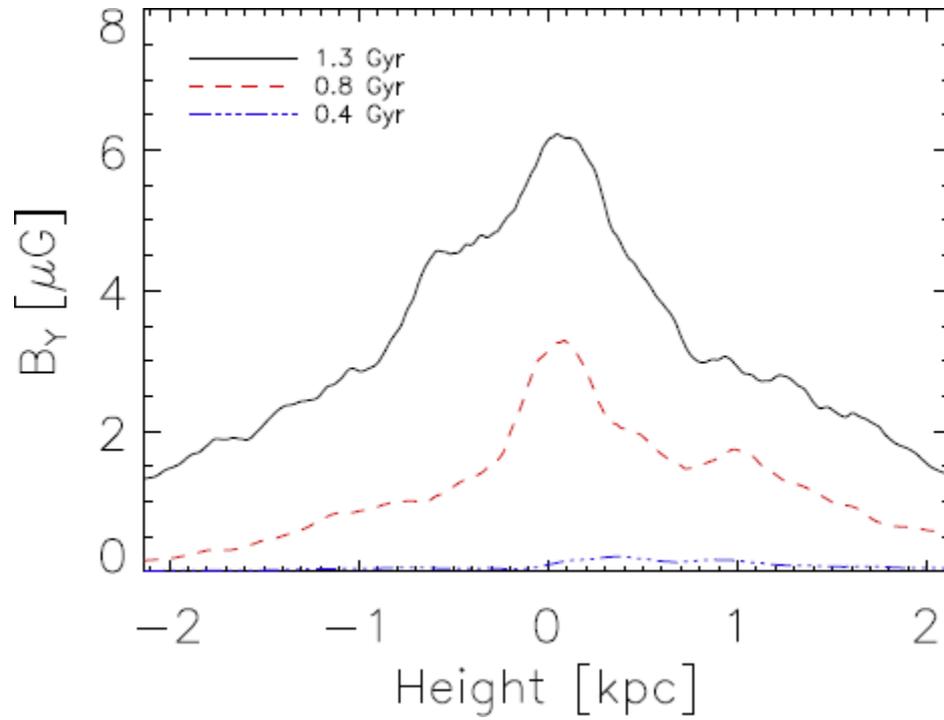
# By maximum evolution 25%



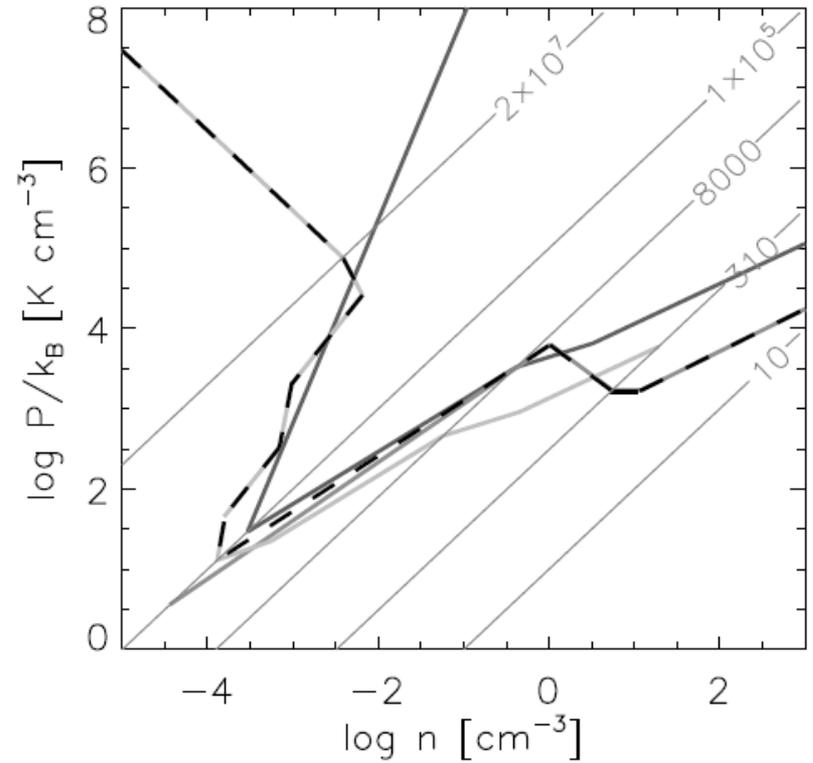
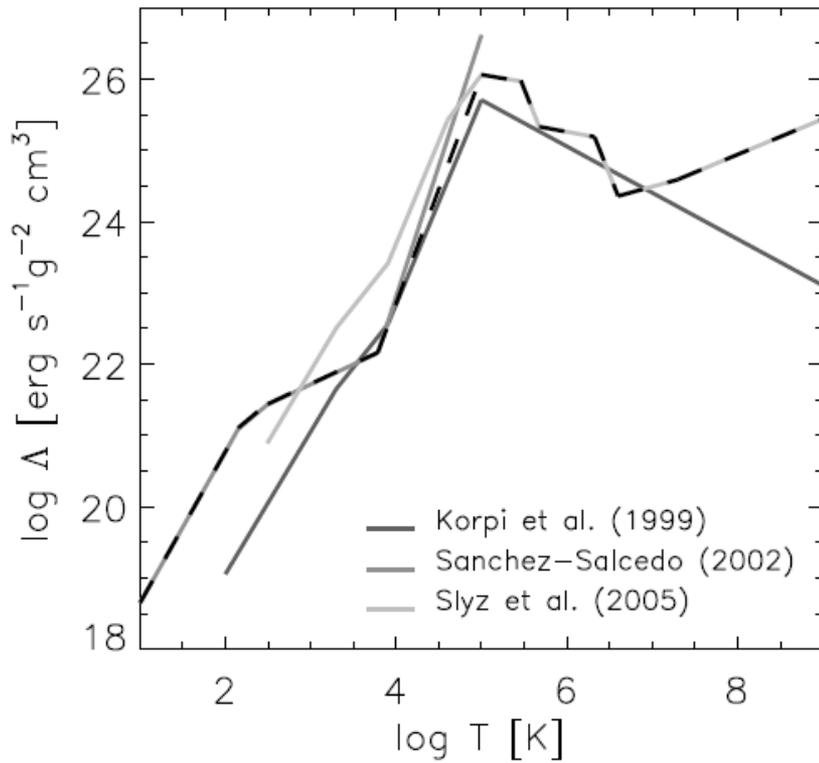
# By maximum evolution 50%



