

Cosmic Rays & Galactic Winds

Lecture given at the Summer School: "Magnetic Fields: From Star-forming Regions to Galaxy Clusters and Beyond"

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Dieter Breitschwerdt - Ringberg Summer School "Magnetic Fields" - 18.7.2011



Lecture Overview

- ★ Introduction
- ★ Cosmic Rays
- ★ Galactic Winds
- ★ Theory & Models
- ★ Self-consistent modeling of galactic outflows & ionization structure
- ★ Electron transport in NGC 891 and NGC 253
- ★ Diffuse radial γ -ray gradient in the Milky Way
- ★ Cosmic Ray Acceleration Beyond the “Knee”
- ★ Summary

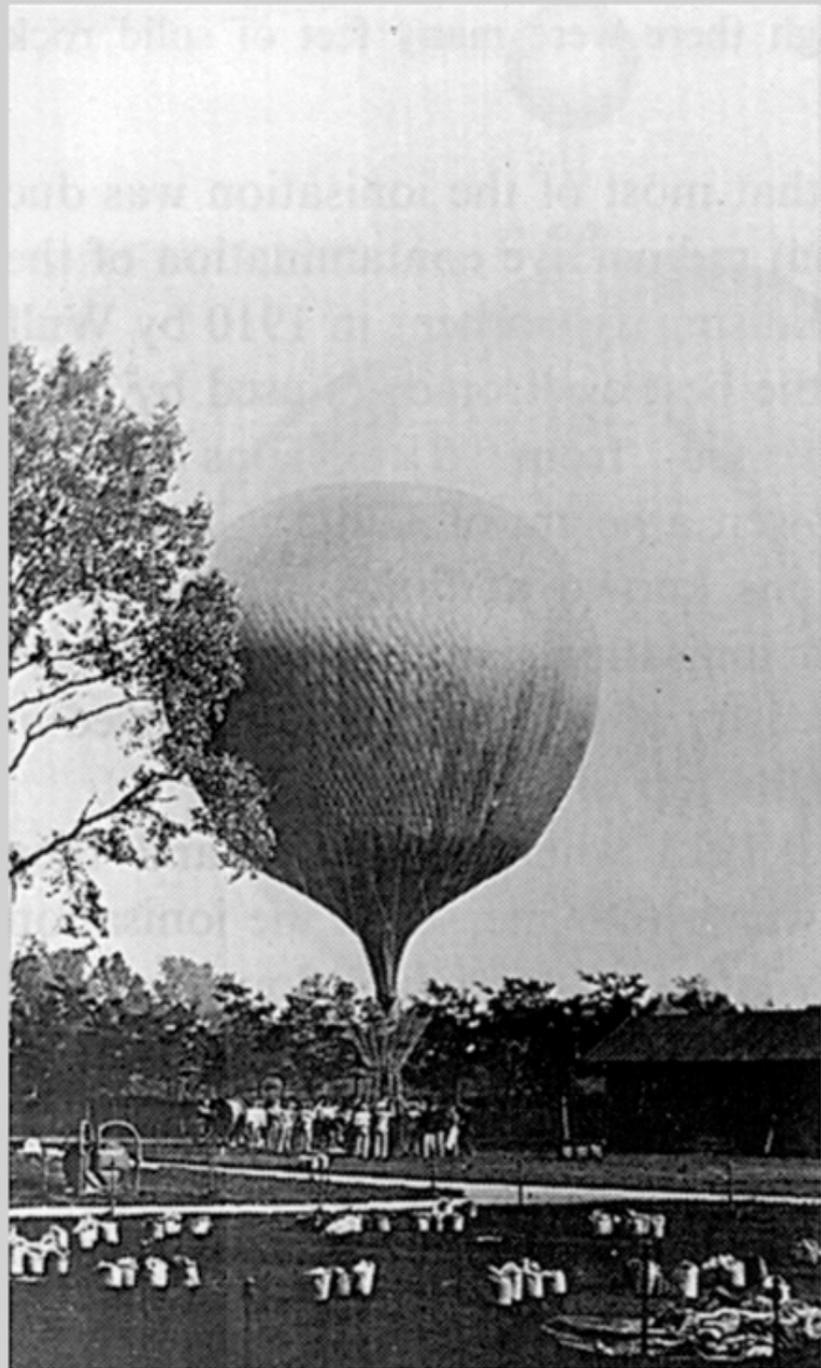
Cosmic Rays

Cosmic Radiation

Includes -

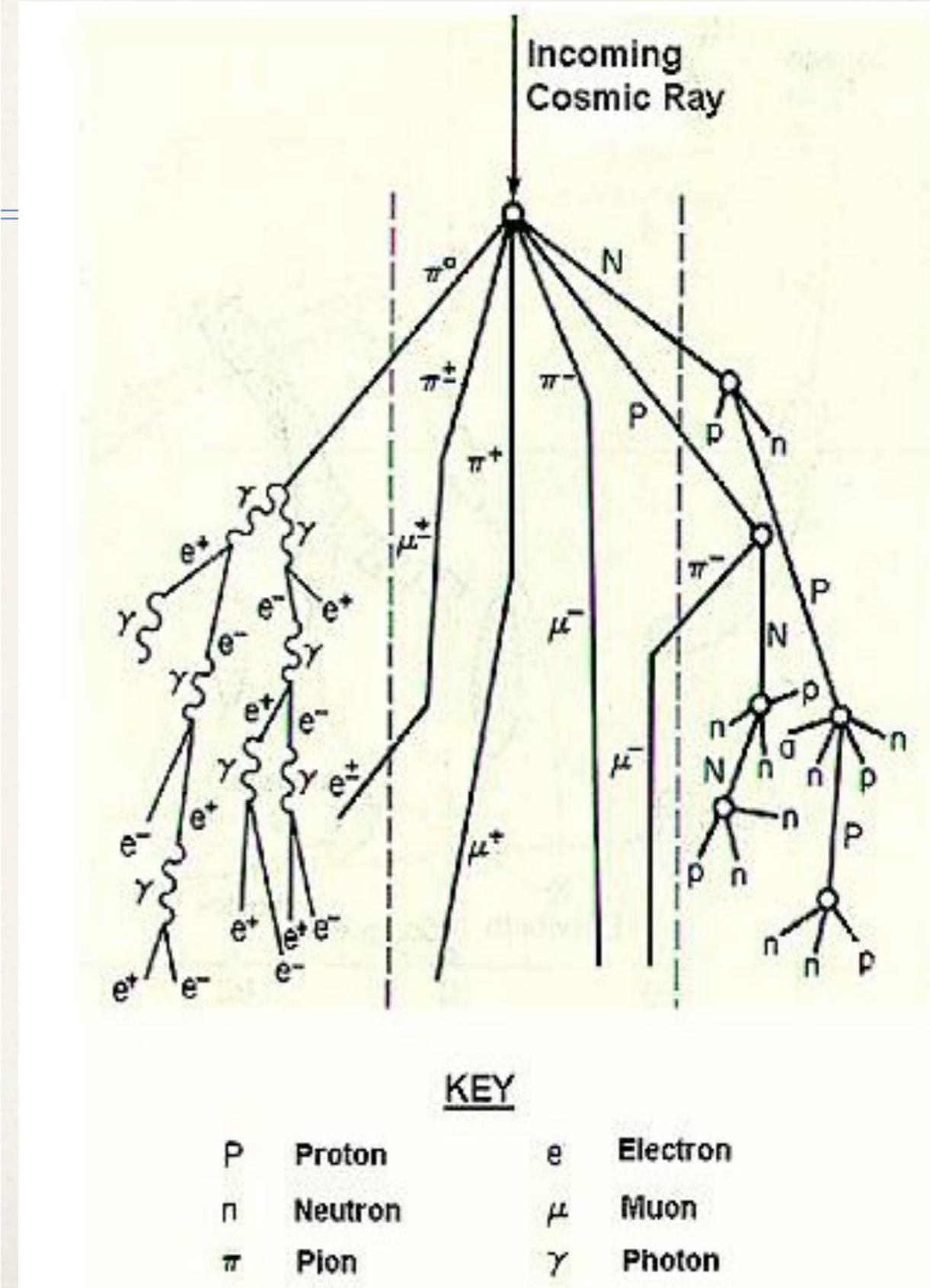
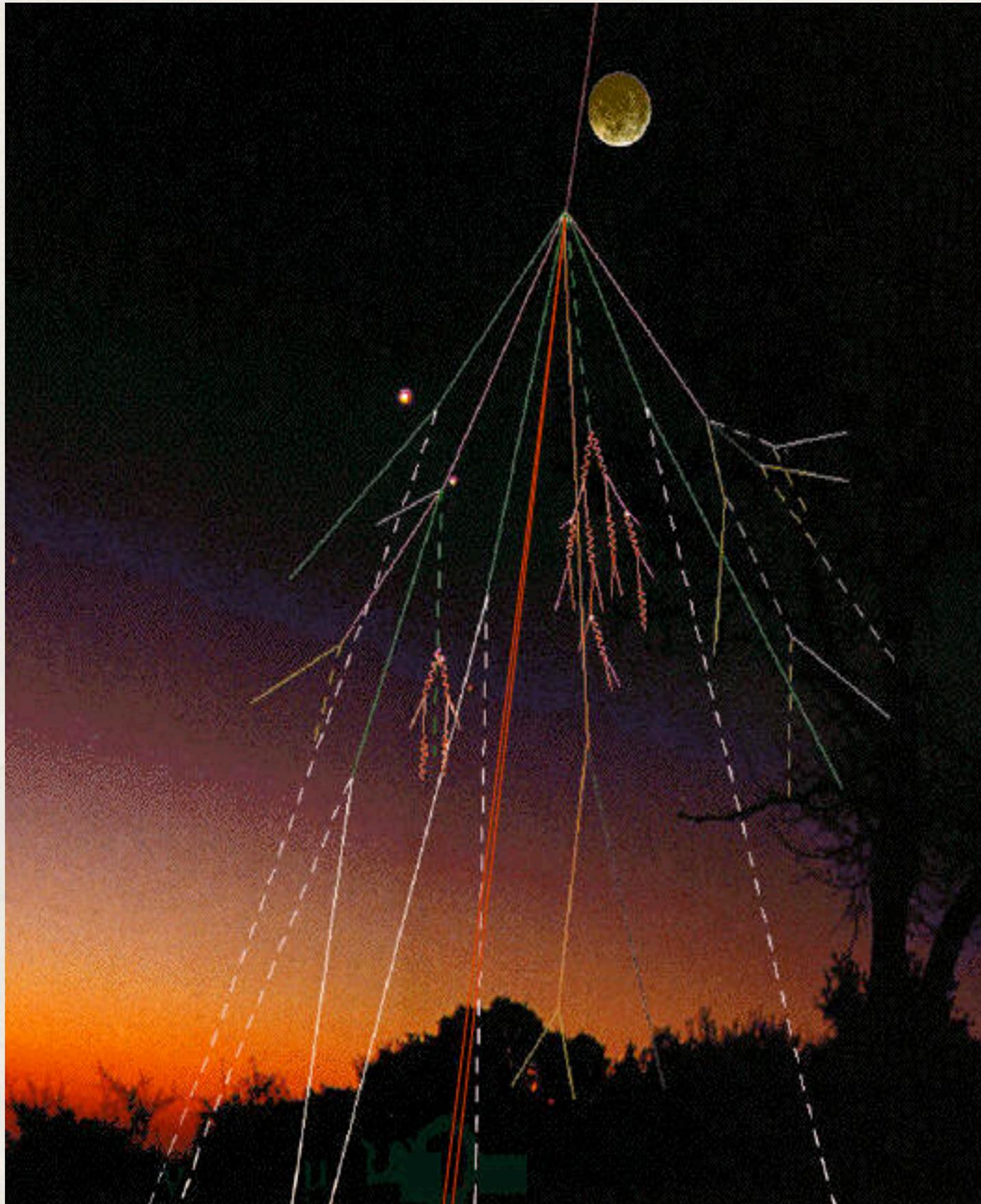
- **Particles** (2% electrons, 98% protons and atomic nuclei)
- **Photons**
- **Large energies** ($10^9 \text{ eV} \leq E \leq 10^{20} \text{ eV}$)
 γ -ray photons produced in collisions of high energy particles

Extraterrestrial Origin



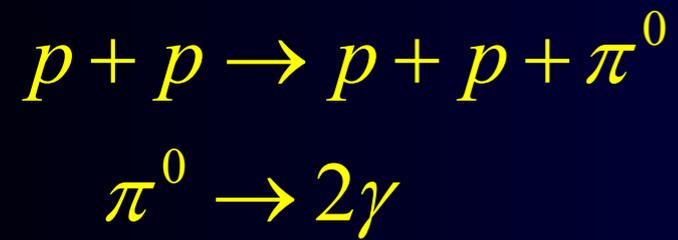
- Increase of ionizing radiation with altitude
- 1912 Victor Hess' balloon flight up to 17500 ft. (without oxygen mask!)
- Used gold leaf electroscope

Cosmic Showers



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Inelastic collision of CRs with ISM



For $E > 100$ MeV dominant process

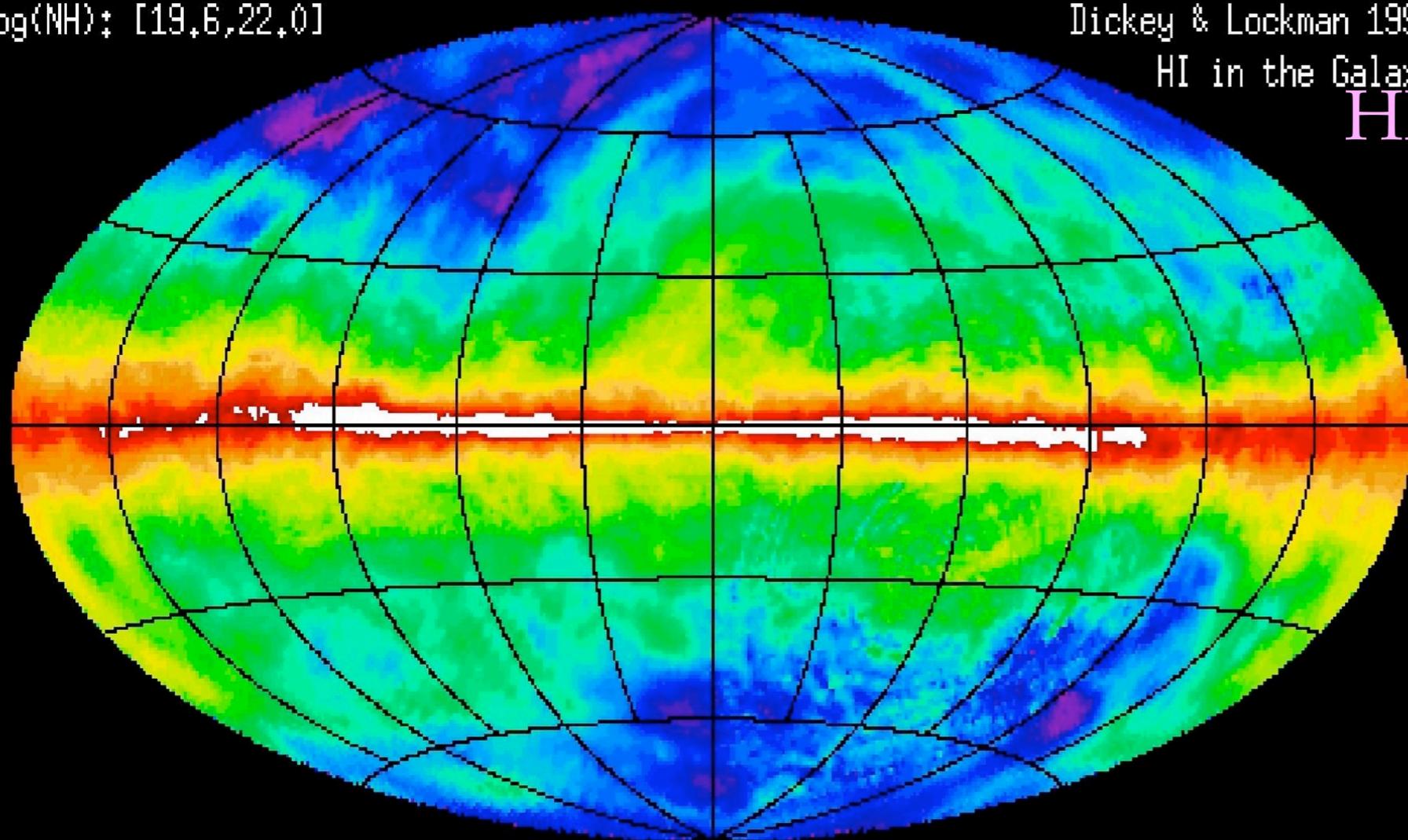
EGRET All-Sky Gamma Ray Survey Above 100 MeV

