

Cosmology and black holes: (towards) the first structures

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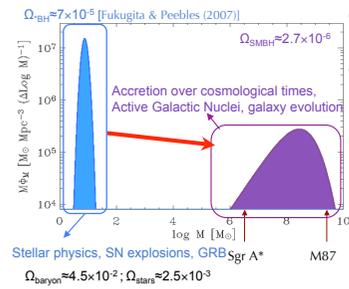
Excellence Cluster "Universe"
Max-Planck Institut für Extraterrestrische Physik
Garching, Germany



SKADS MCCT, 14/4/2008



Black Holes in the local Universe



- Density of stellar mass BHs increases w/galaxy mass
- How does the high-mass (Supermassive) peak grow?

This talk:

- What do we know about evolution of SMBH population (mass function, accretion rates, etc.)
 - A lot! (up to $z \sim 4-5$)
 - Evolution of mass and accretion rate density
 - Constraints on radiative efficiency and avg. BH spin
 - Anti-hierarchical evolution (downsizing)
 - Radiative vs. kinetic energy output
- Frontiers: The first Black holes
 - Eddington or Super-Eddington?
 - The role of mergers
 - The first seeds

Accretion efficiency, Eddington limits

- In order to get close (R), a particle of mass m must get rid of energy
 $E_{ib} = GM_{BH}m/R$
- Efficiency of accretion in liberating rest mass energy:
 $\eta = E_{ib}/mc^2 = R_g/2R_{in}$, with $R_g = GM_{BH}/c^2$ GR $\rightarrow 0.06 < \eta(a) < 0.42$
- BHs grow by accreting mass: $\text{Power}_{\text{released}} = [\eta/(1-\eta)] (dM_{BH}/dt)c^2$
- Self-regulating luminosity

$$L_{\text{Edd}} = 4\pi GMc m_p / \sigma$$

1. $L_{\text{Edd,es}} = 1.3 \times 10^{38} (M_{BH}/M_{\text{sun}}) [\text{XRB, AGN}]$
2. $L_{\text{Edd,\gamma}} = 8 \times 10^{33} (E_{\gamma}/50 \text{ MeV})^{-2} (M_{BH}/M_{\text{sun}}) [\text{GRB}]$

AGN and Cosmology: (very) early developments

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LETTERS TO THE EDITOR

ASTRONOMY

Evidence on the Evolutionary Character of the Universe derived from Recent Red-shift Measurements

Hoyle and Burbidge have recently examined the red-shifts of a number of quasi-stellar radio sources^{1,2} and have plotted their radio flux densities (2) against red shift (z). They conclude that the results are inconsistent with a cosmological interpretation of these red-shifts.

Their conclusions show that this correlation is not valid and that indeed the observed relationship between δ and z provides valuable new information on the evolutionary nature of the Universe.

If, on the other hand, the red-shifts are attributed to motion, the only way of modifying the flux density F_{ν} is by the factor $(1+z)^2$. The flux density F_{ν} is then given by $F_{\nu} = F_{\nu}^0 (1+z)^2$, where F_{ν}^0 is the flux density at the source. The flux density F_{ν} is then given by $F_{\nu} = F_{\nu}^0 (1+z)^2$, where F_{ν}^0 is the flux density at the source.

Longair 1966

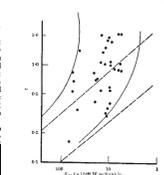
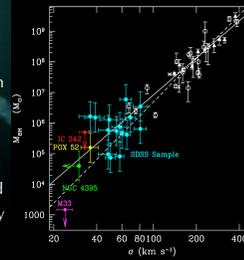


Fig. 1. The red-shift flux density plot against red shift for the sources listed in Table 1 of Burbidge and Hoyle (1966). The flux density F_{ν} is in units of $10^{-26} \text{ W m}^{-2} \text{ Hz}^{-1}$. The red shift z is in units of $1+z^2$.