# By maximum evolution 100%



# Equilibrium curve and cooling function



$$\lambda_F = 2\pi \left[ \frac{\rho^2 \Lambda}{\kappa T} (1 - \beta) \right]^{-1/2}$$

Koyama & Inutsuka (2004)

$$P_m = 2.5$$

$$\eta = 0.2 \times 10^{25} \text{ cm}^2 \text{ s}^{-1}$$

$$\Gamma(z) = \Gamma_0 \times \begin{cases} e^{-\frac{z^2}{2z_0 H_\Gamma}} & \text{if } |z| \leq z_0 \\ e^{\frac{z_0}{2H_\Gamma}} (e^{-\frac{z}{H_\Gamma}} + 10^{-5}) & \text{otherwise,} \end{cases}$$

$$H_\Gamma = 300 \text{ pc and choose } z_0 = H_\Gamma/5.$$

$$\Gamma_0 = 0.015 \text{ erg s}^{-1}$$

photoelectric heating and ionising radiation

from OB stars.

Joung & Mac Low (2006)

$$\Lambda(T) = \Lambda_i T^{\beta_i}$$

$T_i$ [K]	$\Lambda_i$ [erg s <sup>-1</sup> g <sup>-2</sup> cm <sup>3</sup> K <sup>-<math>\beta_i</math>]</sup>	$\beta_i$
10	$3.420 \times 10^{16}$	2.12
141	$9.100 \times 10^{18}$	1.00
313	$1.110 \times 10^{20}$	0.56
6102	$1.064 \times 10^{10}$	3.21
$10^5$	$1.147 \times 10^{27}$	-0.20
$2.88 \times 10^5$	$2.290 \times 10^{42}$	-3.00
$4.73 \times 10^5$	$3.800 \times 10^{26}$	-0.22
$2.11 \times 10^6$	$1.445 \times 10^{44}$	-3.00
$3.98 \times 10^6$	$1.513 \times 10^{22}$	0.33
$2.00 \times 10^7$	$8.706 \times 10^{20}$	0.50

Sánchez-Salcedo, Vázquez-Semadeni & Gazol (2002)

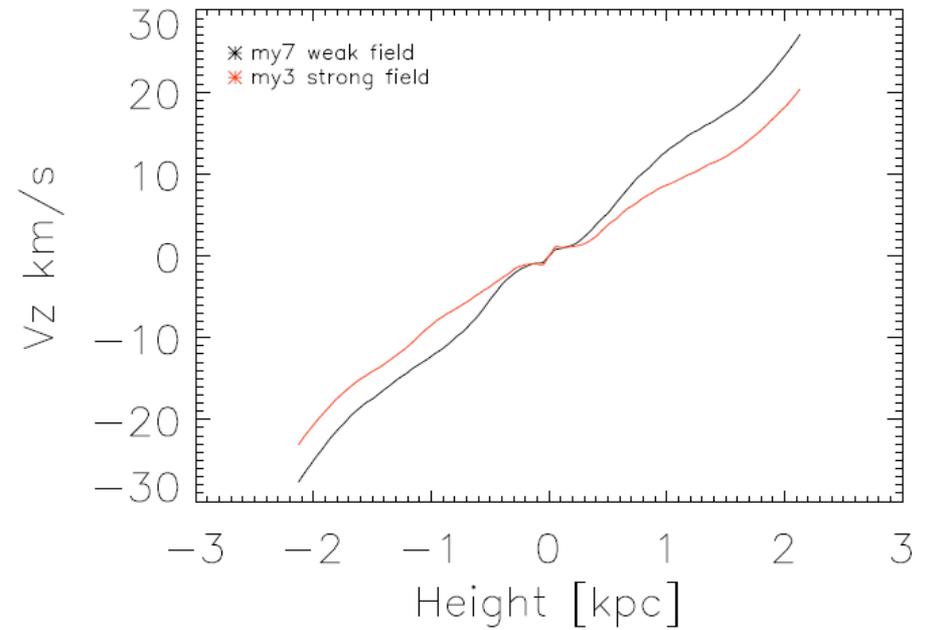
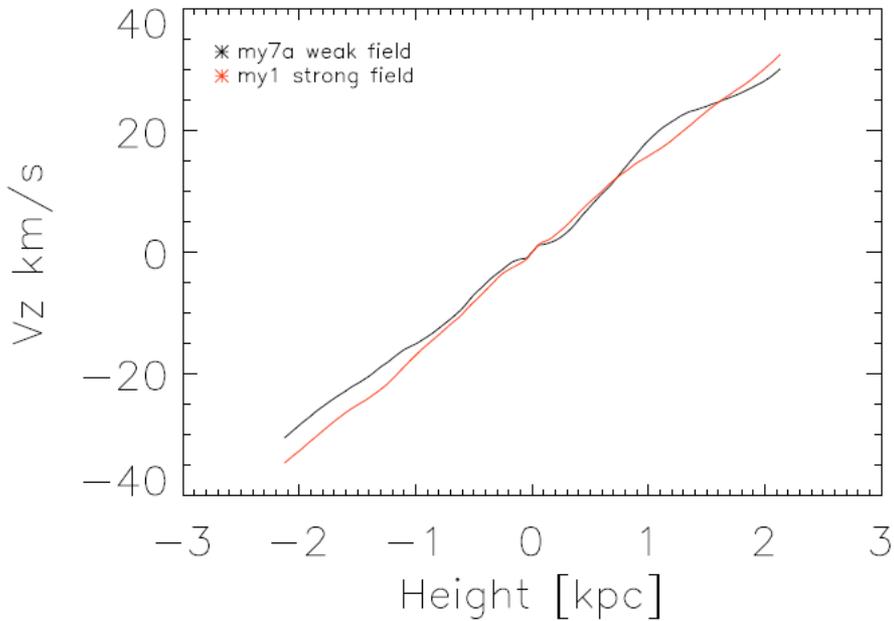
# Flux dependence