$\log(NH): [19.6, 22.0]$

Dickey & Lockman 1990

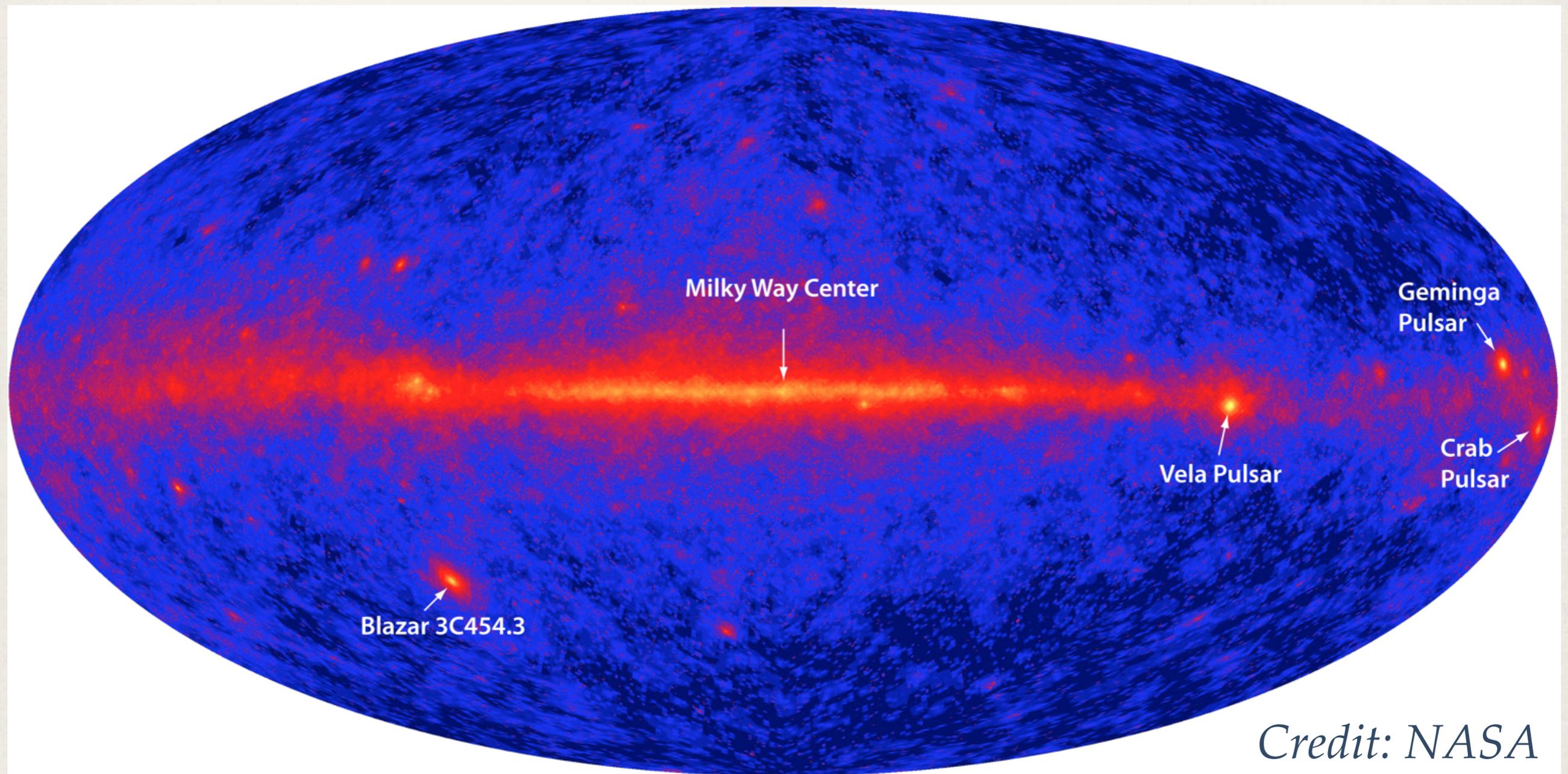
HI in the Galaxy

HI



- Diffuse γ -emission maps the Galactic HI distribution
- Discrete sources at high lat.: AGN (e.g. 3C279)

FERMI LAT All-Sky Map



Credit: NASA

Gamma-ray All-Sky Map (1 year; launch: 11.6.08): 20 MeV - 300 GeV (LAT)

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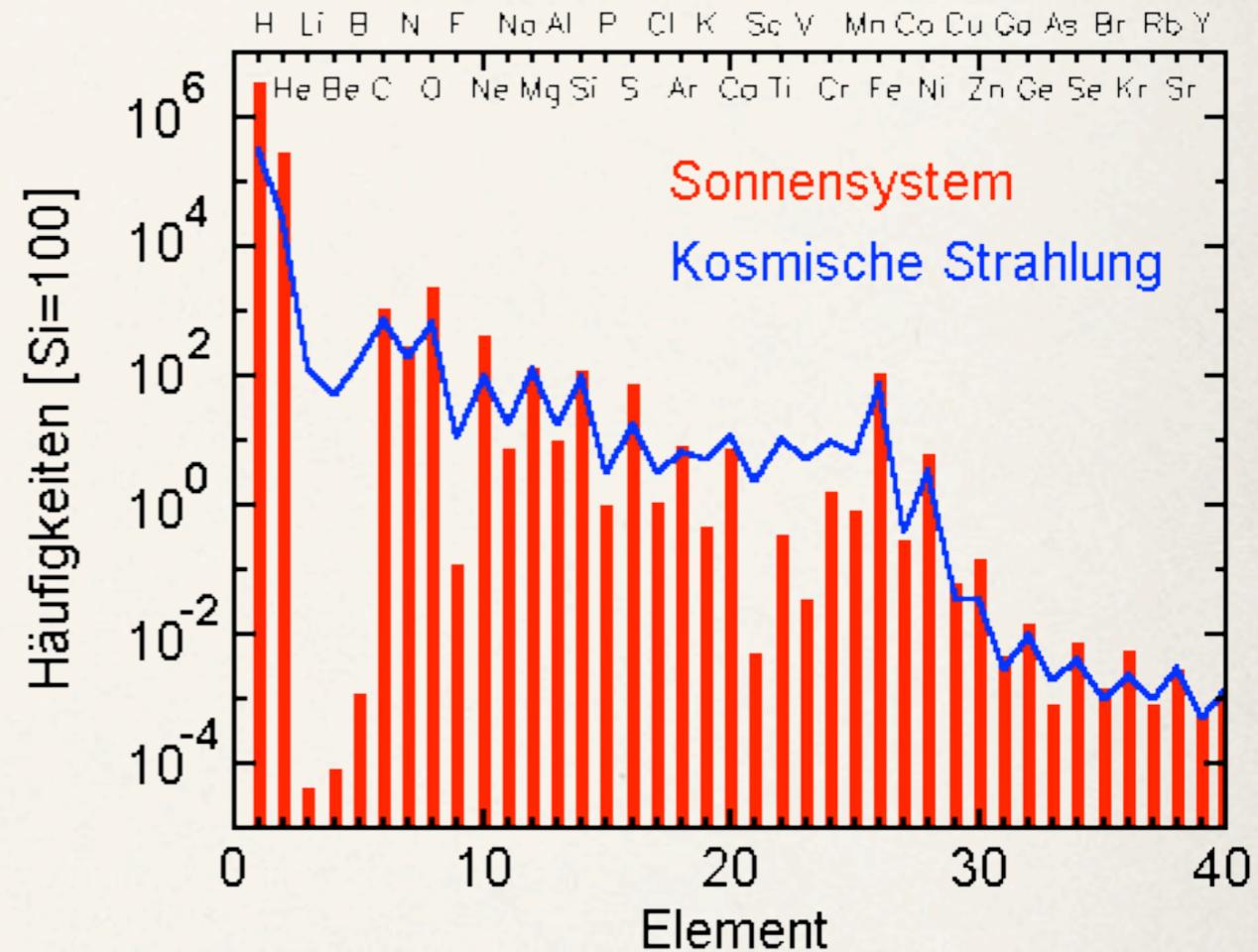
Chemical composition

<u>Groups of nuclei</u>	<u>Z</u>	<u>CR</u>	<u>Universe</u>
Protons (H)	1	700	3000
α (He)	2	50	300
Light (Li, Be, B)	3-5	1	0.00001*
Medium (C,N,O,F)	6-9	3	3
Heavy (Ne->Ca)	10-19	0.7	1
V. Heavy	>20	0.3	0.06

Note: Overabundance of light elements  spallation!

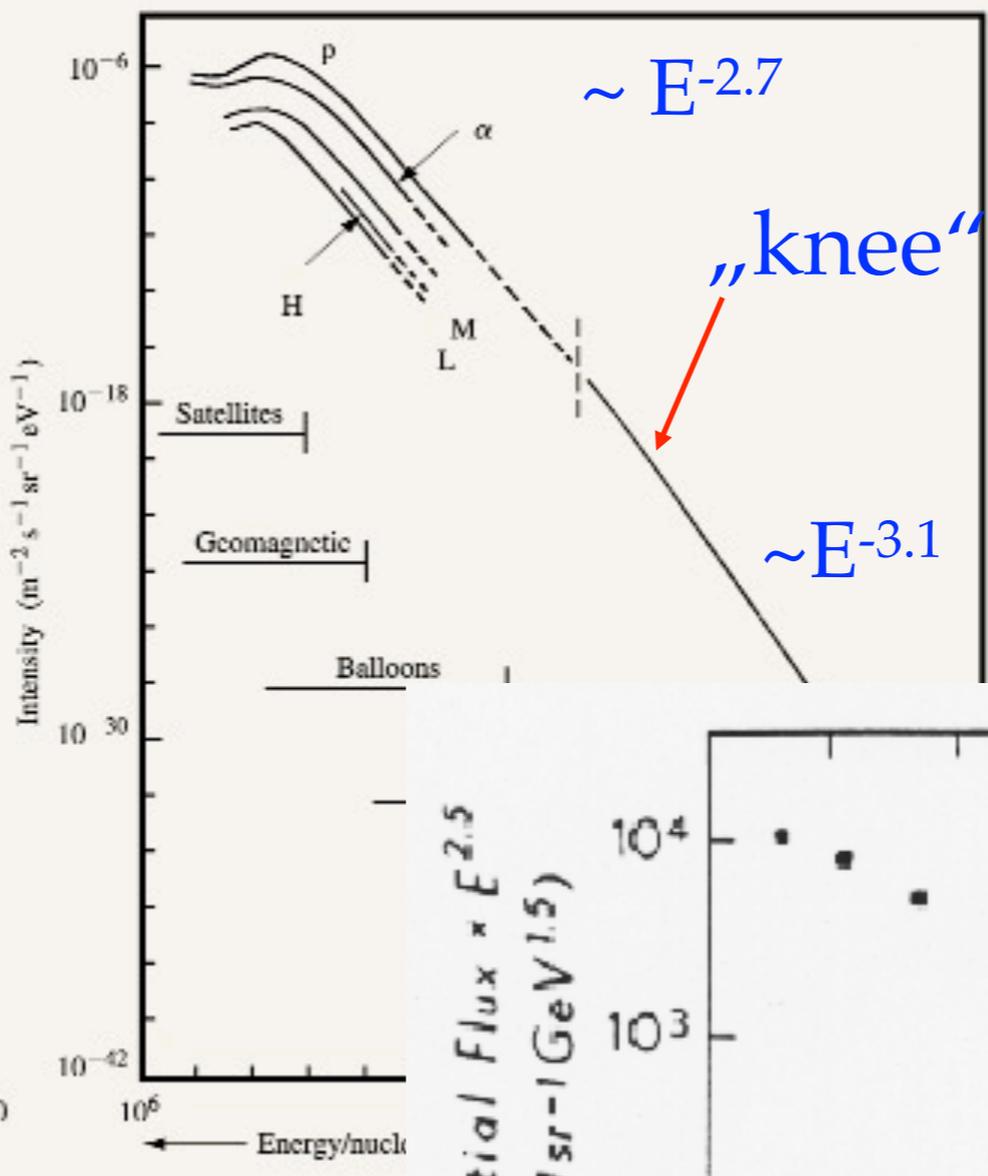
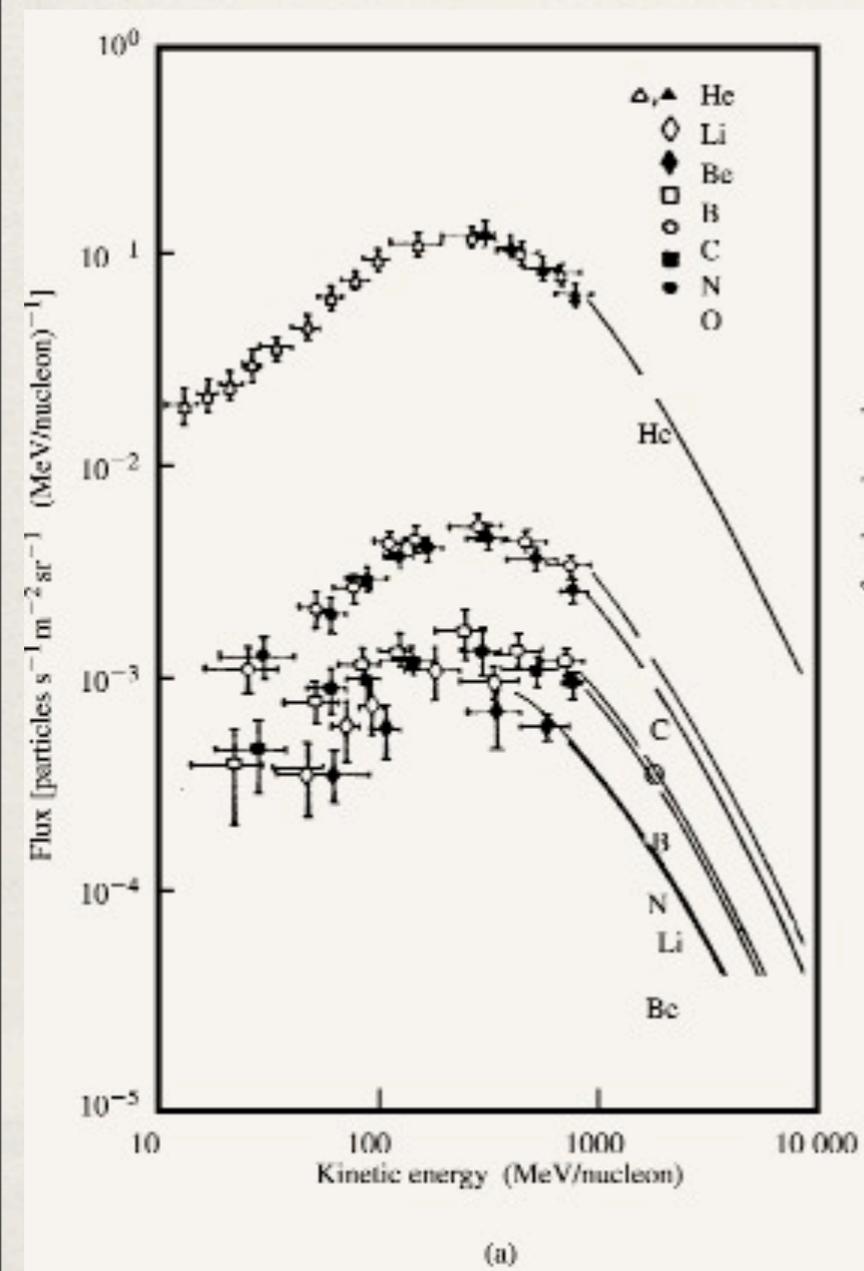
Chemical Composition: Origin of light elements

- Chemical Composition similar to ISM → acceleration of interstellar matter
- Over-abundance of light elements caused by fragmentation of ISM particles in inelastic collision with CR primaries → CR secondaries, e.g. Li, Be, B
- Sc, V, Ti, Cr as spallation products of iron nuclei
- Use fragmentation probabilities and calculate transfer equations by taking into account all possible channels

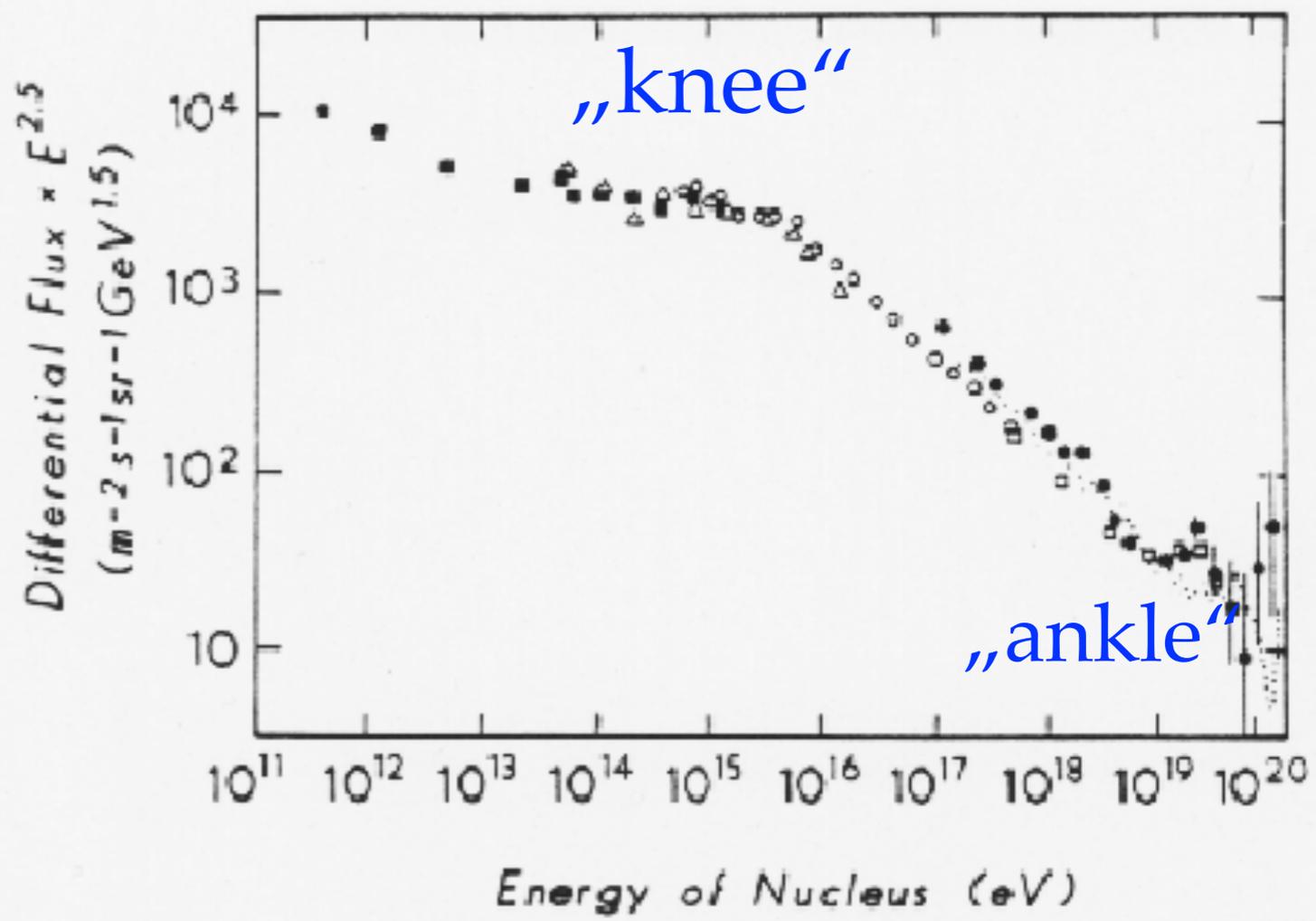


GSFC/NASA: Chemical composition of galactic Cosmic Rays http://imagine.gsfc.nasa.gov/docs/science/know_12/cosmic_rays.html