AGN and Cosmology: a shift of paradigm

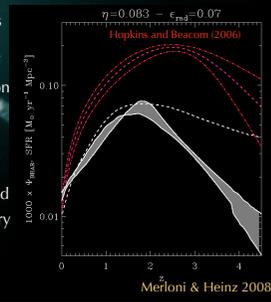
- BH/Galaxy Scaling relations discovered in 2000 (HST)
- Nowadays QSOs/AGN can:
 - Regulate galaxy formation
 - Stop cooling flows
 - “produce” early type galaxies with the right colors
 - Keep the universe ionized
- QSOs/AGN tracks the history of star formation in the Universe

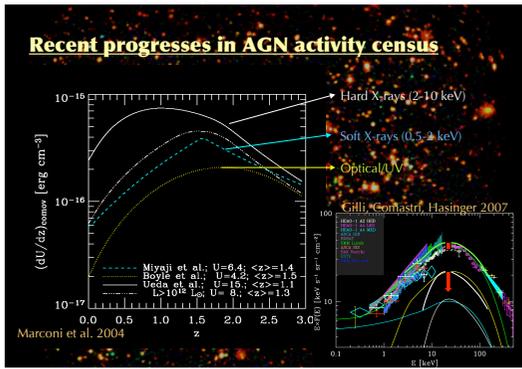


Gebhardt et al. 2000; Ferrarese et al. 2000; Greene et al. 2005

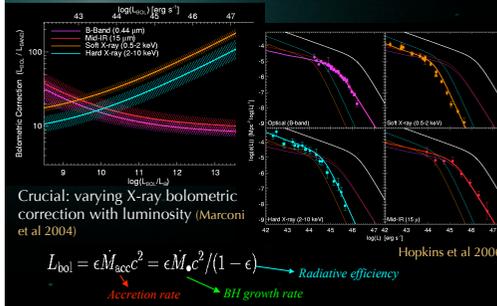
AGN and Cosmology: a shift of paradigm

- BH/Galaxy Scaling relations discovered in 2000 (HST)
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Bolometric corrections robustness



Note: radiative efficiency vs. accretion efficiency

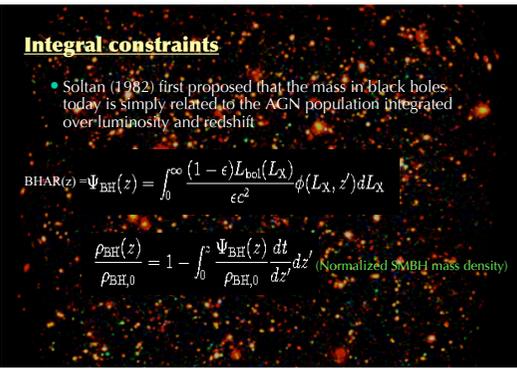
radiative efficiency accretion efficiency (BH spin)

$$\epsilon \equiv \epsilon(a, \dot{m}, \dot{m}_{\text{cr}}) = \eta(a) f(\dot{m}, \dot{m}_{\text{cr}})$$

Non Spinning BH $0.06 \leq \eta(a) \leq 0.42$ Maximally Spinning BH

$$f(\dot{m}, \dot{m}_{\text{cr}}) = \begin{cases} 1, & \dot{m} \geq \dot{m}_{\text{cr}} \\ \dot{m}/\dot{m}_{\text{cr}}, & \dot{m} < \dot{m}_{\text{cr}} \end{cases}$$

Determined by the complex physics of gas accretion



Integral constraints

- Soltan (1982) first proposed that the mass in black holes today is simply related to the AGN population integrated over luminosity and redshift

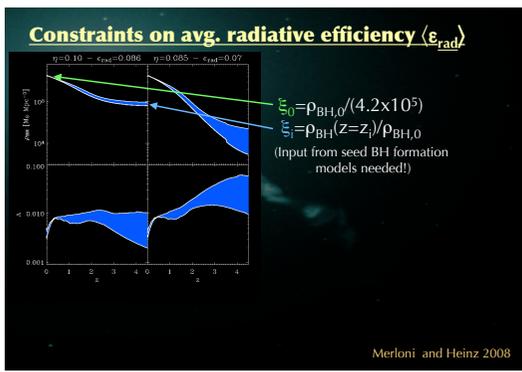
$$\text{BHAR}(z) = \Psi_{\text{BH}}(z) = \int_0^{\infty} \frac{(1 - \epsilon)L_{\text{bol}}(L_X)}{\epsilon c^2} \phi(L_X, z') dL_X$$
$$\frac{\rho_{\text{BH}}(z)}{\rho_{\text{BH},0}} = 1 - \int_0^z \frac{\Psi_{\text{BH}}(z')}{\rho_{\text{BH},0}} \frac{dt}{dz'} \quad (\text{Normalized SMBH mass density})$$

Convergence in local mass density estimates

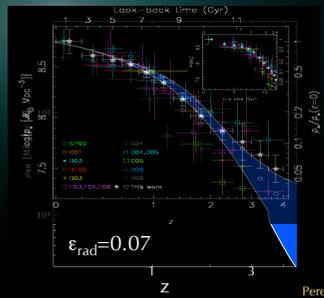
Table 3. Modification of Table 1. Here, the local SMBH mass density estimates have been fully corrected for their dependence on h , and transformed to $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$. While some estimates of ρ_{SMB} appear not to have changed from Table 1, one should note that the quoted dependence on h may have changed. The term $f(h)$ is used to denote that a more complicated dependence on h exists and needs to be taken into account if one is to transform these values to a different Hubble constant (see, e.g. equation 2).