	<b>M1</b>	<b>M3</b>	<b>M7</b>	<b>M7a</b>	<b>Gressel</b>
Time (Gyr)	1.2	1.6	2.5	2.5	1.9
Flux ( $\mu\text{Gcm}^2$ )	0	$10^{41}$	$10^{39}$	0	0
SN rate (%)	25%	25%	25%	25%	25%
$B_{\text{seed}}$ ( $\mu\text{G}$ )	0.1	0.1	0.001	0.001	0.001 ( $B_{\gamma}$ )
Growth time (Myr)	308	420	200	191	254
$E_{\text{mag}} : E_{\text{kin}}$ (final)	0.4	1.3	2.3	0.5	0.6
$\langle B \rangle U$ ( $\mu\text{G}$ )	-0.3	-1.8	-1	0.3	0.07
$\langle B \rangle L$ ( $\mu\text{G}$ )	0.9	1.8	-1	0.07	0.3
$B_{\text{mean}} : B_{\text{rms}}$	1.5	2	2	2.4	1.7

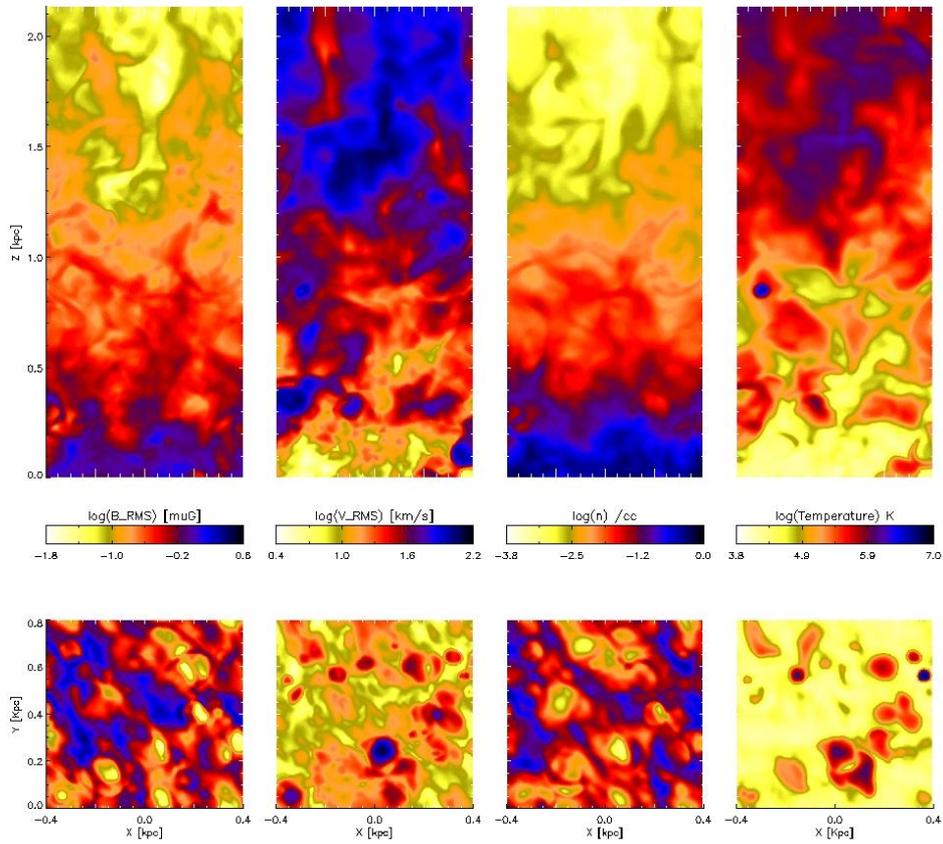
# SNR dependence

	<b>M7</b>	<b>M71</b>	<b>M72</b>
Time (Gyr)	2.5	1.8	1.5
Flux ( $\mu\text{Gcm}^2$ )	$10^{39}$	$10^{39}$	$10^{39}$
SN rate (%)	25%	50%	100%
$B_{\text{seed}}$ ( $\mu\text{G}$ )	0.001	0.001	0.001
Growth time (Myr)	200	190	179
$E_{\text{mag}} : E_{\text{kin}}$	2.3	1.1	0.4
$\langle B \rangle_U$ ( $\mu\text{G}$ )	-1	0.6	0.3
$\langle B \rangle_L$ ( $\mu\text{G}$ )	-1	0.6	0.3
$B_{\text{mean}} : B_{\text{rms}}$	2	1.7	1