Differential Energy Spectrum



- differential energy spectrum power law
- 10⁹ < E < 10¹⁵ eV: -2.7
- 10¹⁶ < E < 10¹⁹ eV: -3.1



- for E > 10¹⁹ eV CRs are extragalactic (r_g ~ 3 kpc for 10¹⁹ eV protons at B ~ 1 μG)

Primary CR energy spectrum

- Power law spectrum for $10^9 \text{ eV} < E < 10^{15} \text{ eV}$:

$$I_N(E) \propto E^{-\gamma} \quad \text{with } \gamma \approx 2.70 \quad \text{or } N(E)dE = KE^{-\gamma} dE$$

- $[I_N] = \text{cm}^{-2} \text{ s}^{-1} \text{ sr}^{-1} (\text{GeV} / \text{nucleon})^{-1}$
- Steepening for $E > 10^{15} \text{ eV}$ with $\gamma = 3.08$ („knee“)
- becoming shallower for $E > 10^{18} \text{ eV}$ („ankle“)
- Below $E \sim 10^9 \text{ eV}$ CR intensity drops due to solar modulation (magnetic field inhibits particle streaming)

- gyroradius:
$$r_g = \frac{\gamma_L m_0 v \sin \vartheta}{ZeB} = \left(\frac{pc}{Ze} \right) \frac{\sin \vartheta}{Bc} = R \frac{\sin \vartheta}{Bc}$$

- R ... rigidity, θ ... pitch angle

Example: CR with $E=1 \text{ GeV}$ has $r_g \sim 10^{12} \text{ cm}$! For $B \sim 1 \mu\text{G}$ @ 10^{15} eV , $r_g \sim 0.3 \text{ pc}$

Important CR facts:

- **CR Isotropy:**

– Energies $10^{11} \text{ eV} < E < 10^{15} \text{ eV}$: $\frac{\delta I}{I} \approx 6 \times 10^{-4}$ (anisotropy)

consistent with CRs streaming away from Galaxy

– Energies $10^{15} \text{ eV} < E < 10^{19} \text{ eV}$:

anisotropy increases \rightarrow particles escape more easily
(energy dependent escape)

Note: @ 10^{19} eV , $r_g \sim 3 \text{ kpc}$

– Energies $E > 10^{19} \text{ eV}$: CRs from Local Supercluster?

particles cannot be confined to Galactic disk

- **CR clocks:**

– CR secondaries produced in spallation (from O and C) such as ^{10}Be have half life time $\tau_H \sim 1.6 \text{ Myr}$ \rightarrow β -decay into ^{10}B

- From amount of ^{10}Be relative to other Be isotopes and ^{10}B and τ_{H} the mean CR residence time can be estimated to be

$$\tau_{\text{esc}} \sim 2 \cdot 10^7 \text{ yr for a 1 GeV nucleon}$$

→ **CRs have to be constantly replenished!**

What are the sources?

- Detailed quantitative analysis of amount of **primaries and secondary** spallation products yields a mean Galactic mass traversed („*grammage*“ x) as a function of rigidity R :

$$x(R) = 6.9 \left(\frac{R}{20 \text{ GV}} \right)^{-\xi} \text{ g/cm}^2, \quad \xi = 0.6$$

for 1 GeV particle, $x \sim 9 \text{ g/cm}^2$

- Mean measured CR energy density:

$$\varepsilon_{\text{CR}} \sim \varepsilon_{\text{mag}} \sim \varepsilon_{\text{th}} \sim \varepsilon_{\text{turb}} \cong 1 \text{ eV/cm}^3$$

- If all CRs were **extragalactic**, an extremely high energy production rate would be necessary (more than AGN and radio galaxies could produce) to sustain high CR background radiation
 - assuming **energy equipartition** between B-field and CRs
radio continuum observations of starburst galaxy M82 give $\epsilon_{CR}(M82) \sim 100\epsilon_{CR}(Galaxy)$
 - CR production rate proportional to star formation rate
- *no constant high background level!*
- *CR interact strongly with B-field and thermal gas*

CR propagation

- High energy nucleons are ultrarelativistic -> light travel time from sources

$$\tau_{lc} \sim L/c \approx 3 \times 10^4 \text{ yr} \ll \tau_{esc}$$

- CRs as charged particles *strongly coupled to B-field*
- *B-field: $\vec{B} = \langle \vec{B}_{reg} \rangle + \delta\vec{B}$ with strong fluctuation component $\delta\vec{B}$*
→ MHD (Alfvén) waves

→ Cross field *propagation by pitch angle scattering*

→ random walk of particles!

- CRs **DIFFUSE** through Galaxy with mean speed

$$\langle v_{diff} \rangle \sim L/\tau_r \approx 10 \text{ kpc} / 2 \times 10^7 \text{ yr} = 490 \text{ km/s} \sim 10^{-3} c$$

- Mean gas density traversed by particles

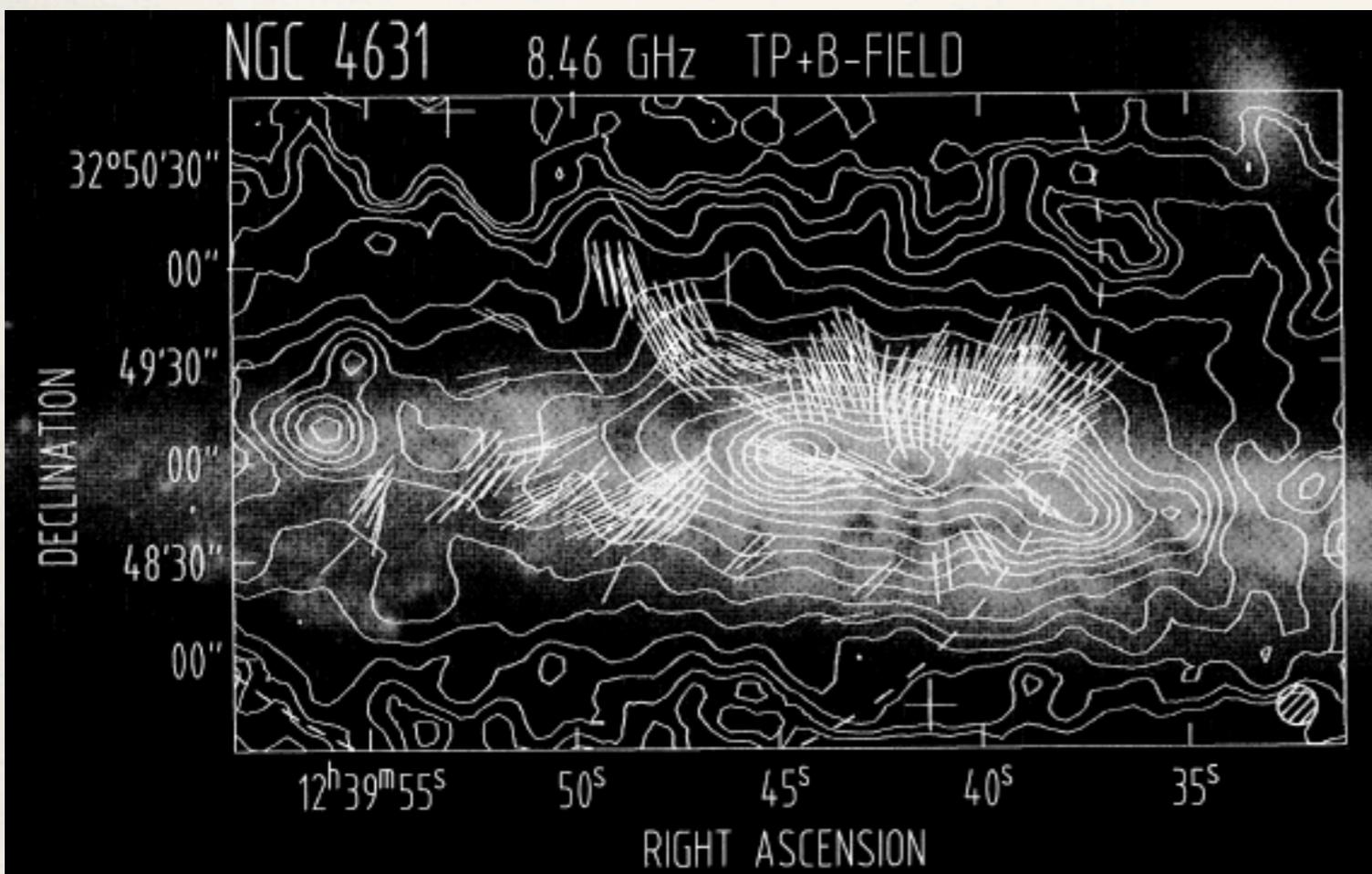
$$\langle \rho_h \rangle \approx x / c\tau_{esc} \sim 5 \times 10^{-25} \text{ g/cm}^3 \sim \frac{1}{4} \langle \rho_{ISM} \rangle$$

- particles spend most time **outside the Galactic disk** in the Galactic *halo!* **→** „*confinement*“ volume
 - CR „height“ ~ 4 times h_g ($=250$ pc) ~ 1 kpc
 - CR diffusion coefficient:

$$\kappa \sim h_{CR} \times L / \tau_{esc} \approx 5 \times 10^{28} \text{ cm}^2 / \text{s}$$
 - Mean free path for CR propagation: $\lambda_{CR} \sim 3\kappa / c \sim 1 \text{ pc}$
→ strong scattering off magnetic irregularities!
- Analysis of radioactive isotopes in meteorites: CR flux roughly constant over last 10^9 years

CR origin

- CR electrons ($\sim 1\%$ of CR particle density) must be of Galactic origin due to strong **synchrotron losses** in Galactic magnetic field and **inverse Compton** losses



- Note: radio continuum observations of edge-on galaxies show halo field component

Golla, G. et al.

- Estimate of total Galactic CR energy flux:

$$F_{CR} \sim \epsilon_{CR} \frac{V_{conf}}{\tau_{esc}} \approx 10^{41} \text{ erg/s}$$

Note: only $\sim 1\%$ radiated away in γ -rays!

- Enormous energy requirements leave as most realistic Galactic CR source supernova remnants (SNRs)
 - Available hydrodynamic energy:

$$F_{SNR} \sim \nu_{SN} E_{SNR} \approx \frac{3}{100 \text{ yr}} \times 10^{51} \text{ erg} \approx 10^{42} \text{ erg/s}$$

→ about 10% of total SNR energy has to be converted to CRs

– Promising mechanism: diffusive shock acceleration

- Ultrahigh energy CRs must be extragalactic

$$r_g \geq 100 \text{ kpc} > R_{gal} \text{ (for } E \sim 10^{20} \text{ eV)}$$

Galactic Winds



- ★ **Galaxies** are essential building blocks of the Universe
- ★ **feedback processes** in the disk and halo become ever more important for their appearance and evolution → **Galactic Cosmic Matter Cycle**
- ★ **star formation** generates hot plasma, “metals”, CRs, B-fields in disk & halo
- ★ no hydrostatic halo → **superbubbles, outflows** (fountain & winds)

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Parker's Wind Theory

$$\frac{1}{r^2} \frac{d}{dr} (\rho u r^2) = q \quad \text{Mass}$$

$$\rho u \frac{du}{dr} = -\frac{dP}{dr} - qu - g \quad \text{Momentum}$$

$$\frac{1}{r^2} \frac{d}{dr} \left[\rho u r^2 \left(\frac{1}{2} u^2 + \frac{\gamma}{\gamma - 1} \frac{P}{\rho} \right) \right] = Q \quad \text{Energy}$$

$$c_g^2 = \gamma \frac{P}{\rho}, \quad P = A(S) \rho^\gamma, \quad M = \frac{u}{c}$$

Assumption:

- steady state flow
- polytropic gas: $P \sim \rho^\gamma$
- spherical geometry
- no sources and sinks, i.e. $q=Q=0$

Wind Equation & Solution

Für $q = Q = 0$ und $g = \frac{GM}{r^2}$

derivation of „wind equation“

$$\frac{1}{u} \frac{du}{dr} = \frac{\frac{2c_g^2}{r} - \frac{GM}{r^2}}{u^2 - c^2}$$

Inner Boundary Condition:

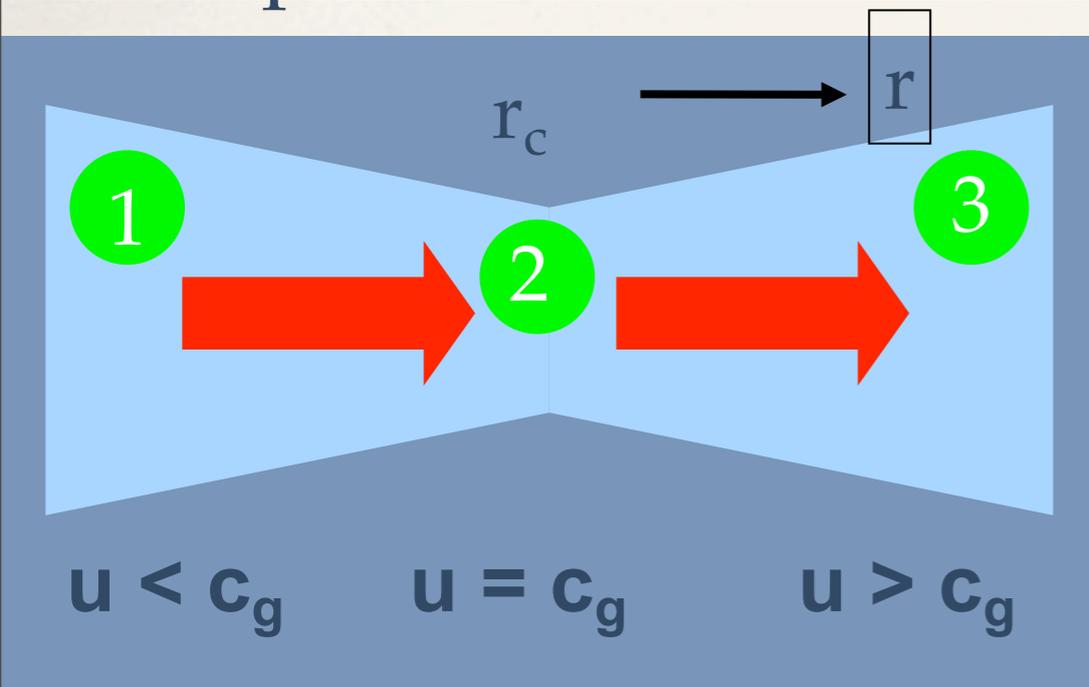
$$r \rightarrow 0 : u \rightarrow 0$$

Outer Boundary Condition:

$$r \rightarrow \infty : P \rightarrow 0$$

smooth transsonic solution

Principle: de Laval Nozzle



For accelerating flow: $\frac{du}{dr} > 0$

$$2 \quad r = r_c = \frac{GM}{2c_g^2} : u = c_g \quad \text{Critical Point}$$

Numerator: $c^2 (1/A) dA/dr$

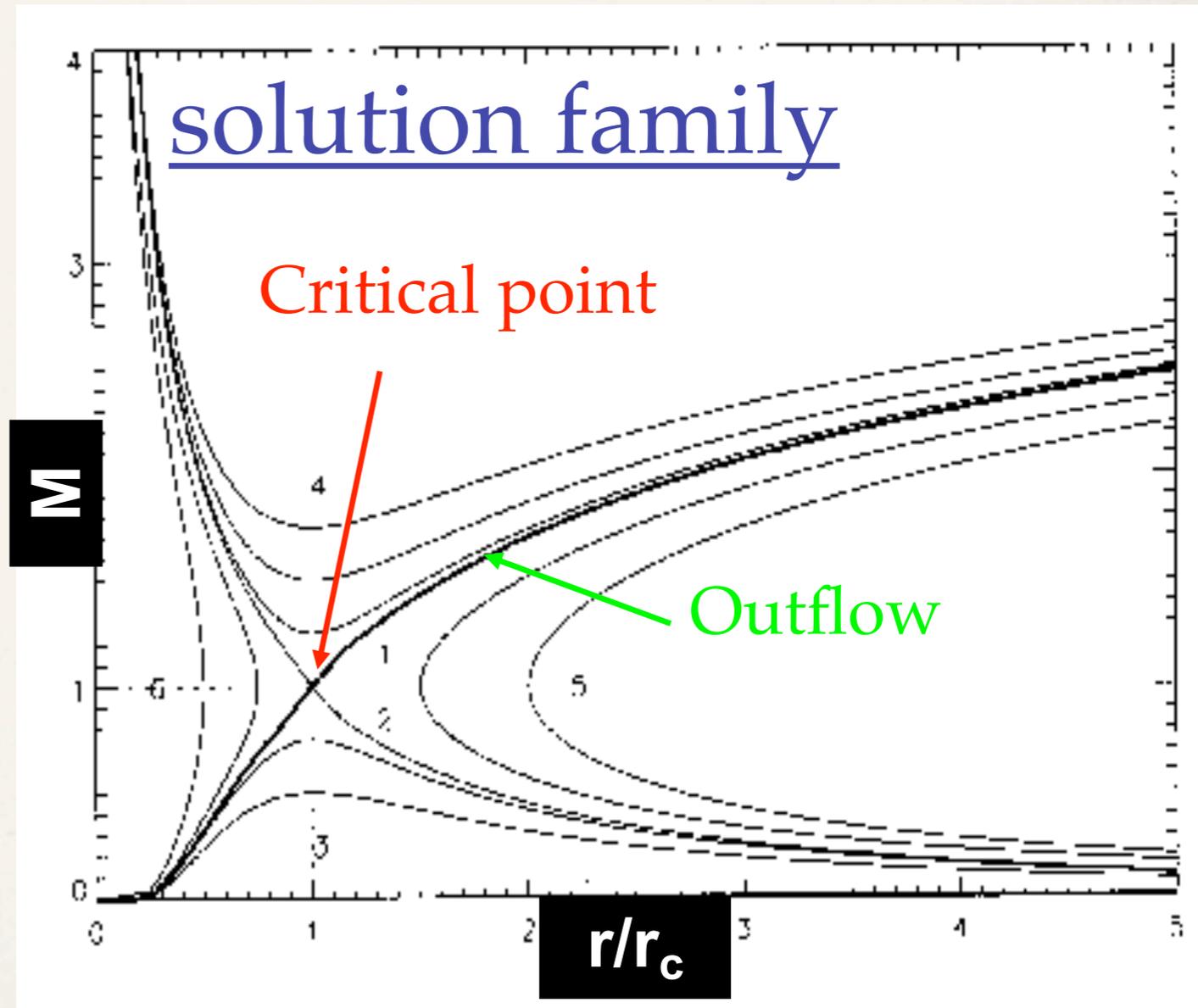
- Equations are integrated inwards and outwards from critical point (Parker's solar wind solution)
- Mass loss rate and terminal velocity are derived from Bernoulli's equation:

$$\frac{1}{2}u^2 + \frac{\gamma}{\gamma-1} \frac{P}{\rho} - \frac{GM}{r} = \text{const.}$$

equating at $r=r_c$ and $r \rightarrow \infty$:

$$\frac{1}{2}c_g^2 + \frac{c_g^2}{\gamma-1} - \frac{GM}{r_c} \approx \frac{1}{2}u_\infty^2$$

$$\Rightarrow u_\infty = \sqrt{\frac{5-3\gamma}{\gamma-1}} c_g, \quad \dot{m} = 4\pi r_c^2 \rho_c c_g$$