Study	Method	$\rho_{\text{SMB}} (E/S0)$	$\rho_{\text{SMB}} (S0)$	$\rho_{\text{SMB}} (\text{total})$
		$10^5 M_{\odot} \text{ Mpc}^{-3}$	$10^5 M_{\odot} \text{ Mpc}^{-3}$	$10^5 M_{\odot} \text{ Mpc}^{-3}$
Graham et al. (2007)	$M_{\text{BH}} \propto$	$(3.46 \pm 1.16) h_0^2$	$(0.60 \pm 0.46) h_0^2$	$(4.44 \pm 1.67) h_0^2$ ✓
Wyithe (2000)	$M_{\text{BH}} \propto$	$(1.08 \pm 0.38) h_0^2$ ✓
Fiksel & Poehle (2014)	$\rho_{\text{SMB}} \propto$	$(1.4^{+1.1}_{-0.7}) h_0^2$	$(1.3^{+1.1}_{-0.7}) h_0^2$	$(1.1 \pm 0.5) h_0^2$ ✓
Marconi et al. (2004)	$M_{\text{BH}} (z, \sigma)$	$3.20 h_0^{0.8} f(h)$	$1.3 h_0^{0.8} f(h)$	$(4.6^{+1.1}_{-0.8}) h_0^{0.8} f(h)$ ✓
Shankar et al. (2004)	$M_{\text{BH}} L$	$(4.3^{+1.1}_{-0.7}) h_0^{0.8} f(h)$	$(1.8^{+0.7}_{-0.5}) h_0^{0.8} f(h)$	$(6.0^{+1.1}_{-0.8}) h_0^{0.8} f(h)$ ✓
Shankar et al. (2004)	$M_{\text{BH}} \sigma$	$(3.4^{+1.1}_{-0.7}) h_0^2$	$(1.4^{+0.5}_{-0.3}) h_0^2$	$(4.8^{+1.1}_{-0.8}) h_0^2$ ✓
McLure & Dunlop (2004)	$M_{\text{BH}} L$	$(4.8 \pm 0.7) h_0^{0.8} f(h)$
Wyithe & Loeh (2003)	$M_{\text{BH}} \sigma$
Alar & Richstone (2002)	$M_{\text{BH}} \sigma$	$(4.3 \pm 1.0) h_0^{0.8} f(h)$	$(1.4 \pm 1.3) h_0^{0.8} f(h)$	$(2.1^{+2.4}_{-1.3}) h_0^2$ ✓
Yu & Tremaine (2002)	$M_{\text{BH}} \sigma$	$(2.0 \pm 0.2) h_0^2$	$(0.9 \pm 0.2) h_0^2$	$(6.0 \pm 2.0) h_0^{0.8} f(h)$ ✓
Mariti & Ferrarese (2001)	$\rho_{\text{SMB}} \propto$	$(2.9 \pm 0.4) h_0^2$ ✓
Salcedo et al. (1999)	$\rho_{\text{SMB}} \propto$	$6.2 h_0^2$	$2.0 h_0^2$	$4.0 h_0^2$ ✓
				$8.2 h_0^2$ ✓

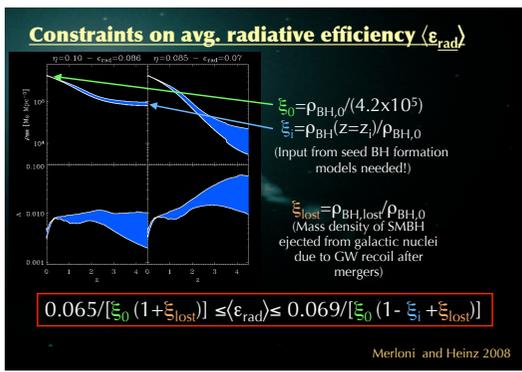
Graham and Driver (2007)