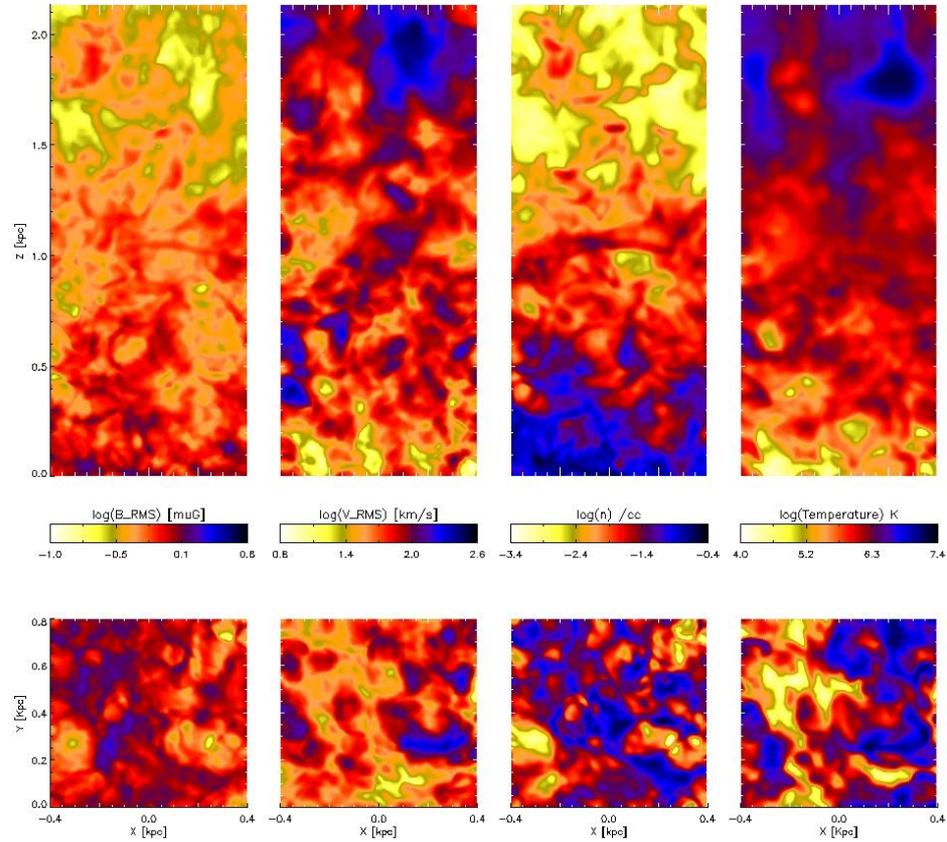
# Vertical wind



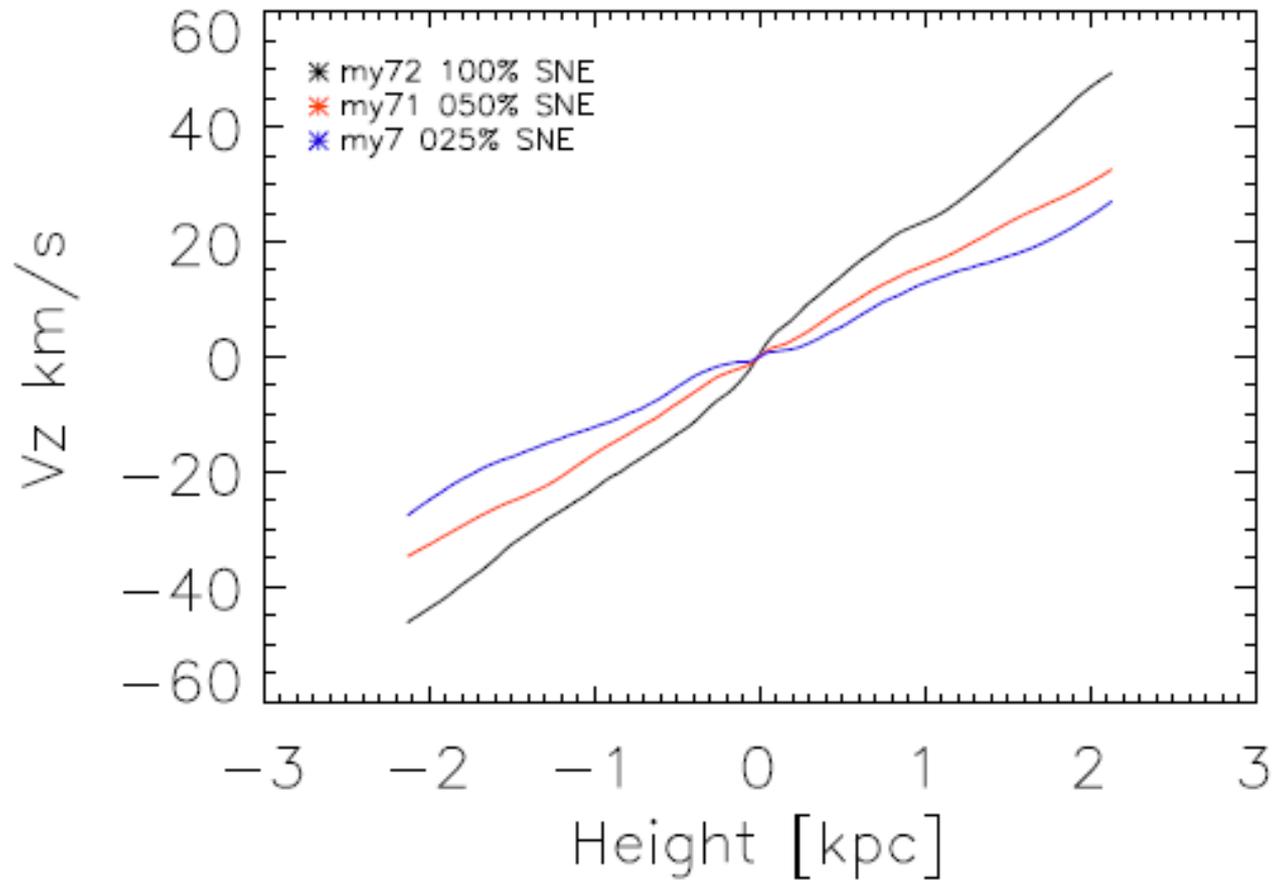
# My71



# My72



# Vertical wind (SNR dependence)



# Model equations

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{v}) = 0$$

$$\begin{aligned} \partial_t (\rho \mathbf{v}) + \nabla \cdot [ \rho (\mathbf{v} \otimes \mathbf{v}) + p^* \mathbf{I} - (\mathbf{B} \otimes \mathbf{B}) ] = \\ - 2\Omega \rho \boldsymbol{\epsilon} \times \mathbf{v} + 2\Omega^2 q x \hat{x} + \Omega g(z) \hat{z} + \nabla \cdot \boldsymbol{\tau} \end{aligned}$$