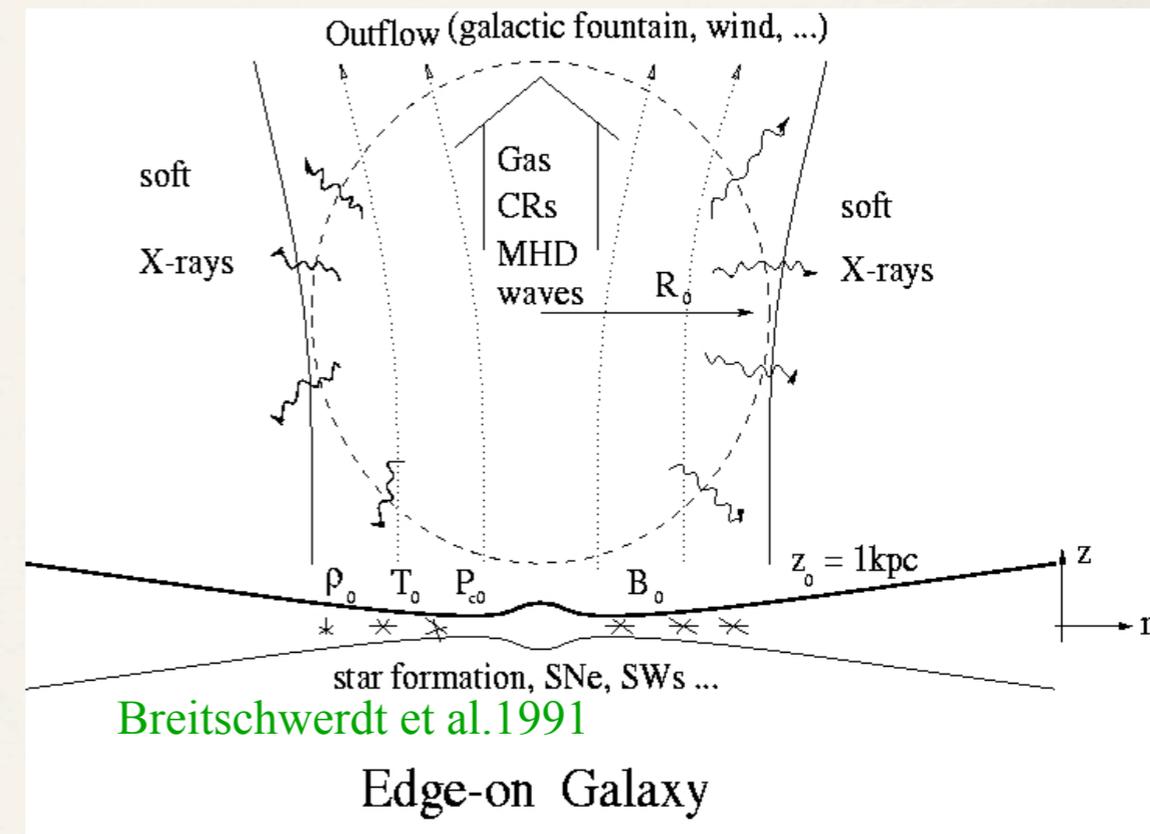


Modeling galactic winds (I)

- * **model edge-on (starburst) galaxies:** e.g. NGC 253, NGC 3079
- * underlying **galactic wind model: steady-state outflow** driven by **thermal gas, CR and wave pressures** (cf. Breitschwerdt et al. 1991)
- * **dynamically and thermally self-consistent** modelling:
 - ➔ Outflow changes ρ and T
 - * this modifies ionization structure
 - * which in turn modifies cooling function $\Lambda(\rho, T)$
 - ➔ which changes outflow
- * flow is described in a flux tube given by:

$$A(z) = A_0 \left[1 + \left(\frac{z}{H} \right)^2 \right]$$

Mass-loaded wind flow!!!



Top: steady-state galactic wind model, in which gas, CRs and waves drive an outflow with a smooth subsonic-supersonic transition if there is strong coupling between CRs and gas

Self-consistent Modeling of Outflows & Ionization Structure

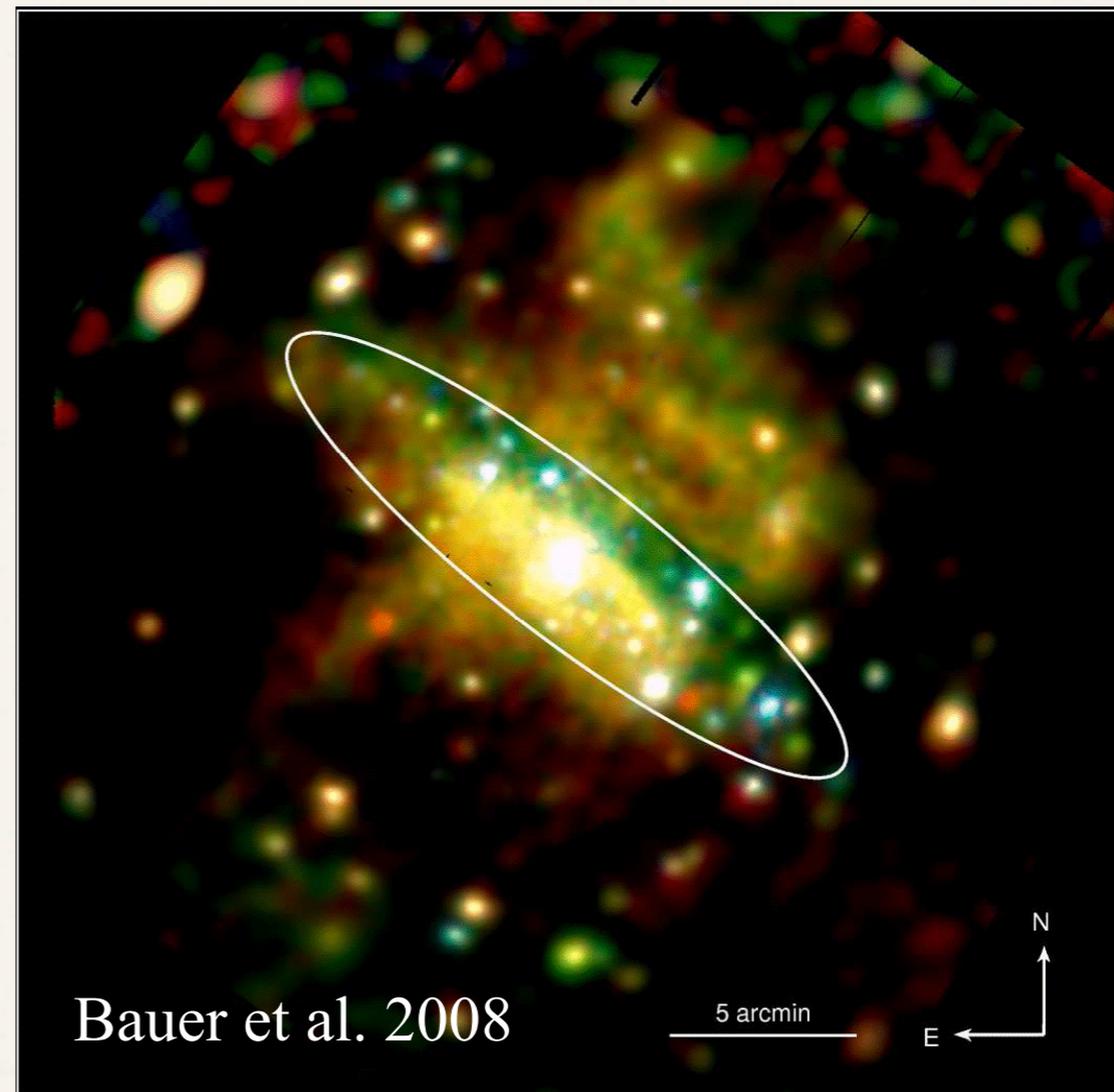
Procedure:

- ❖ **Generating** an outflow model and follow time-dependent evolution of ions (NEI)
 - ❖ **Binning** of high-resolution unabsorbed synthetic (model) spectrum into e.g. EPIC pn channels (for XMM-Newton)
 - ❖ **Folding** spectrum through detector response matrix
- ➔ **Treating observed and synthetic spectrum equally!**
- ❖ **Fitting** synthetic spectrum in XSPEC (X-ray spectral fitting routine) to observational data
 - ❖ **Comparing** with observed spectrum and iterate outflow model if necessary until convergence

Modeling galactic halos with outflow (I): NCG 253



- 2MASS mosaic of NGC253
- Shows also extranuclear SB

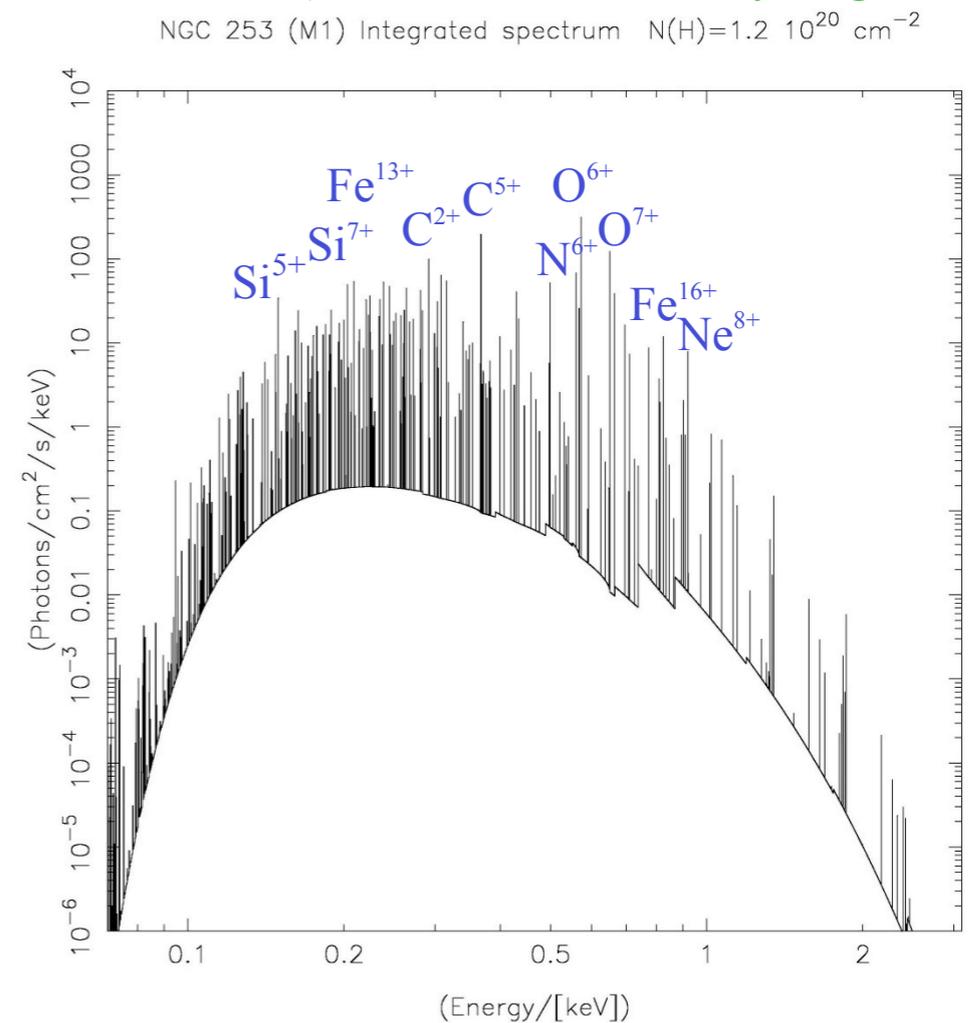


- XMM EPIC pn: Soft X-ray halo of NGC253 (0.2 – 0.5 keV)

Modeling galactic halos with outflow (IV): NCG 253

- ❖ **NEI spectrum** mimics a “multi-temperature” halo by its characteristic lines, but is physically radically different from it
- ❖ **Reason:** Sum of CIE spectra cannot represent the specific **thermodynamic path** of a true NEI spectrum

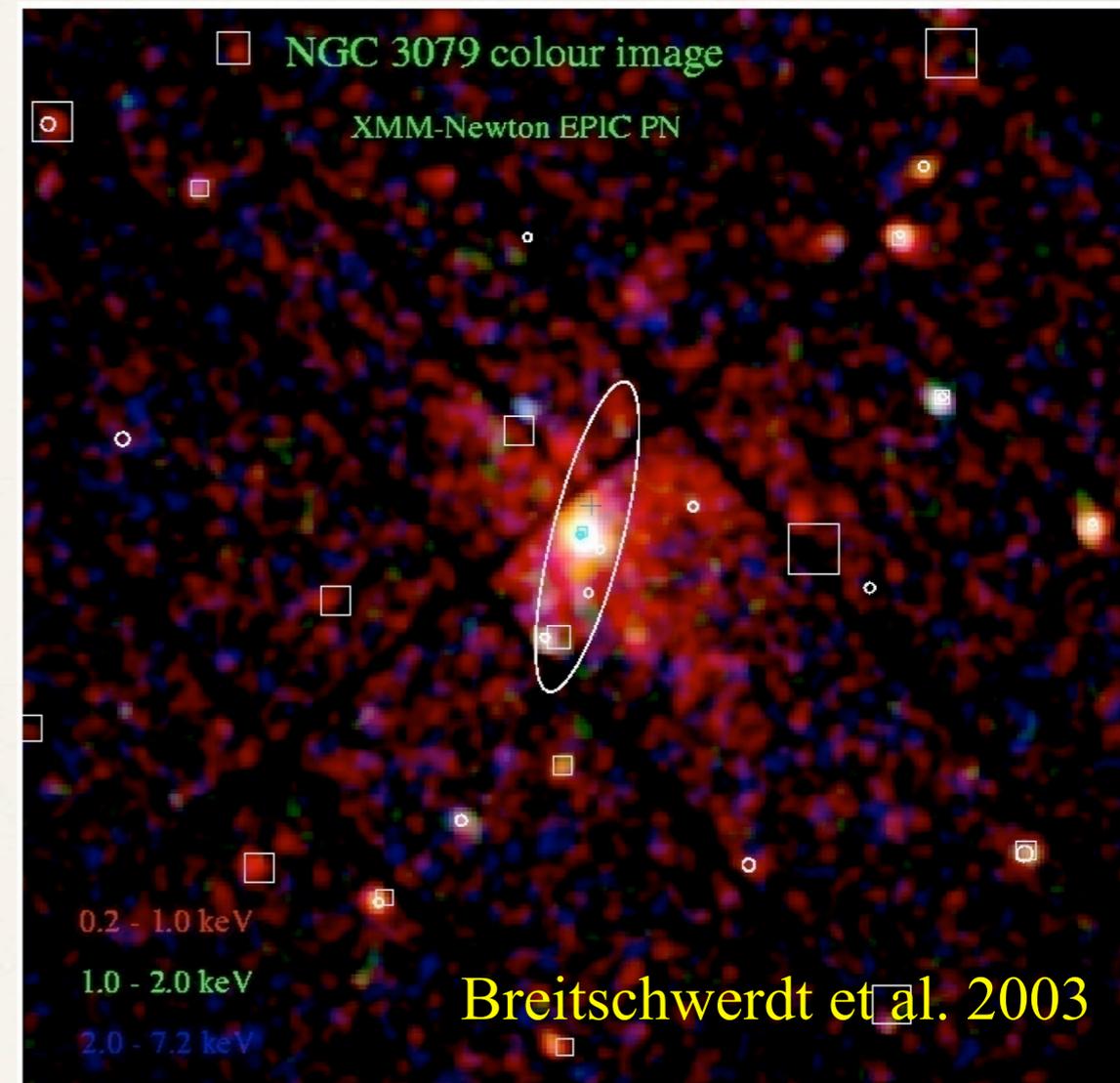
NEI Model (Breitschwerdt & Freyberg 2003)