SMBH vs TOTAL stellar mass densities



Perez-Gonzalez et al.(2007)

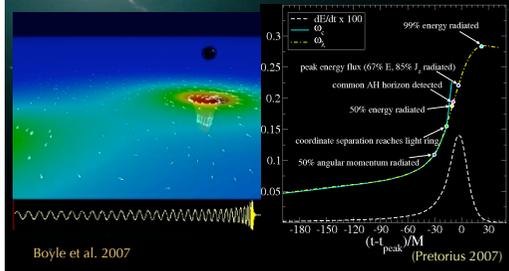


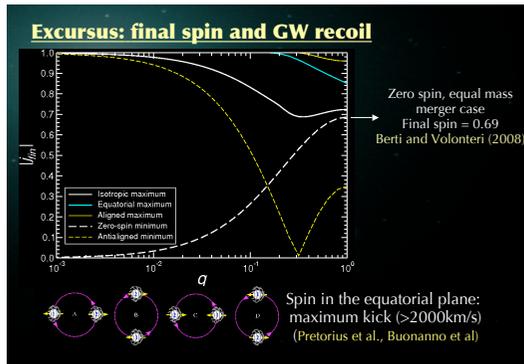
Excursus: SMBH mergers



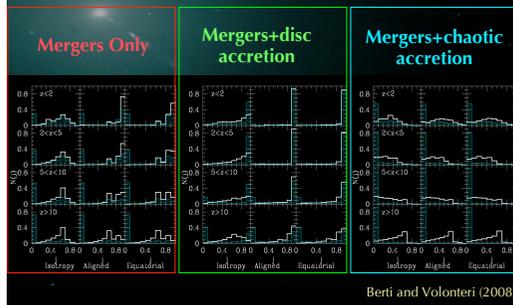
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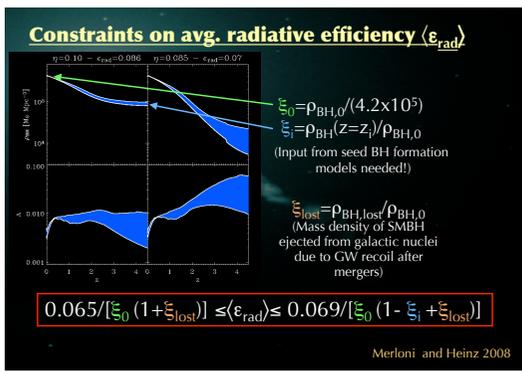
- Recently, great progress in the general relativistic simulations of coalescing Black Holes (Pretorius 2007)



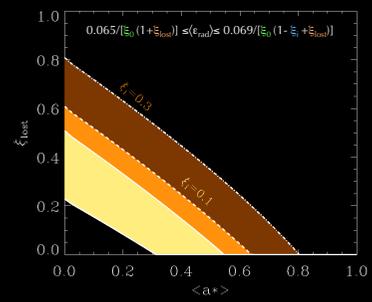


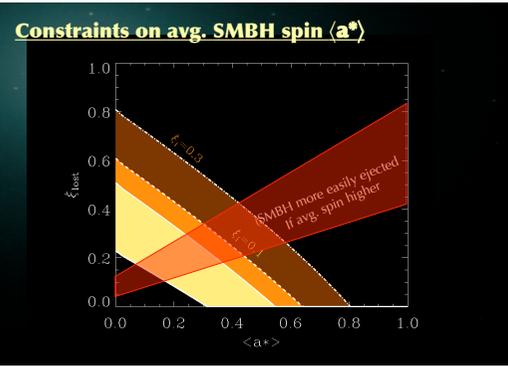
SMBH spin evolution: accretion vs. mergers

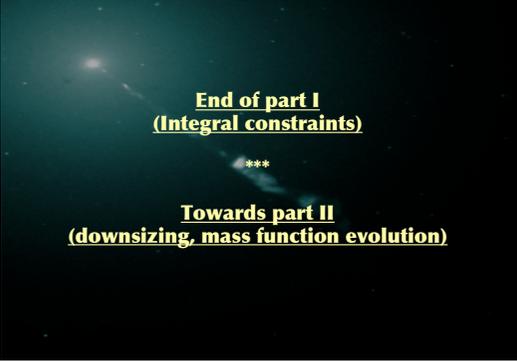




Constraints on avg. SMBH spin ($\langle a \rangle$)