$$\begin{aligned} \partial_t \boldsymbol{\epsilon} + \nabla \cdot [ (e + p^*) \mathbf{v} - \boldsymbol{\epsilon} \cdot \mathbf{B} \mathbf{B} ] = \\ 2\Omega^2 q x \boldsymbol{\epsilon} \cdot \mathbf{v} + \Omega g(z) \boldsymbol{\epsilon} \cdot \mathbf{v} \\ + \nabla \cdot [ \mathbf{v} + \eta \boldsymbol{\epsilon} \times \mathbf{B} ] + \kappa \nabla T \\ + \Gamma_{SN} - \rho^2 \Lambda \boldsymbol{\epsilon} + \rho \Gamma \boldsymbol{\epsilon} \end{aligned}$$

$$\partial_t \mathbf{B} - \nabla \times \boldsymbol{\epsilon} \times \mathbf{B} - \eta \nabla \times \mathbf{B} = \bar{0}$$

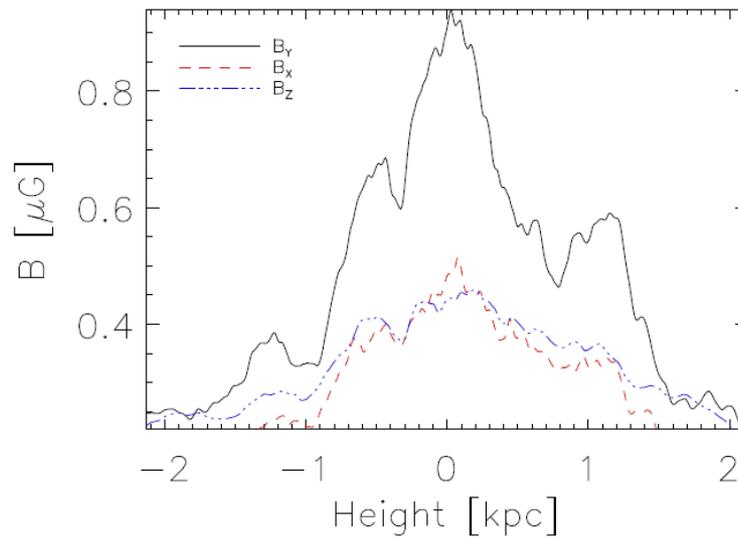
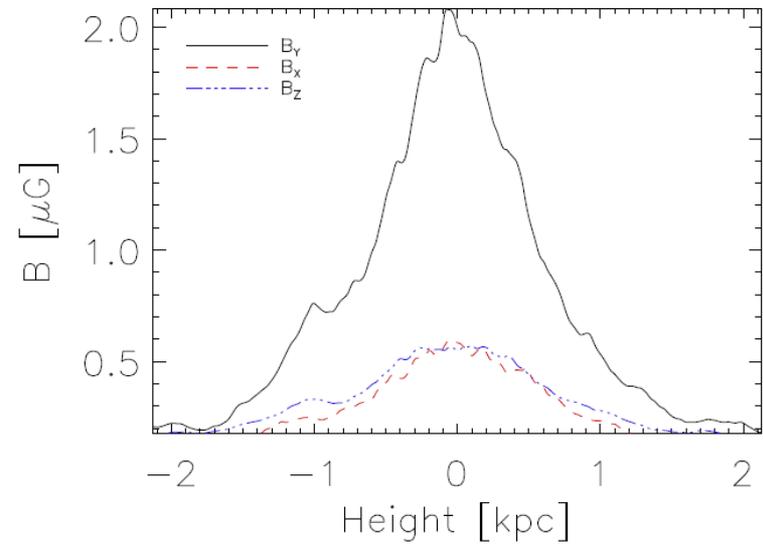
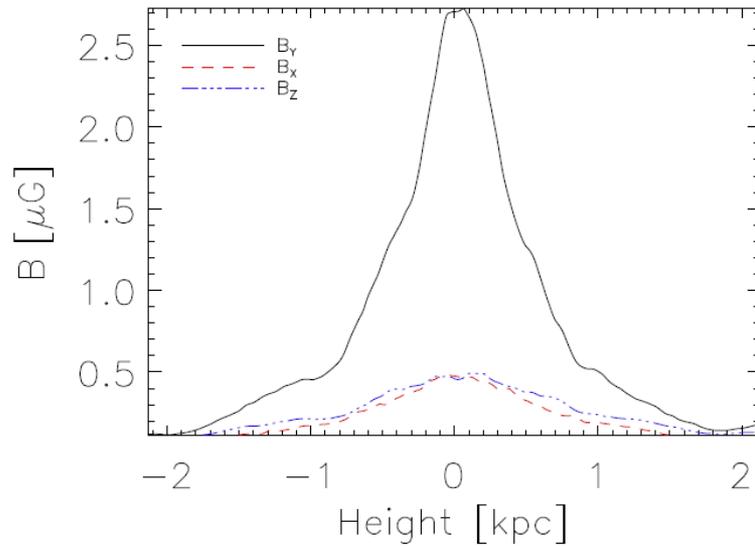
Where,  $p^* = p + B^2/2$