Top: Integrated spectrum of a dynamically and thermally self-consistent NEI simulation; the spectrum is a composite of continuum and lines, which are characteristic for the plasma history → spectrum will be folded through detector response

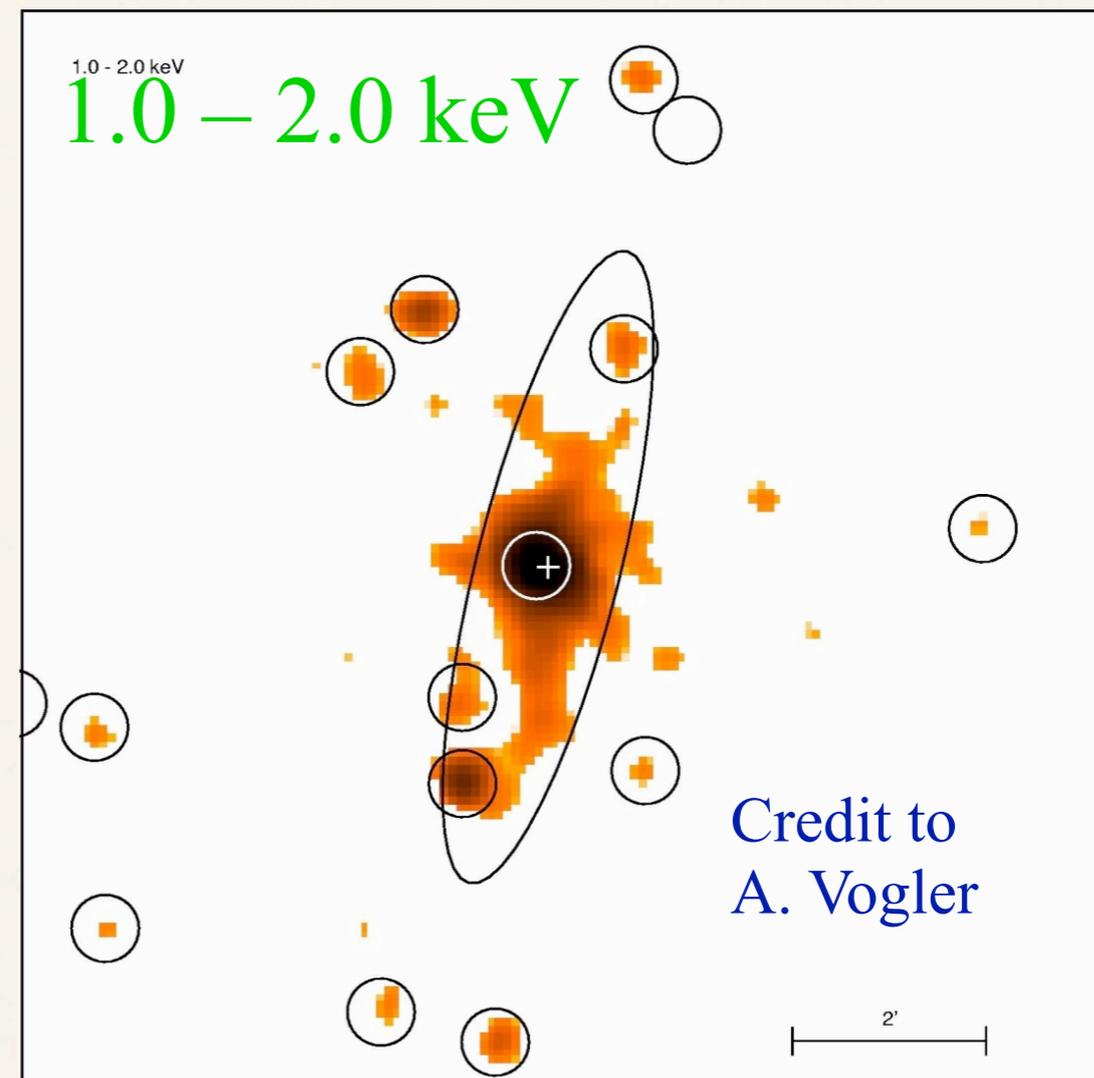
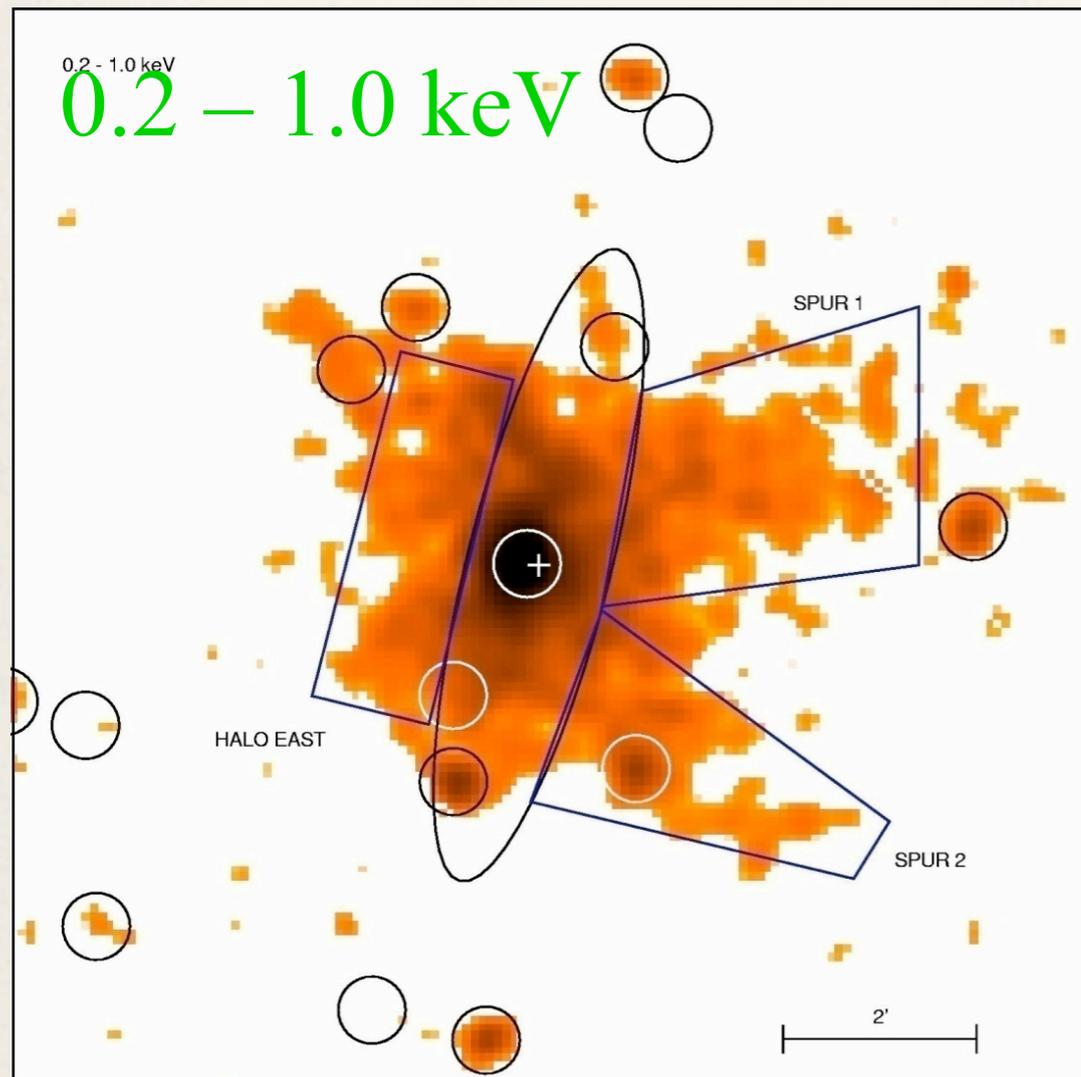
Modeling galactic halos with outflow (V): NCG 3079

- ❖ **NGC 3079**: starburst LINER
SBc galaxy
- ❖ Distance ~ 17 Mpc, Inclination
 $\sim 85^\circ$
- ❖ **Low foreground absorption**
 $\text{Log}(N(\text{H})) = 19.9 \rightarrow$ important
for recording soft X-ray
photons since photoelectric
absorption $\sim E^{-3}$



Top: NGC 3079, XMM-Newton image (EPIC pn camera); the optical disk is indicated by the D_{25} ellipse; the exposure was 25 ksec.

Modeling galactic halos with outflow (VI): NCG 3079



- large extended soft halo emission
- $0.2 \leq E \leq 1.0$ keV

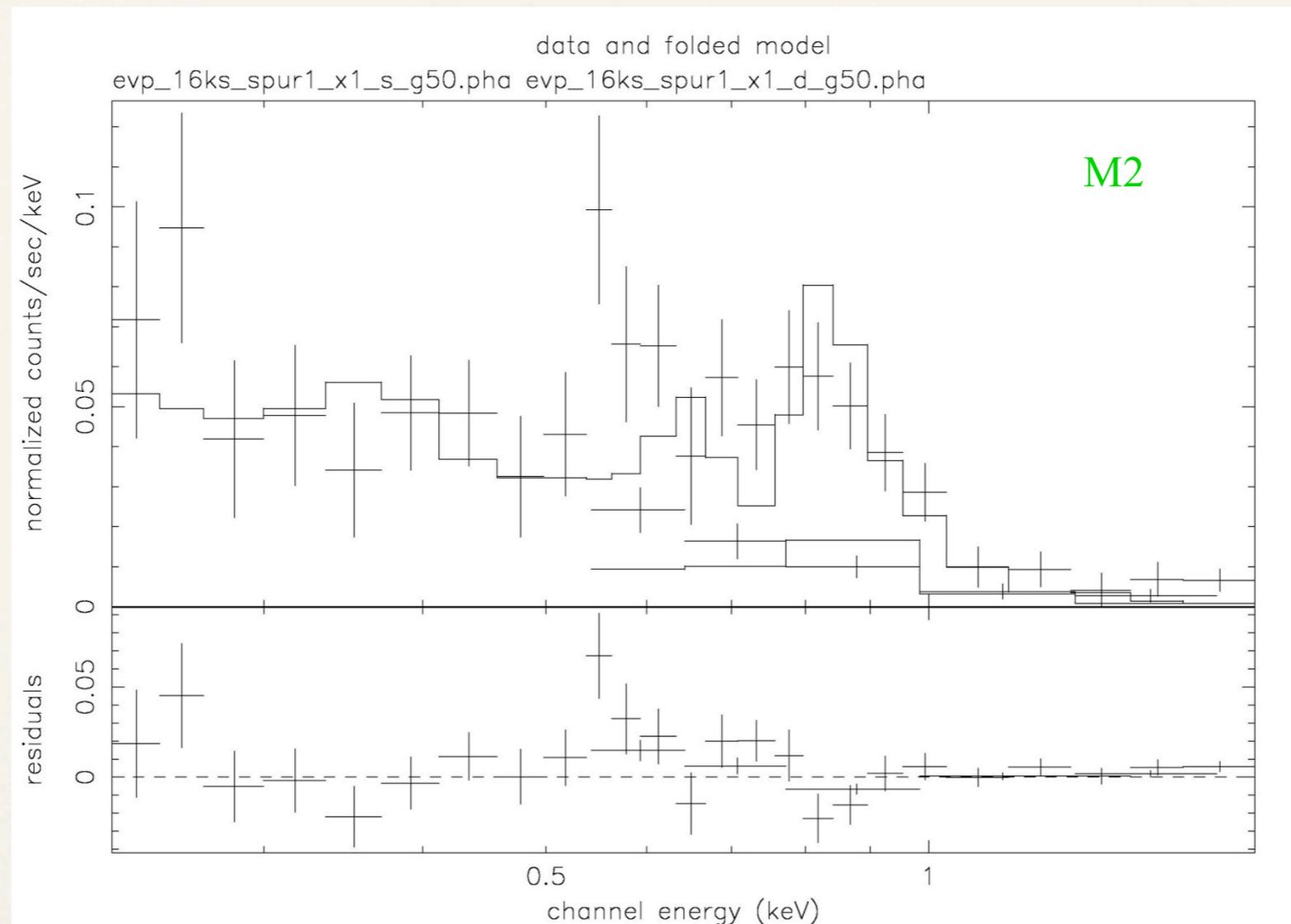
- morphology: soft X-ray spurs
- hard emission largely confined to disk

Modeling galactic halos with outflow (VII): NCG 3079

Model	n_0 [cm ⁻³]	T_0 [10 ⁶ K]	B_0 [μG]	u_0 [km/s]
M1	$3 \cdot 10^{-3}$	3.0	2.0	200.3
M2	$3 \cdot 10^{-3}$	5.0	5.0	312.9
M3	$3 \cdot 10^{-3}$	4.0	5.0	260.4
M4	$4.2 \cdot 10^{-3}$	3.7	5.0	234.1
M5	$5 \cdot 10^{-3}$	3.6	5.0	220.0

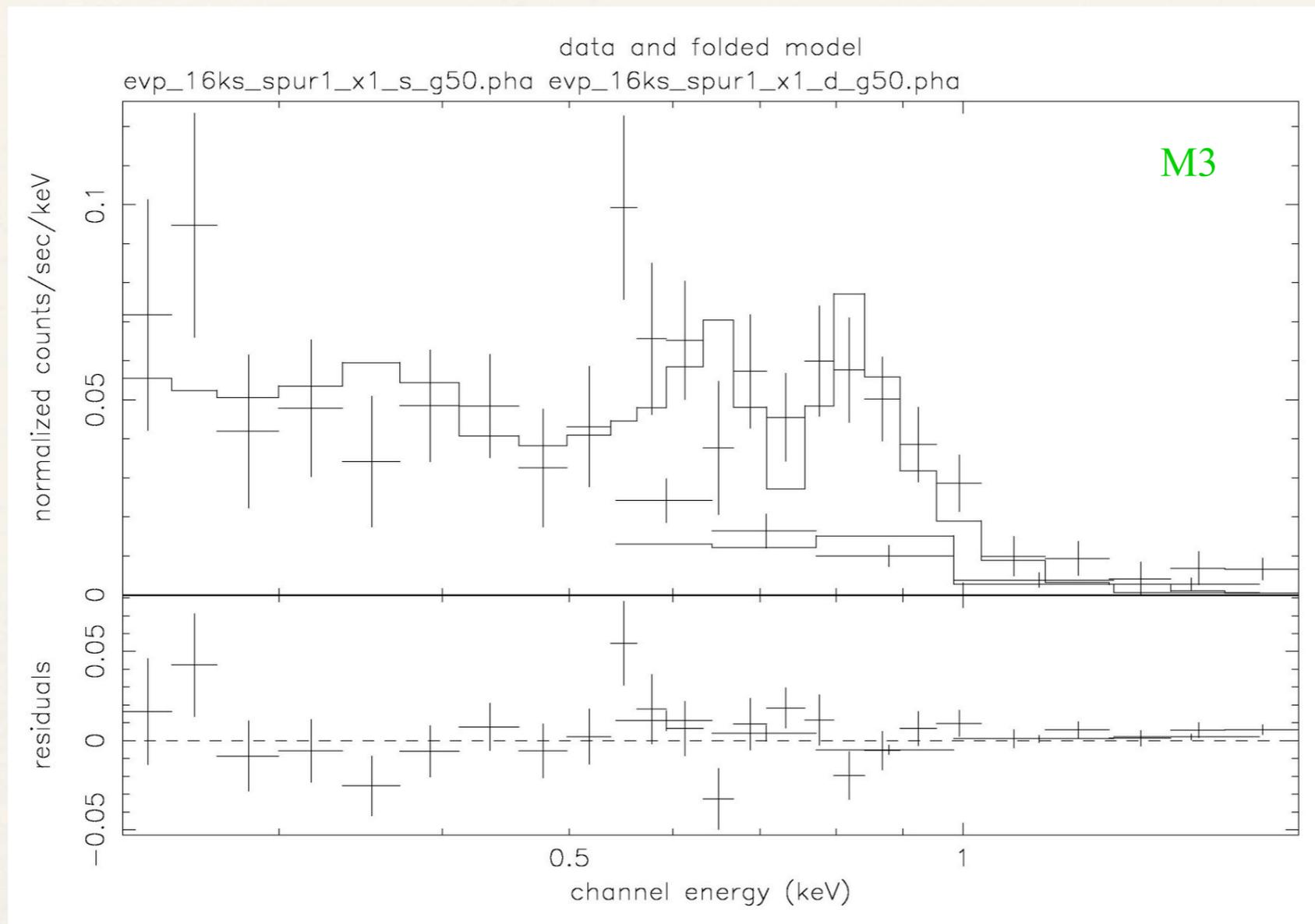
- changing inner boundary conditions, where wind is emanating