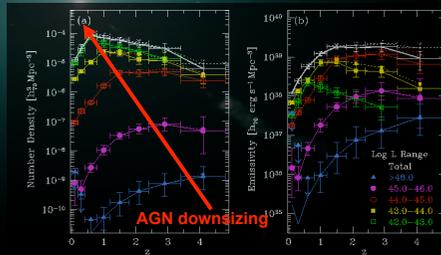




End of part I
(Integral constraints)

Towards part II
(downsizing, mass function evolution)

AGN/SMBH downsizing: clues from X-rays



Ueda et al. 2003; Fiore et al. 2003; Barger et al. 2005; Hasinger et al. 2005

Unveiling the growth of SMBH

1. BH mass can only increase*
2. BHs do not transform into something else as they grow**
3. BHs are like teenagers: they clearly let us know when they grow up (AGN as signposts of BH growth)

*Accreted BH laser mergers can introduce a loss term in the continuity equation

**Merging BH alter the Mass function

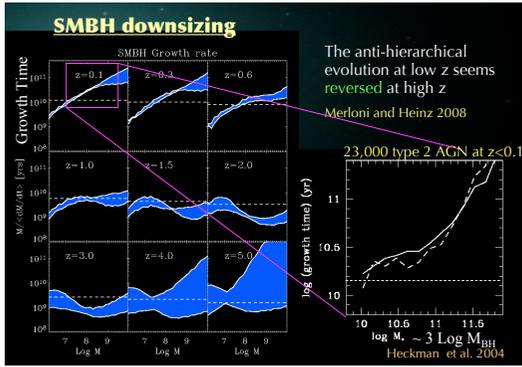
Continuity equation for SMBH growth

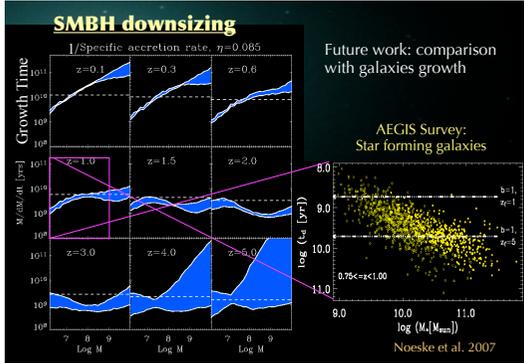
Cavaliere et al. (1973), Small & Blandford (1992), Marconi et al. (2004)

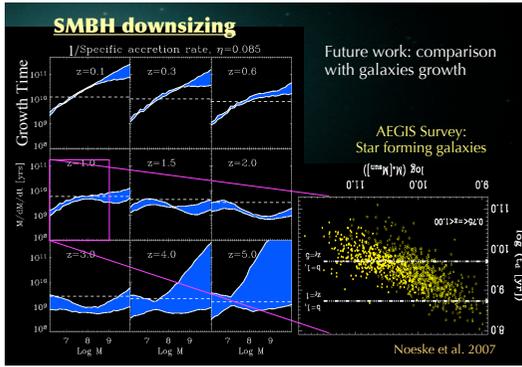
$$\frac{\partial \Phi_M(M, t)}{\partial t} + \frac{\partial [\Phi_M(M, t) \cdot \dot{M}(M, t)]}{\partial M} = 0.$$

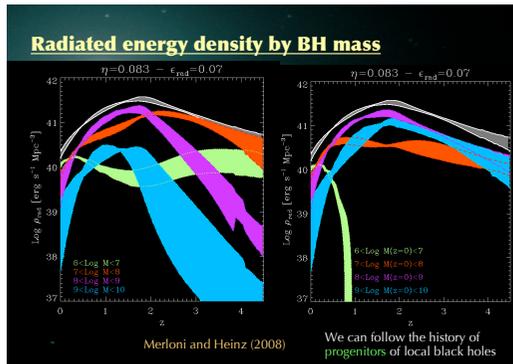
- 1- Need to know simultaneously mass function $\Phi_M(M, t_0)$ and accretion rate distribution
- 2- At any z (or t), it is possible to combine mass function and bolometric luminosity function to calculate accretion rate distribution

Picture from Di Matteo et al. (2007)





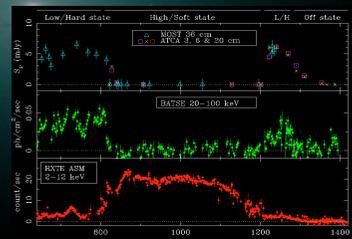




AGN downsizing: changing accretion modes

- SMBH must accrete at lower (average) rates at later times
- Accretion theory (and observations of X-ray Binaries) indicate that
 - The energy output of an accreting BH depends crucially on its accretion rate
 - Low-accretion rate systems tend to be “jet dominated”
- In the recent Cosmology jargon: Quasar mode vs. Radio mode (explosive vs. gentle)