Thermal energy,  $\varepsilon = e - \rho v^2/2 - B^2/2$

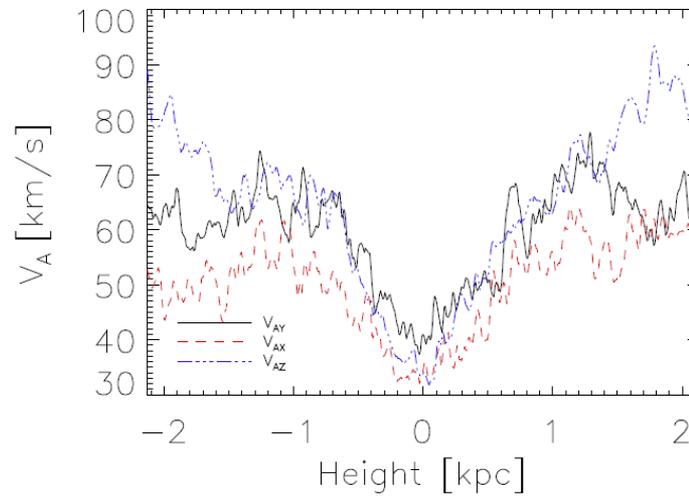
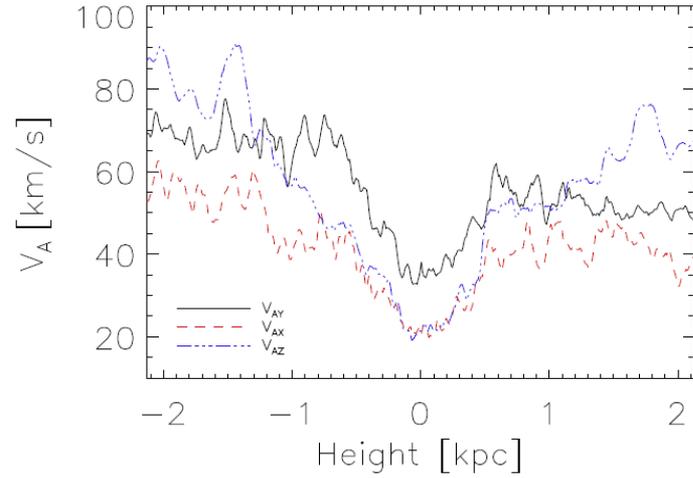
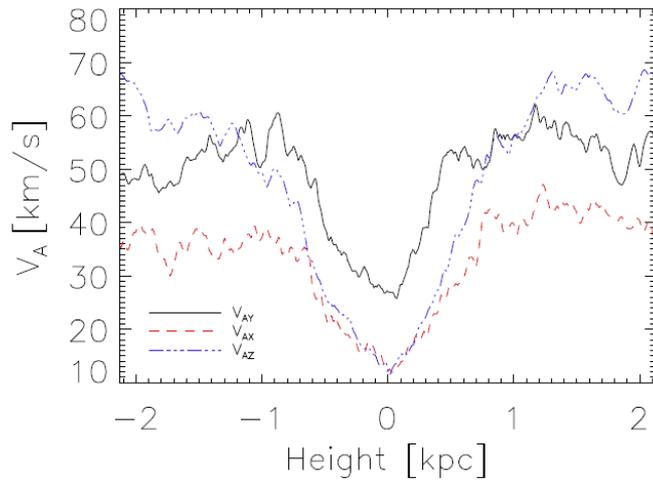
Viscous stress tensor,  $\tau = \tilde{\nu} \left[ \bar{\nabla} \otimes \bar{v} + (\bar{\nabla} \otimes \bar{v})^T - \frac{2}{3} (\bar{\nabla} \cdot \bar{v}) \bar{I} \right]$