Modeling galactic halos with outflow (VII): NCG 3079

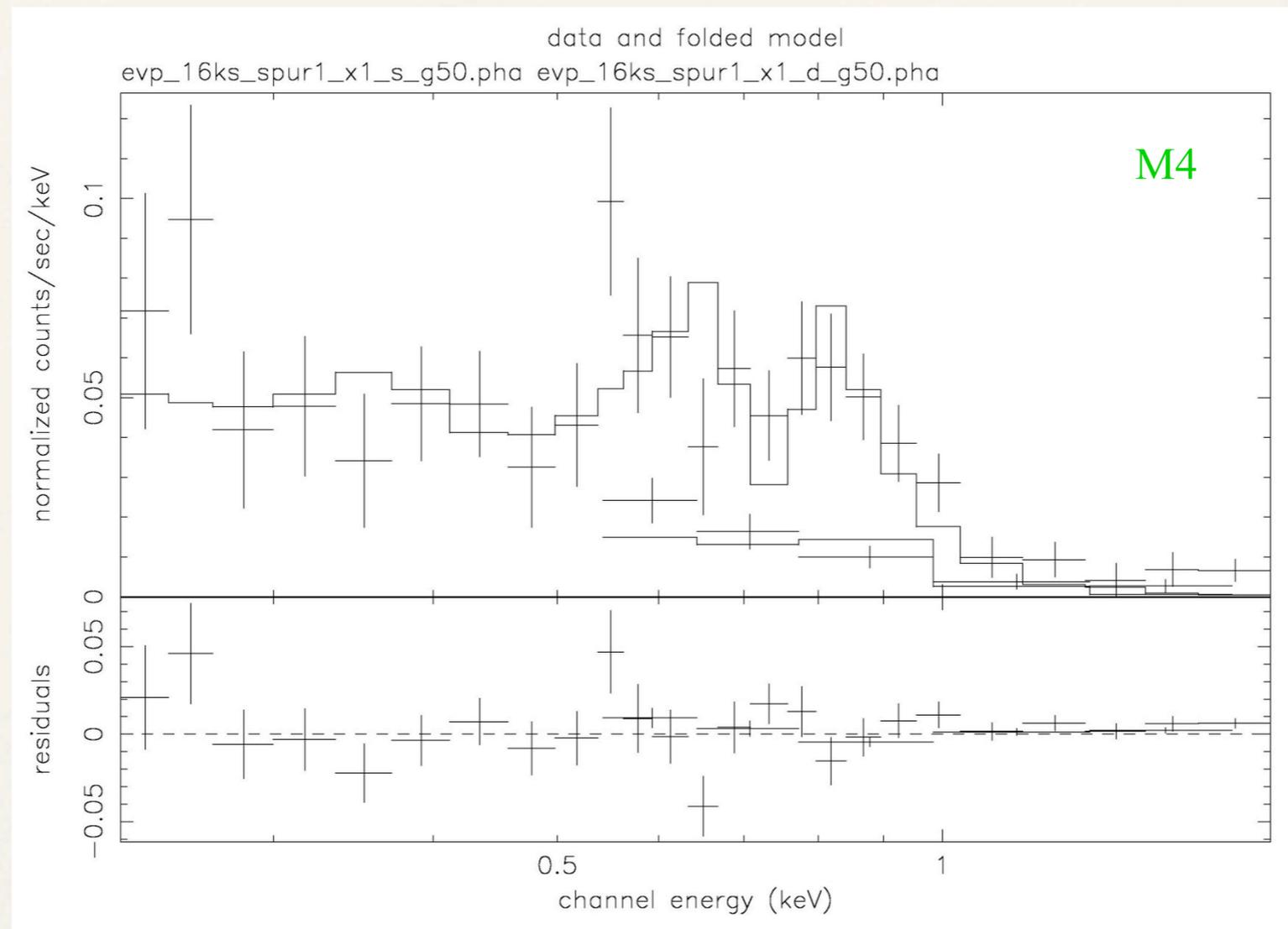


- bad fit in the 0.5 - 0.8 keV region
- too much emission in the 0.8 - 1 keV region → T too high

Modeling galactic halos with outflow (VII): NCG 3079



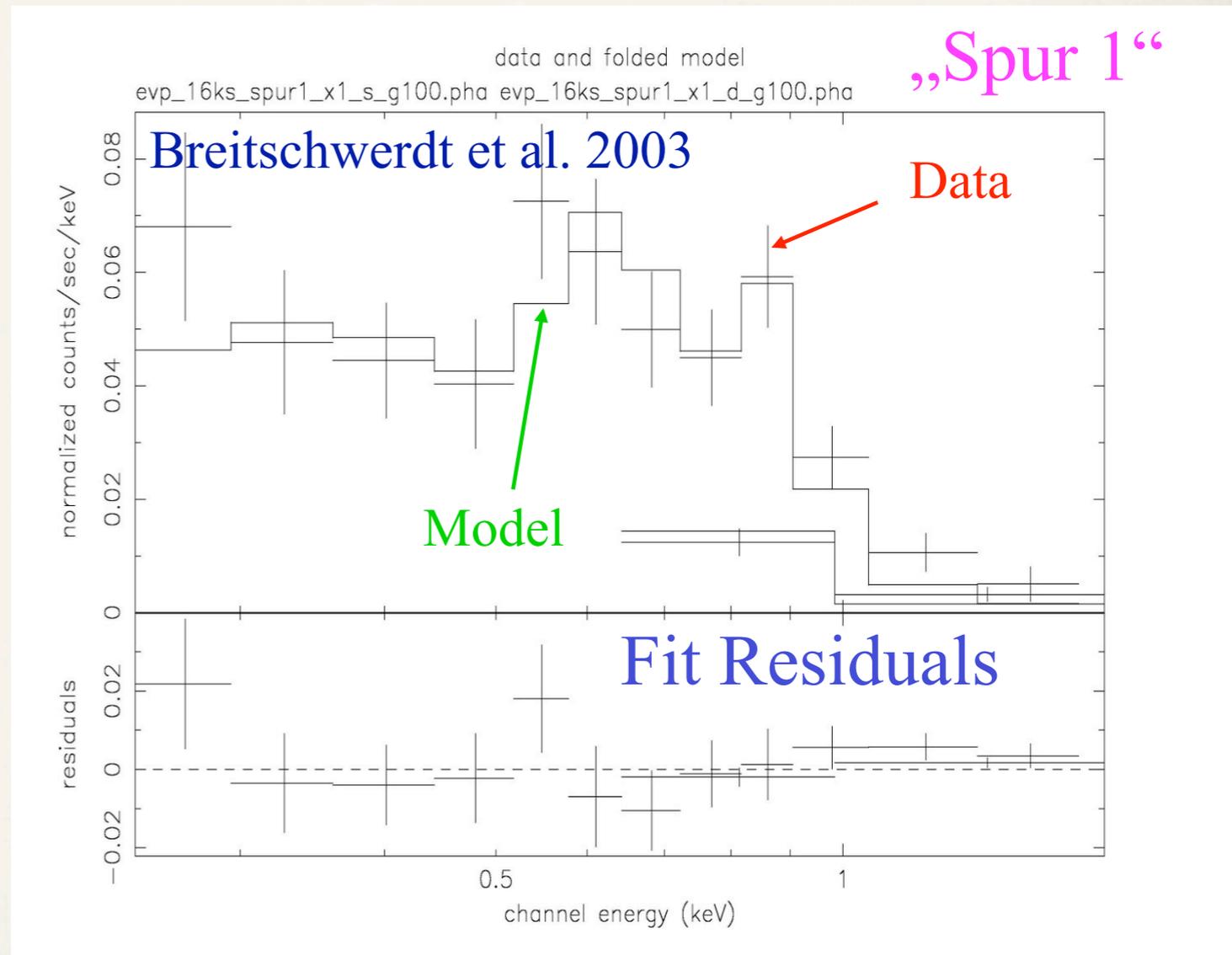
Modeling galactic halos with outflow (VII): NCG 3079



- better fit: T still slightly too high

Modeling galactic halos with outflow (VII): NCG 3079

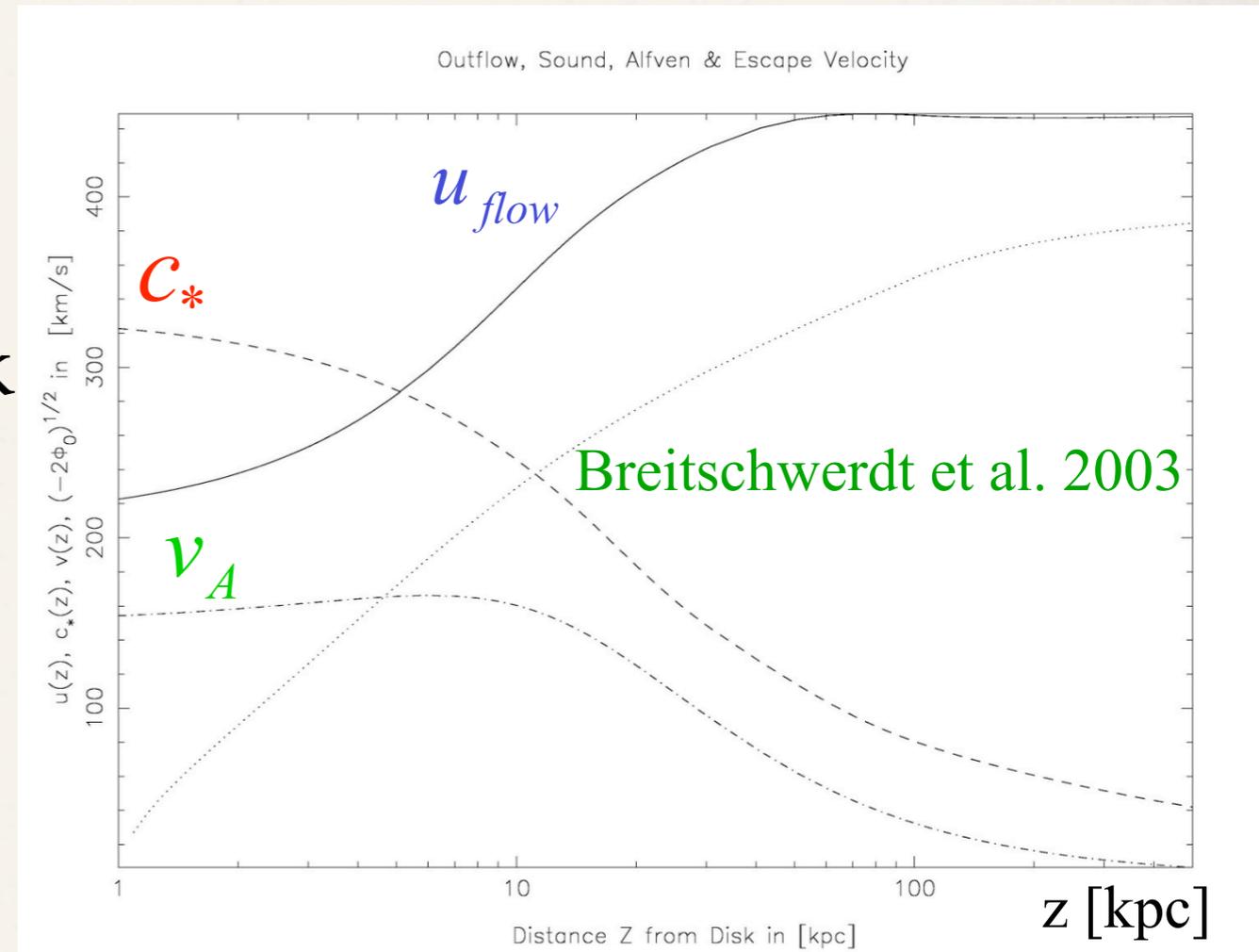
- ❖ **best fit model:**
- ❖ galactic wind with gravitational potential (including dark matter halo)
- ❖ $n_0 = 5 \cdot 10^{-3} \text{ cm}^{-3}$, $T_0 = 3.6 \cdot 10^6 \text{ K}$, $B_0 = 5 \mu\text{G}$, $u_0 = 220 \text{ km/s}$
- ❖ foreground absorption: $N(\text{H}) = 3.9 \cdot 10^{20} \text{ cm}^{-2}$
- ❖ goodness of fit: $\chi^2_{\text{red}} = 1.2$
- ❖ derived mass loss rate: $dm/dt = \rho_0 u_0 = 0.055 M_{\text{sol}}/\text{yr}/\text{kpc}^2$
- ❖ **mass loss rate in “spur” region** ($R = 8\text{kpc}$): $dM/dt = \pi/2 R^2 \rho_0 u_0 = 3.5 M_{\text{sol}}/\text{yr}$



Top: Comparison between data and dynamically and thermally self-consistent galactic wind model → 5 iterations were necessary to achieve an acceptable fit

Modeling galactic halos with outflow (V): NGC 3079

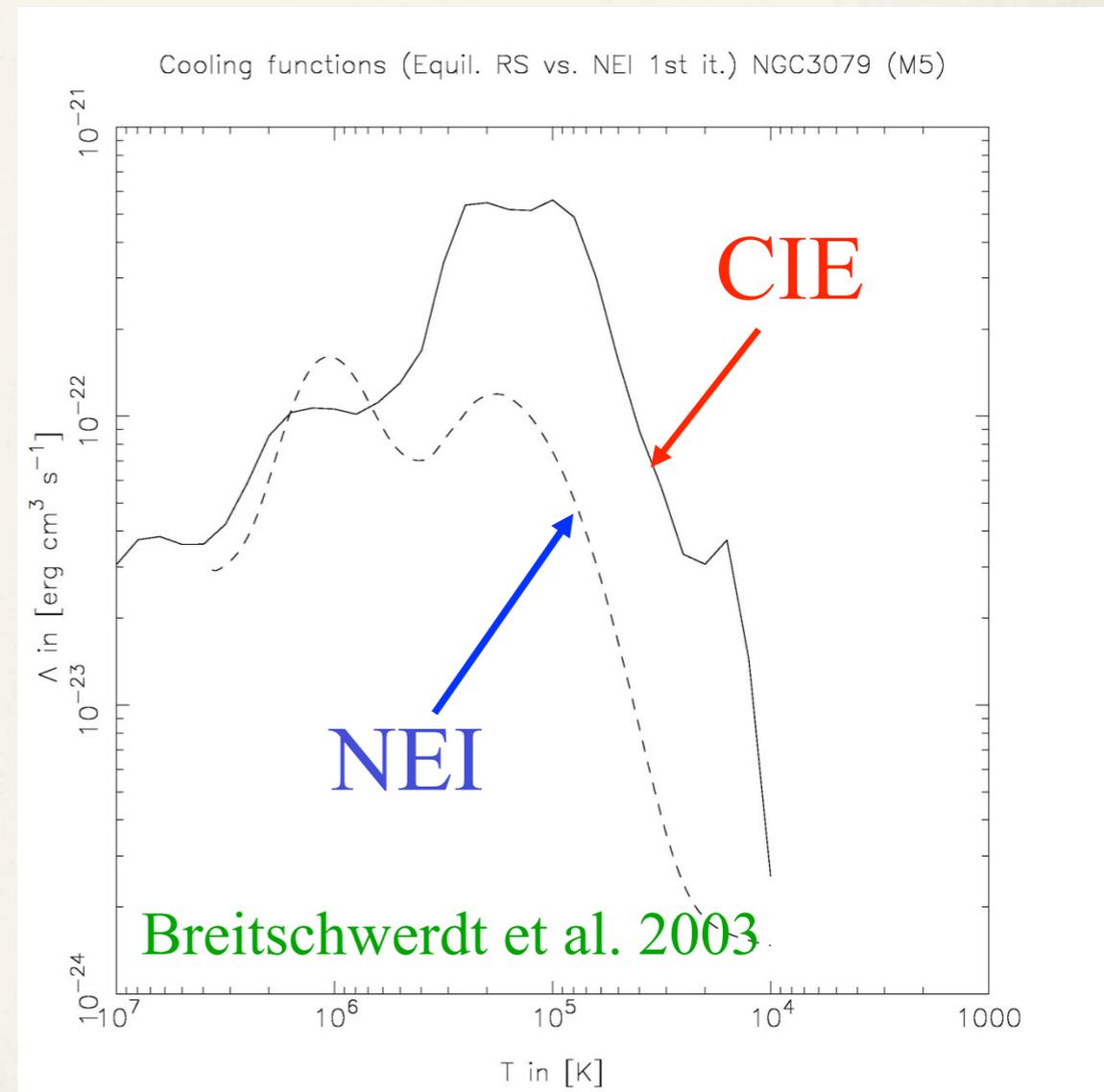
- ❖ **NGC 3079: smooth subsonic-supersonic transition** \rightarrow critical point ($M=1$) in the flow at $z \sim 5$ kpc from the disk
- ❖ superbubble gas injected at the inner boundary ($z_0 = 1$ kpc) with initial velocity $u_0 = 220$ km/s (subsonic, but super-alfvenic)
- ❖ terminal velocity ~ 450 km/s



Top: Derived outflow characteristics for the best fit model

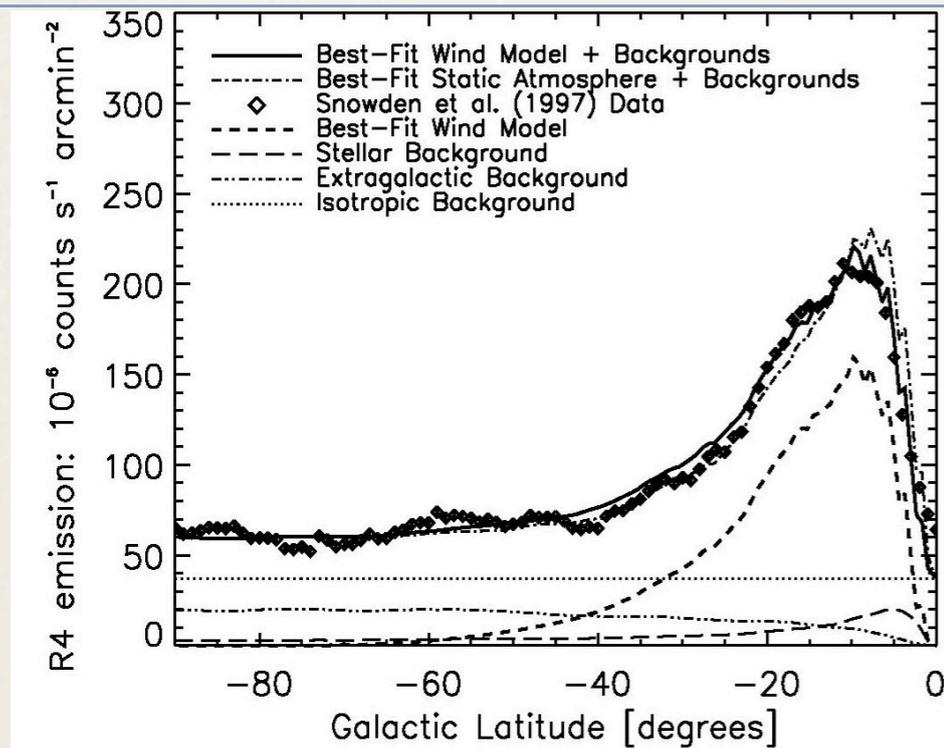
Modeling galactic halos with outflow (V): NCG 3079