BH transients: window on accretion physics

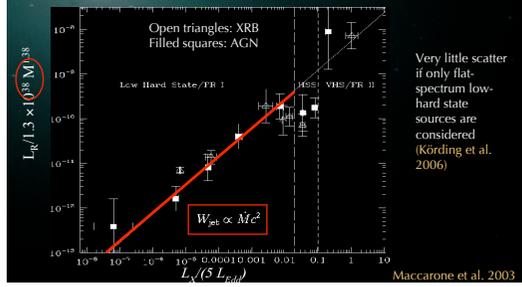


GX 339-4

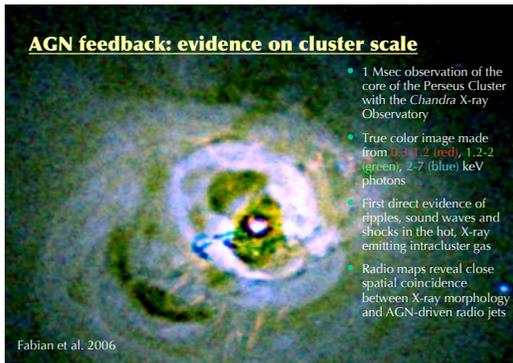
Fender et al, 1999

Radio cores scaling with M and mdot

A "fundamental plane" of active BHs [Merloni et al. 2003; Falcke et al. 2004]



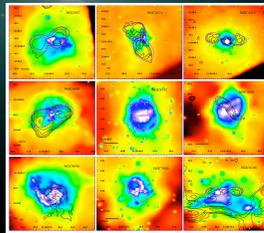
AGN feedback: evidence on cluster scale



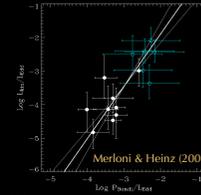
- 1 Msec observation of the core of the Perseus Cluster with the *Chandra* X-ray Observatory
- True color image made from 0.5-1.2 (red), 1.2-2 (green), 2-7 (blue) keV photons
- First direct evidence of ripples, sound waves and shocks in the hot, X-ray emitting intracluster gas
- Radio maps reveal close spatial coincidence between X-ray morphology and AGN-driven radio jets

Fabian et al. 2006

Estimating Jet power in nearby ellipticals

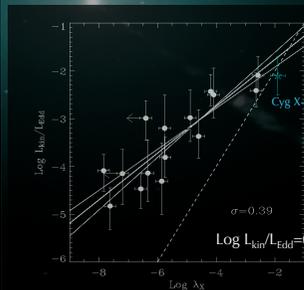


- Allen et al. (2006): correlation between Jet kinetic power and Bondi power



Chandra 0.5-8keV (colour), 1.5GHz VLA radio (contours), bubbles (magenta ellipses).
Allen et al. (2006), Birzan et al. (2004), Rafferty et al. (2006)

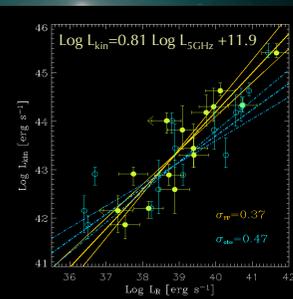
Low Power AGN are jet dominated



- By studying the nuclear properties of the AGN we can establish a link between jet power and accretion power
- The observed slope (0.50 ± 0.045) is perfectly consistent with radiatively inefficient "jet dominated" models (see E. Churazov's talk)

Merloni and Heinz (2007)

Core Radio/ L_{kin} relation: effects of beaming



Slope=0.81

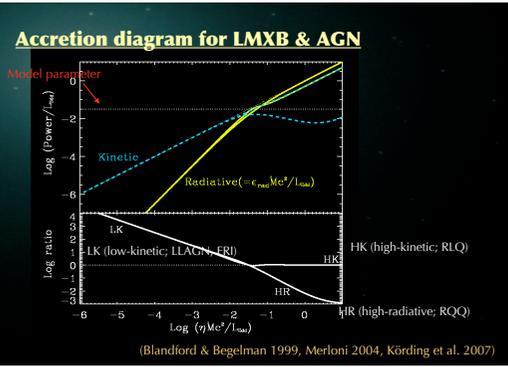
Slope=0.54

Observed L_R (beaming)
Derived from FP relation

Monte Carlo simulation:
Statistical estimates of
mean Lorentz Factor $\Gamma \sim 8$

Not a distance effect:
partial correlation analysis
 $P_{\text{null}} = 2 \times 10^{-4}$

Merloni and Heinz (2007)



What is the “radio mode” of AGN?