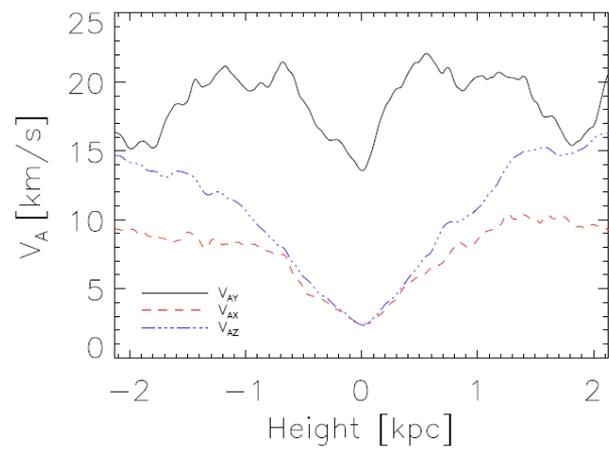
Shear parameter,  $q = d \ln \Omega / d \ln R = -1$

# Mean $|B|$

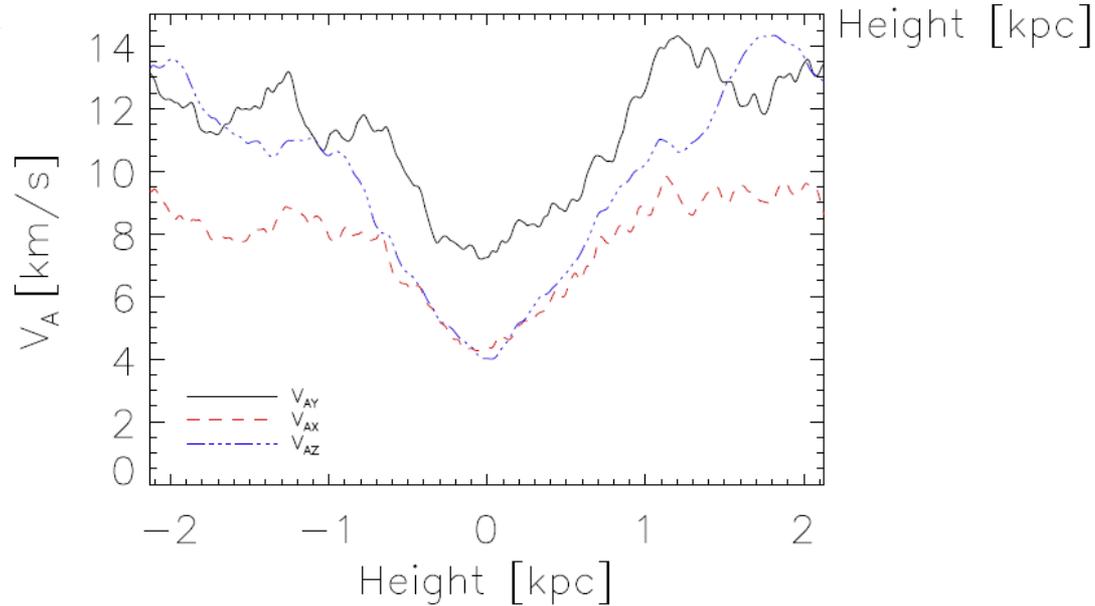
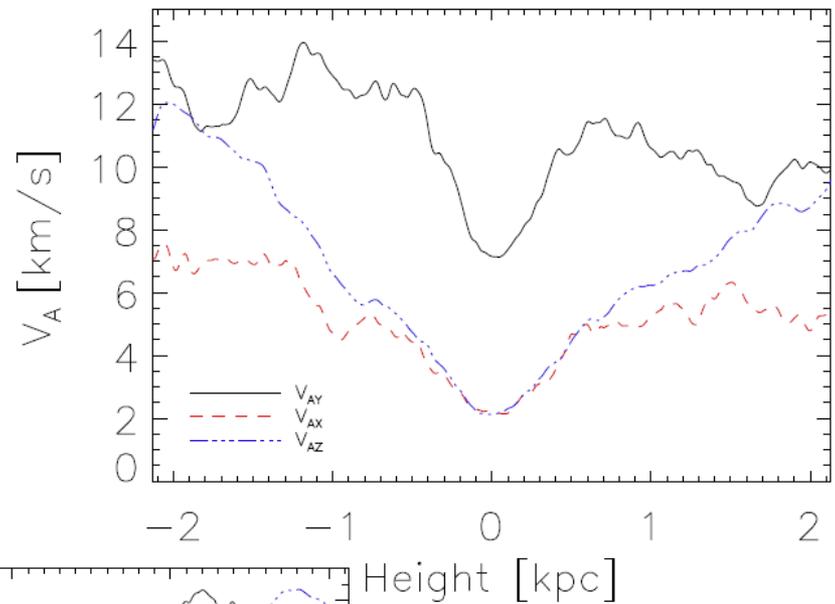
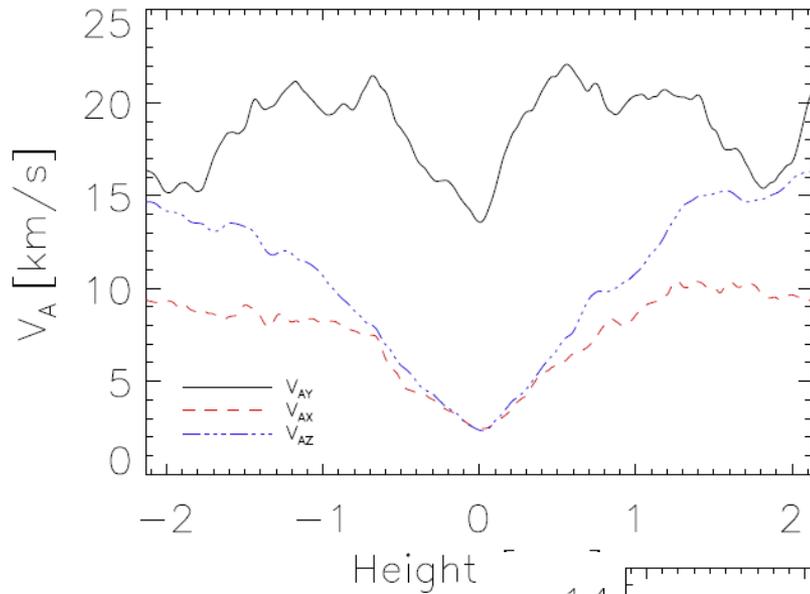