- ❖ **Cooling Function:** Cooling curve depends on the ionization state of the plasma
- ❖ in case of a fast adiabatically expanding flow the difference between CIE and NEI cooling curves is striking
- ❖ whereas the CIE cooling curve peaks at $\sim 10^5$ K, the NEI curve in this particular model has a maximum $\sim 10^6$ K, where OVII, OVIII lines are abundant due to **delayed recombination**

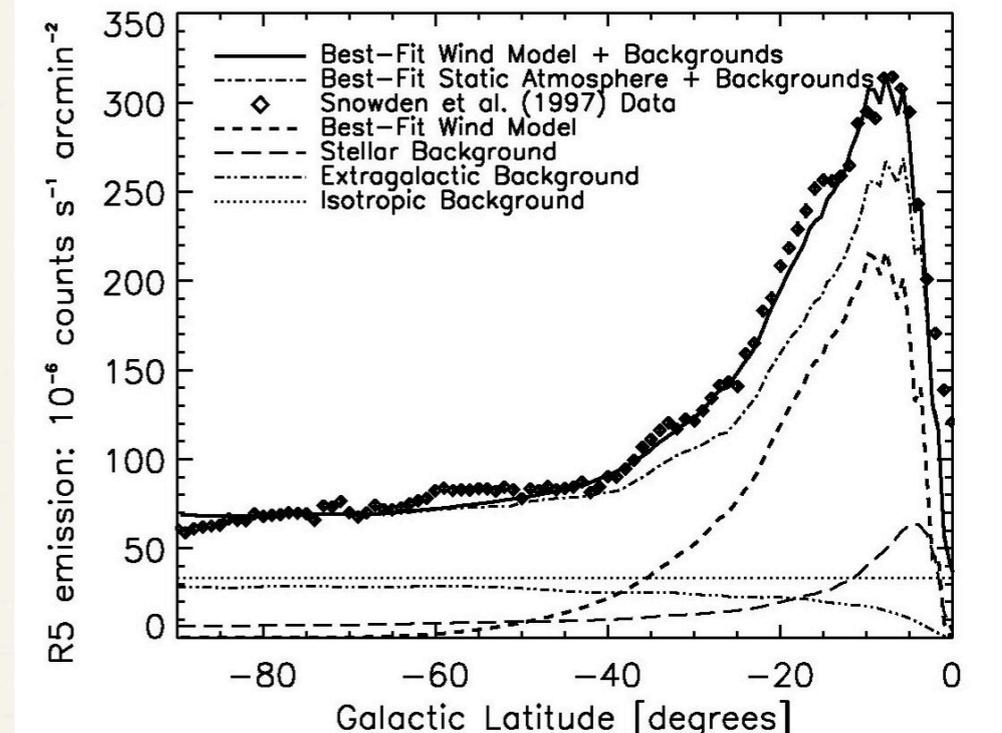


Top: Comparison between cooling curve for CIE and for a dynamical NEI model for the starburst galaxy NGC 3079

Cosmic Ray driven Wind in the Milky Way?



Top: ROSAT PSPC observations (Snowden et al. 1997); shown is ROSAT R4 band (0.64 keV)

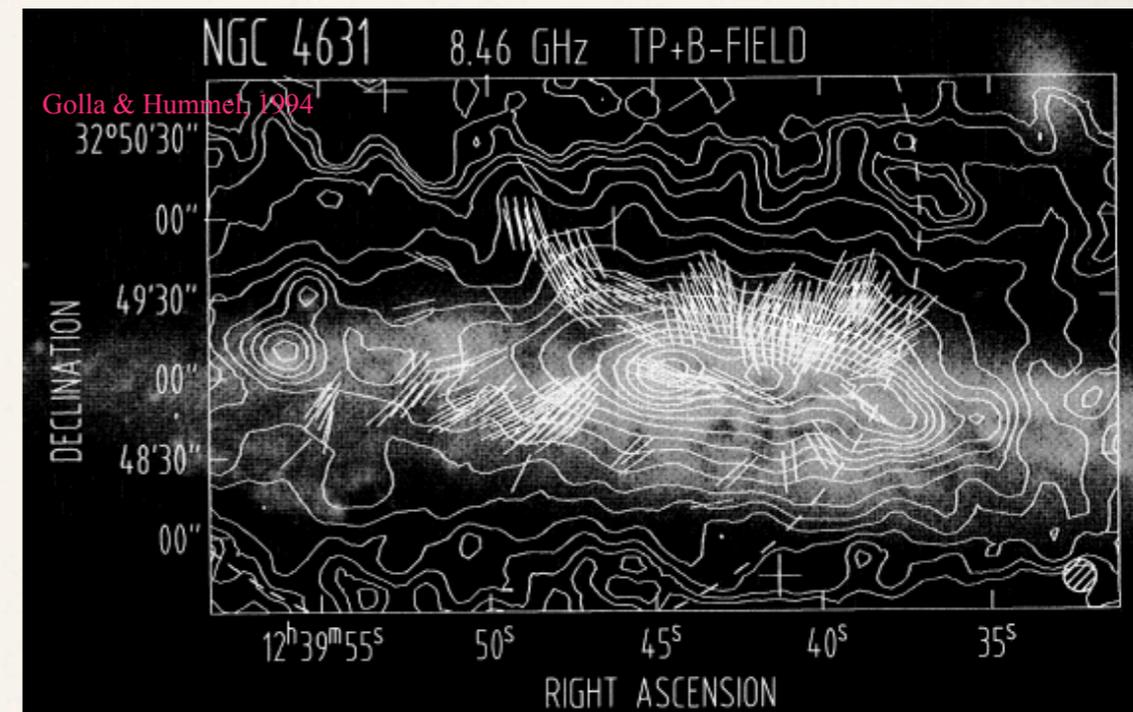


Top: ROSAT PSPC observations (Snowden et al. 1997); shown is ROSAT R5 band (0.85 keV)

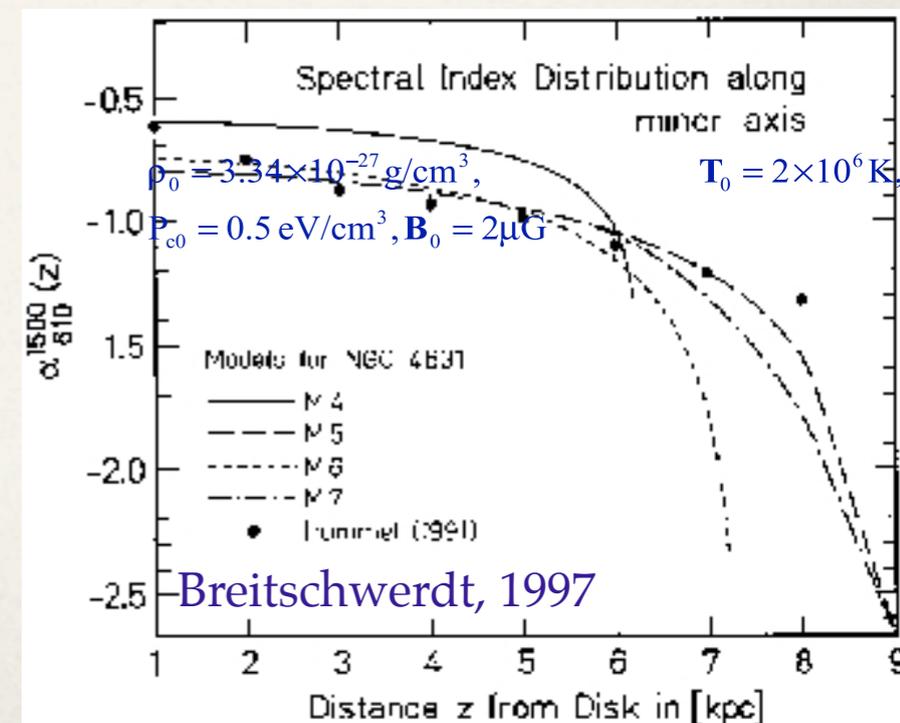
- ★ To fit ROSAT PSPC data in a given region of the MW sky, Everett et al. (2008) show that a **CR driven wind** gives a statistically much better fit than a hydrostatic halo (especially for ROSAT R5 band)
- ★ for lower halo here CIE is a good approximation, since deviation from NEI still small
- ★ fiducial wind model for Milky Way: $n_0 = 6.9 \cdot 10^{-3} \text{ cm}^{-3}$, $T_0 = 2.9 \cdot 10^6 \text{ K}$, $u_0 = 173 \text{ km/s}$

CR electron transport in galactic halos (I): NCG 4631

- ❖ **Non-thermal radio emission of NGC 4631:** significant linear polarization for $z \leq 5$ kpc
- ❖ noticeable B-field component perpendicular to galactic disk
- ❖ **Modelling:** solve **diffusion-advection** transport equation for electrons incl. **synchrotron** and **inverse Compton** losses
- ❖ **radio spectral index** variation is a measure of dominant transport process: flat curve is indicative of accelerating **advection** flow, compensating for increasing losses with time \rightarrow galactic wind



Top: radio map superimposed with polarization vectors on optical image
Right: self-consistent spectral index variation in galactic wind model



CR electron transport in galactic halos (II): NGC 253

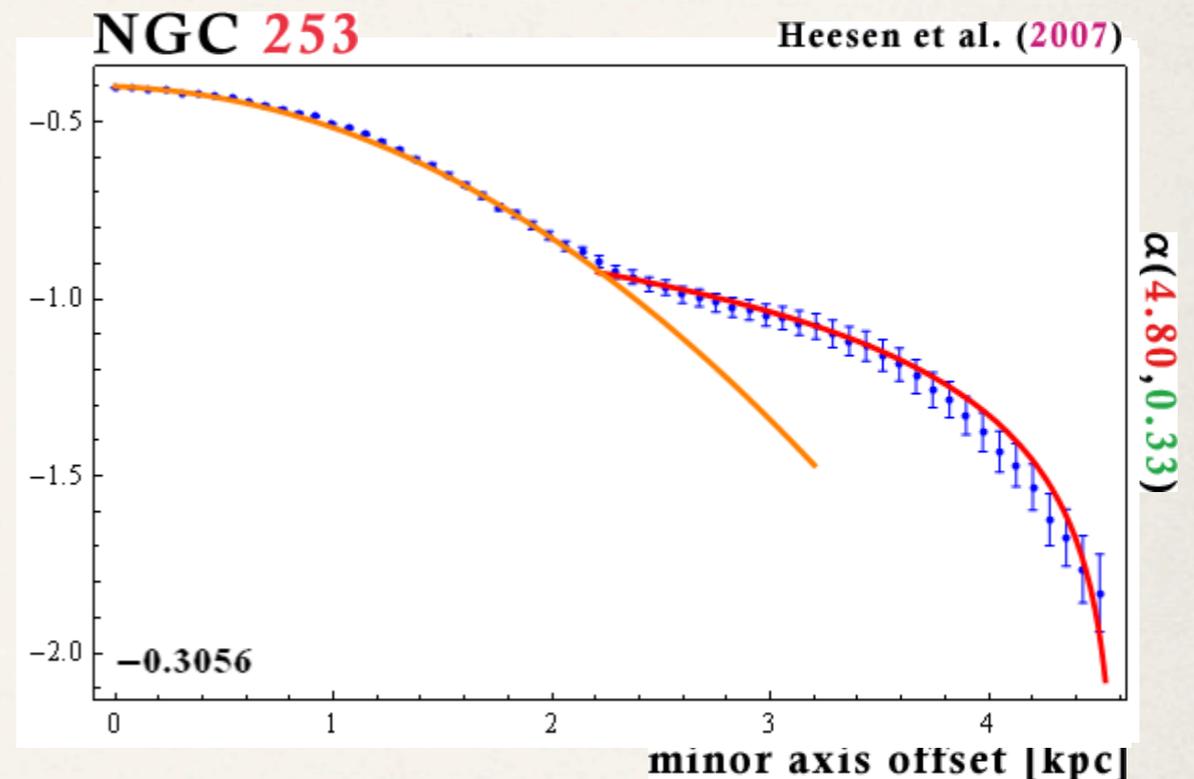
- ❖ Non-thermal radio emission of NGC 253:
- ❖ CR electron transport equation

$$\begin{aligned}
 & - \frac{\partial}{\partial z} \left(D(E, z) \frac{\partial N(E, z)}{\partial z} - u(z) N(E, z) \right) \\
 & - \frac{\partial}{\partial E} \left(\frac{1}{3} \frac{du}{dz} E N(E, z) - \frac{dE}{dt} N(E, z) \right) = Q(E, z)
 \end{aligned}$$

$$Q(E, z) = K_0 E^{-\gamma_0} h_g \delta(z)$$

- ❖ spectral index close to sources up to vertical distances from disk of $z \sim 1-2$ kpc dominated by **diffusion**
- ❖ for $z \geq 1-2$ kpc transport dominated by **advection**
- ❖ transport mechanism varies locally in agreement with local superbubble break-out from galactic disk

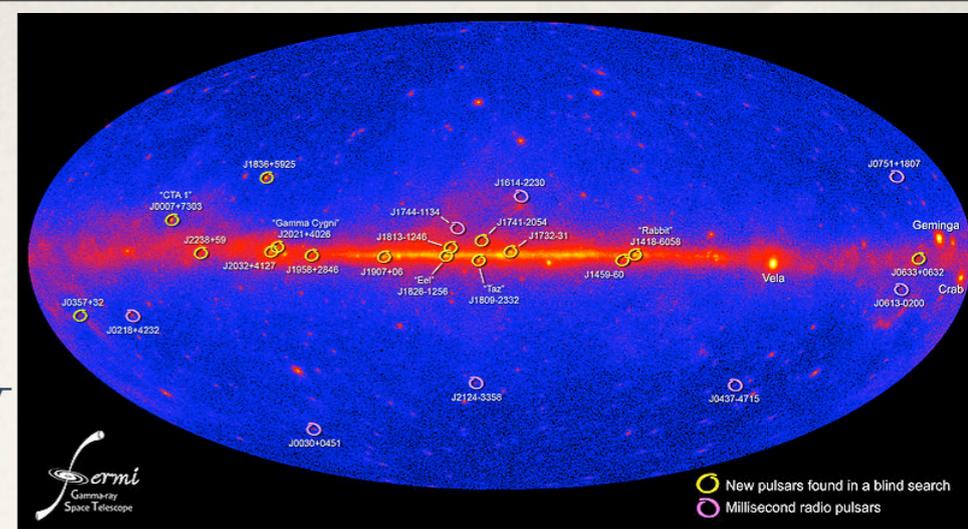
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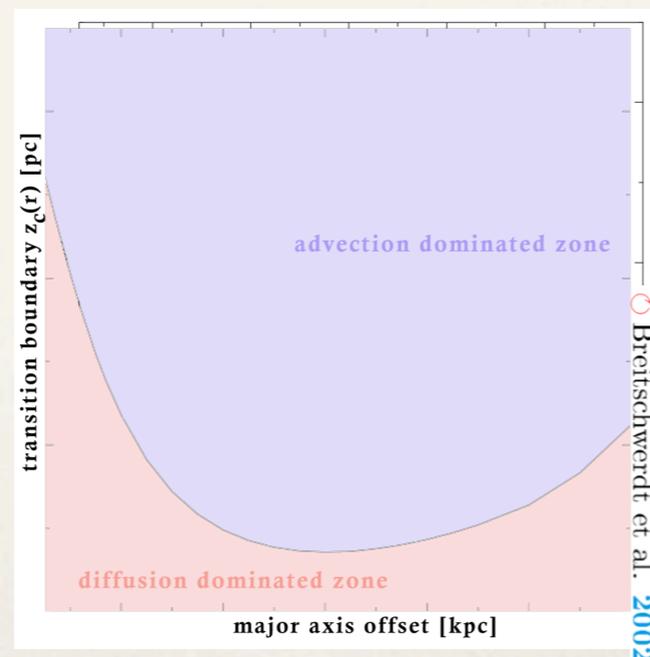
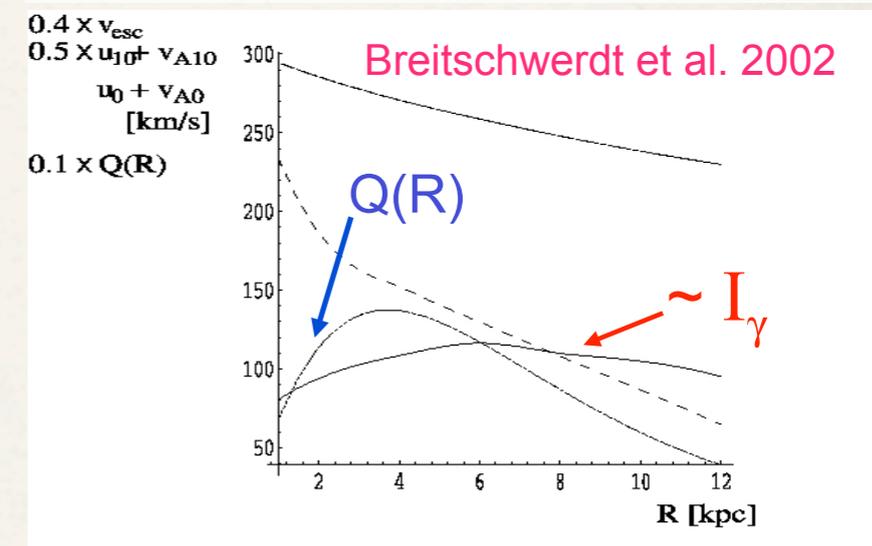
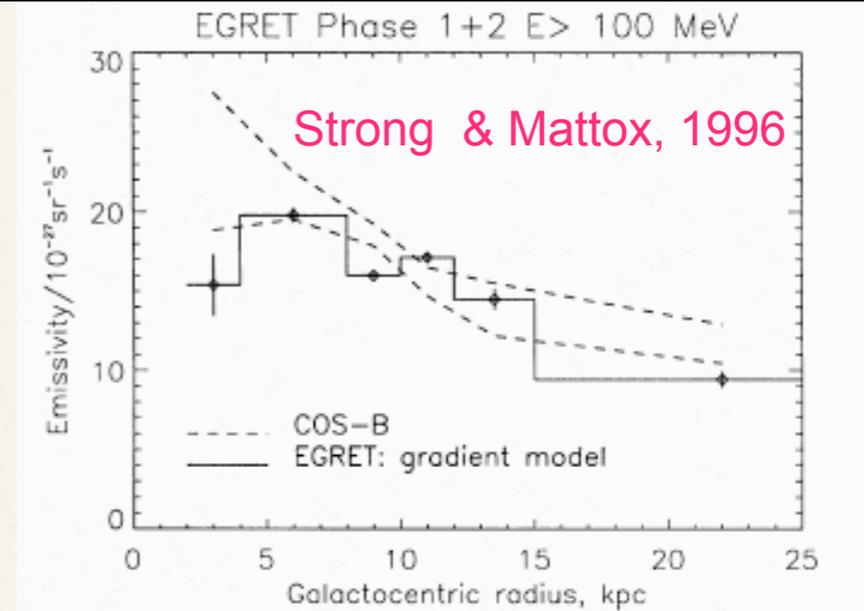
Advection **Diffusion** **Diffusion-Advection**

Top: Comparison between the model (including a galactic wind) and observations (blue dots with error bars) of the starburst galaxy NGC 253; data from Heesen et al.