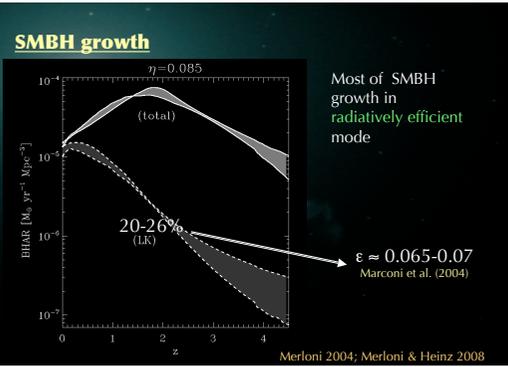
- The energy source that counterbalance cooling in the cores of groups and clusters. Prevents overproduction of massive galaxies at late times (a **FEEDBACK** mode; Croton et al. 2006; Bower et al. 2006)
 - CANNOT be associated to QSOs: their number density declines too fast
- A **FEEDING** mode (hot gas vs. cold gas, Hardcastle et al. 2007)
- **HERE**: The physical state of ALL black holes at low accretion rate ~less than a few % of the Eddington rate (an **ACCRETION** mode)

Continuity equation for SMBH growth

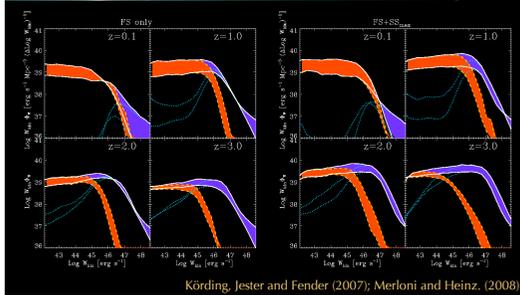
Cavaliere et al. (1973), Small & Blandford (1992), Marconi et al. (2004)

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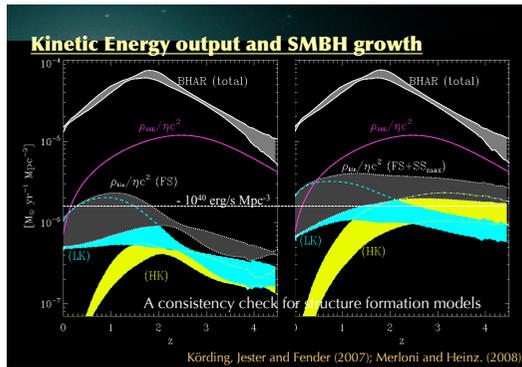
- 1- Need to know simultaneously mass function $\Phi_M(M, t_0)$ and accretion rate distribution
- 2- At any z (or t), it is possible to combine mass function and bolometric luminosity function to calculate accretion rate distribution
- 3- For each accretion mode, use observed scaling relations between (core) radio and X-ray (bolometric) luminosity to "couple in" the evolution of radio LF