# Maximum $|V_A|$

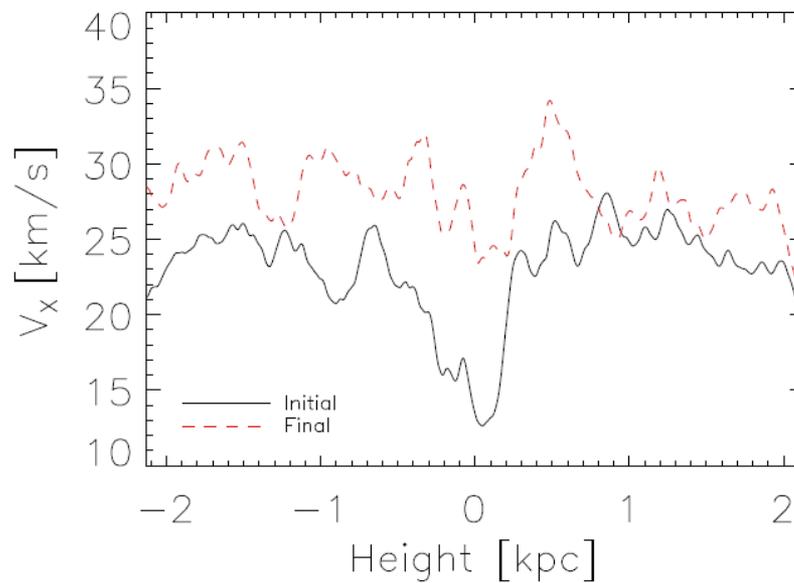
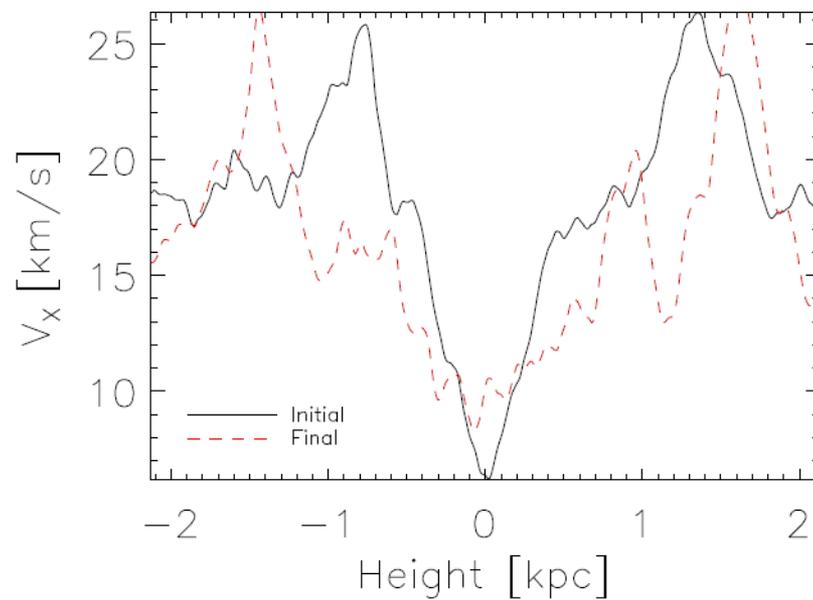
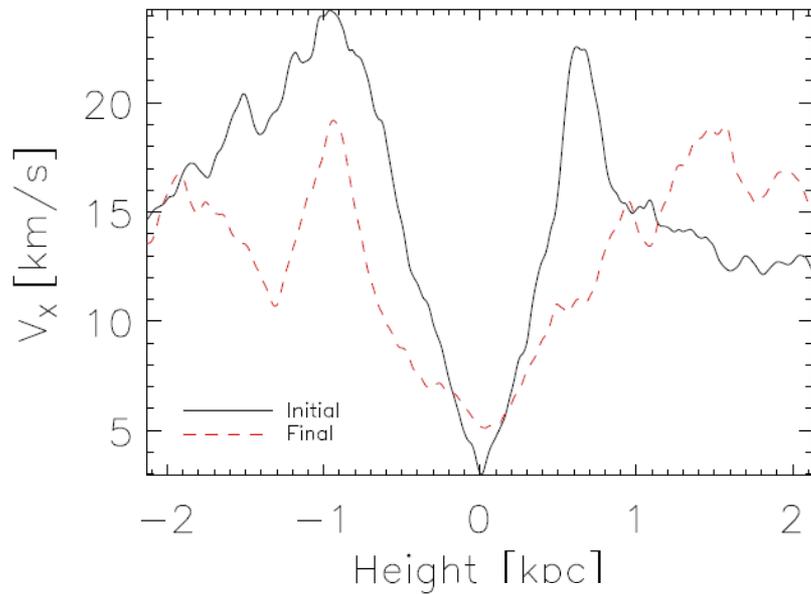




# Mean $|V_A|$



# Mean $|V_x|$



# Mean $|V_z|$