Radial Diffuse γ -ray Gradient in the Galaxy



- Diffuse γ -ray gradient in the Milky Way: Cos-B, EGRET measured **shallow** gradient of diffuse γ -emission
- for $E > 100$ MeV γ -rays are mainly due to π_0 -decay, which are due to interaction between CR protons and HI atoms \rightarrow γ -emission should follow **CR source** distribution (SNRs and pulsars), which peaks at $R \sim 4$ kpc
- diffusion model can only marginally reproduce γ -ray gradient for huge CR halo
- simple model**: radially dependent **diffusion-advection** boundary due to local radial variations in **star formation**

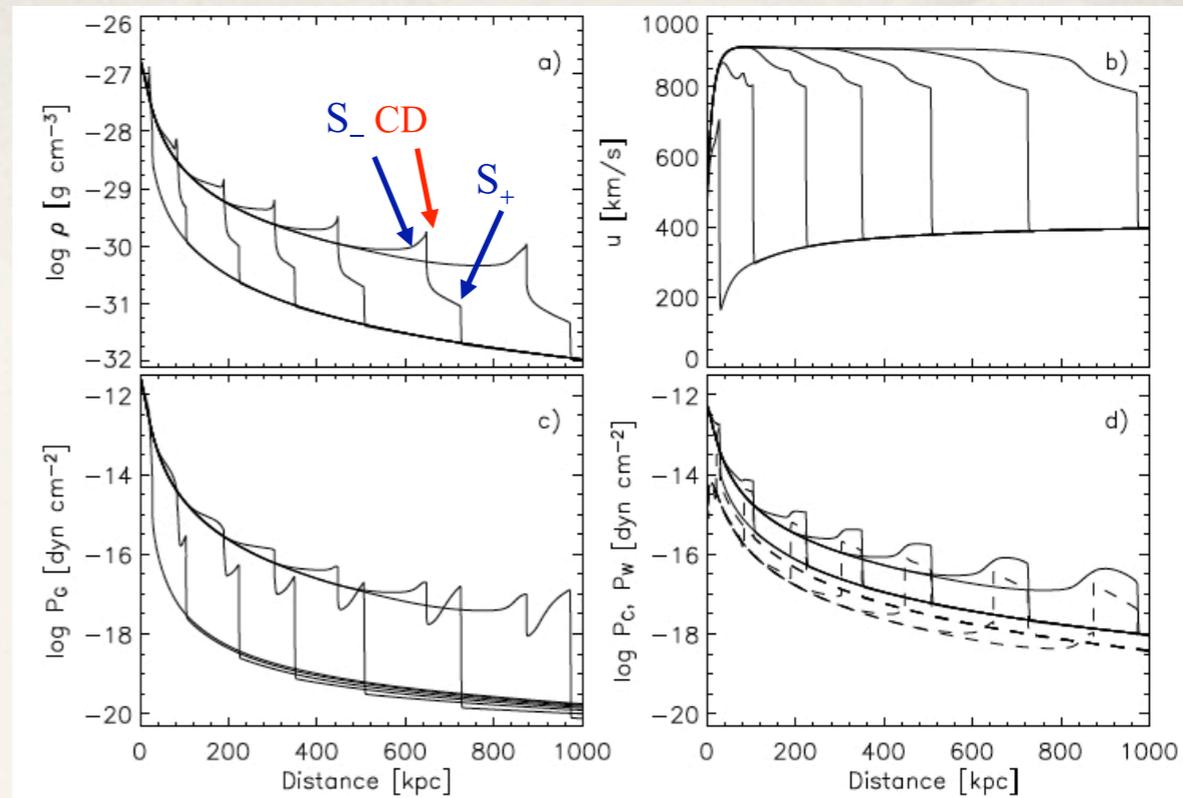


Top: Fermi γ -ray all-sky survey
Middle: diffuse γ -ray gradient according to COS-B and EGRET observations
Bottom: model (galactic wind) of diffuse γ -ray emission

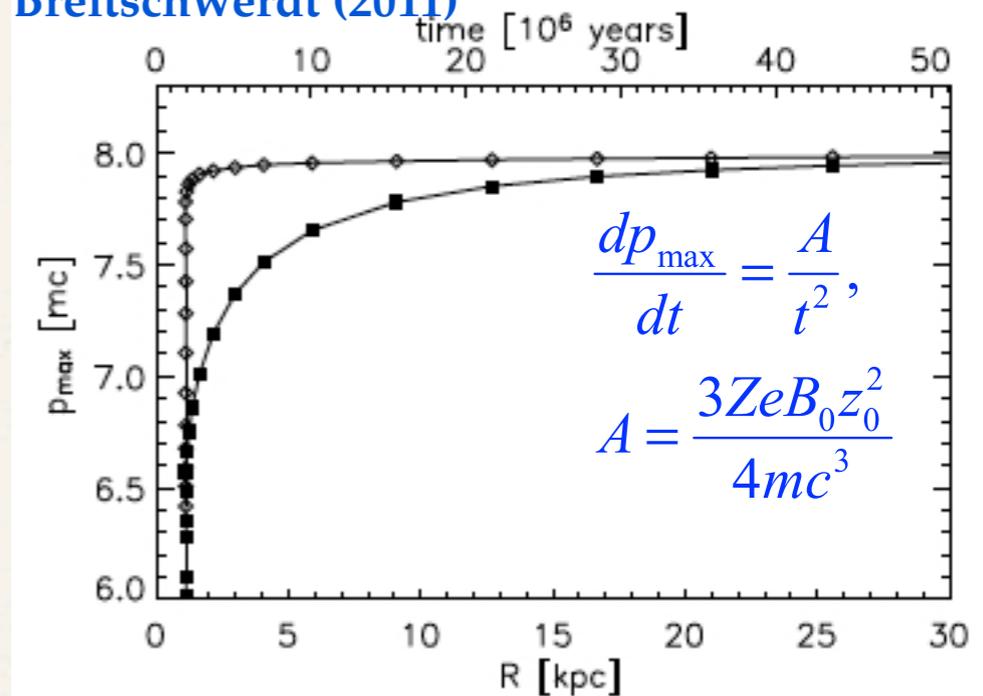
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CR acceleration beyond the “Knee” (I)

- ❖ **Time-dependent galactic wind** calculations (Dorfi & Breitschwerdt 2010) confirm **stationary** wind solutions as time-asymptotic flow
- ❖ for starburst galaxies we use **time-dependent boundary conditions**, reflecting the duration of a starburst → increase of CR & gas pressures by a factor of 10 at $z=z_0$
- ❖ double shock structure in the galactic wind region → **post-acceleration** of galactic CRs (**1st order Fermi**)
- ❖ particles are convected downstream of **forward shock**, i.e. towards the galactic disk
- ❖ particle acceleration modifies shock → **subshock** → shock strengthens as propagating down a density gradient
- ❖ within a few $10^6 - 10^7$ yr, particles can reach energies up to $10^{17} - 10^{18}$ eV



Dorfi & Breitschwerdt (2011)

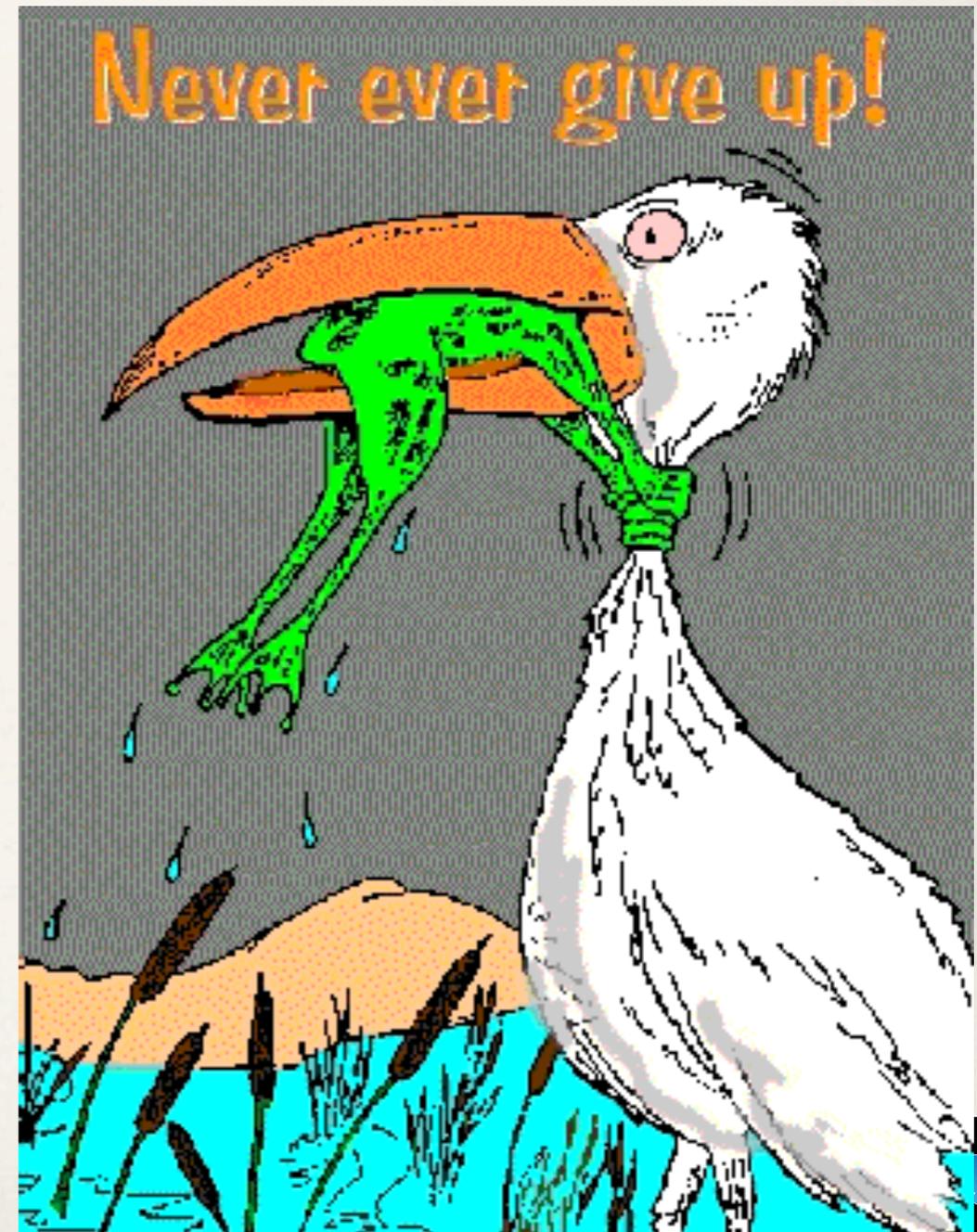


Top: density (a), velocity (b), gas (c), CR and wave pressures (d) for shocks in a galactic wind of the Milky Way for $3 \cdot 10^7 \leq t \leq 10^9$ yr
Bottom: maximum momentum of particles post-accelerated in galactic wind for forward (filled squares) and forward + reverse shock (diamonds)

Summary & Conclusions

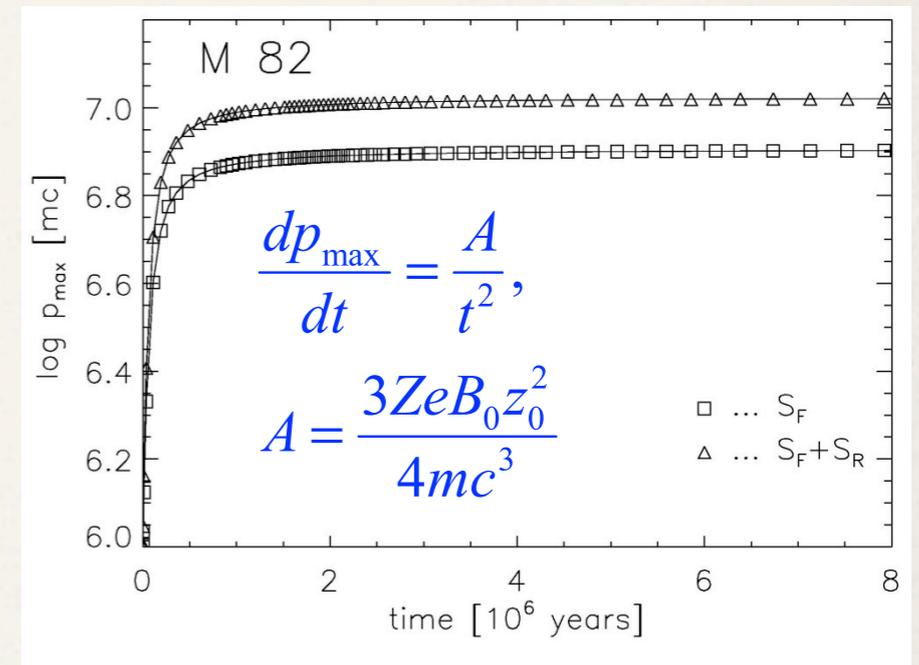
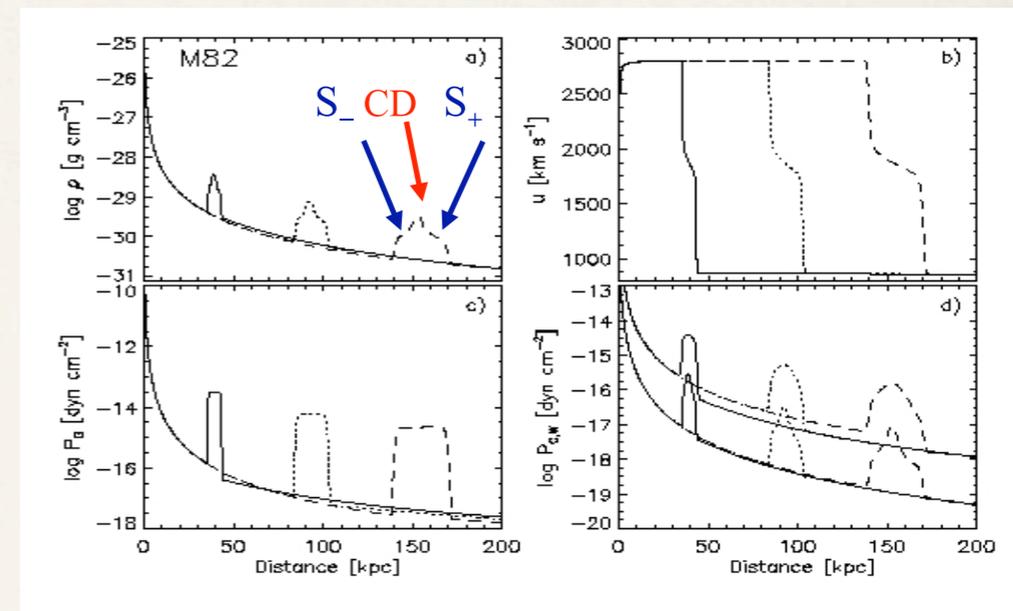
- * Cosmic Rays (CRs) are essential component of ISM
- * CR energy density comparable to thermal plasma, magnetic field, turbulent motions
- * CR composition similar to ISM but with light element enhancement (spallation)
- * CR differential energy spectrum consists of power laws
- * 10% of total SN energy is turned into CRs
- * Diffusive shock acceleration as most probable mechanism
- * CRs leave Galaxy → replenishment in $\sim 2 \cdot 10^7$ yr
- * CRs couple to plasma via self-excited waves → CR transfer momentum to gas
- * CR driven outflows → galactic winds
- * enhanced star formation and **superbubbles** drive thermal **galactic winds**
- * winds can explain **X-ray and radio halos** → **thermally & dynamically self-consistent models** → flattening of **radio spectral index** due to **advection flow**
- * **CR acceleration** beyond the “knee” → wind shocks can explain energies & spectrum!

Thank you for your attention!



CR acceleration beyond the “Knee” (I)

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