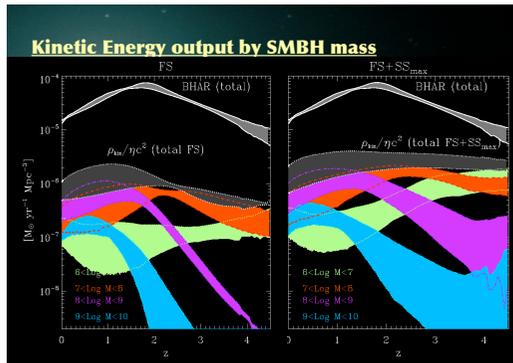


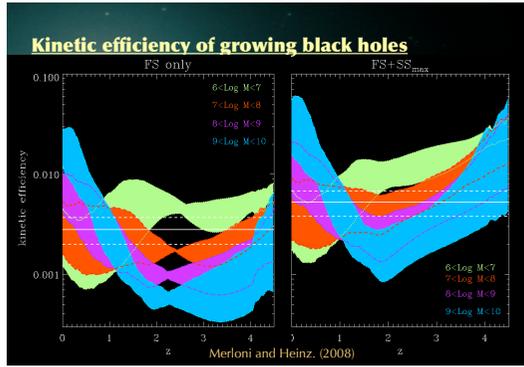
AGN Kinetic Luminosity function

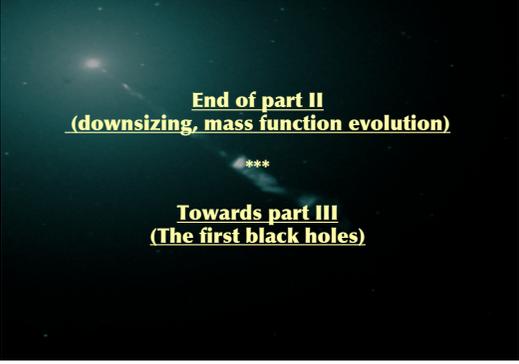


Körding, Jester and Fender (2007); Merloni and Heinz, (2008)



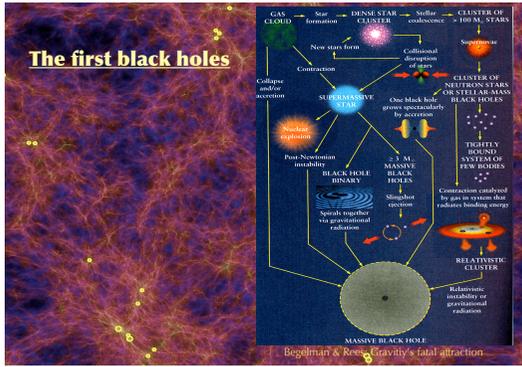




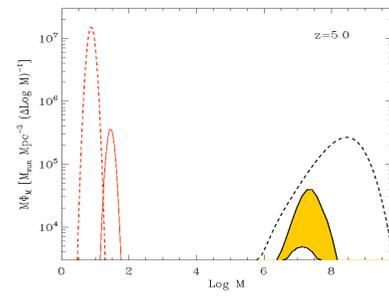


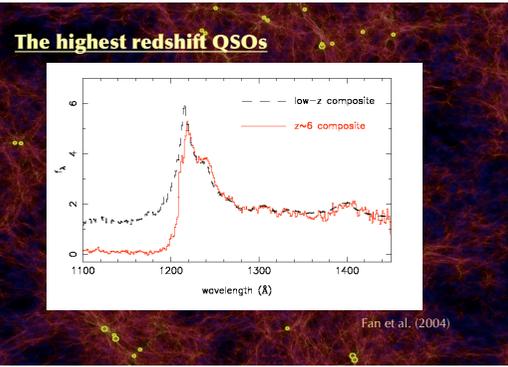
End of part II
(downsizing, mass function evolution)

Towards part III
(The first black holes)

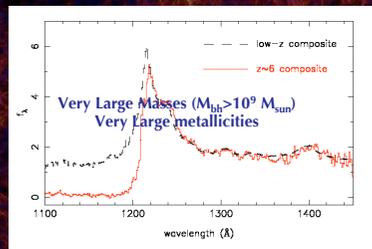


Black Holes mass function evolution





The highest redshift QSOs



Fan et al. (2004)

The highest redshift QSOs: the time problem

Available time from $z=30$ till $z=6$ is about 0.8 Gyr.
Assume SMBH growth at the **Eddington limit**:

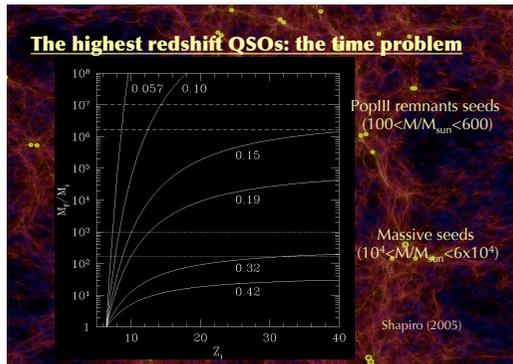
$$dM/dt = (1 - \eta) L_{\text{edd}} / (\epsilon_{\text{rad}} c^2)$$

Assuming, for simplicity, $\eta = \epsilon_{\text{rad}}$

$$M(t) = M(0) \exp [(1 - \eta) / \eta * t / t_{\text{edd}}]$$

With $t_{\text{edd}} = 0.45$ Gyr

↓
Upper limit on η !



The highest redshift QSOs: efficiency problem

BH growth by accretion via standard thin disc from an initial state with $M=M_i$ and spin $a=a_i$

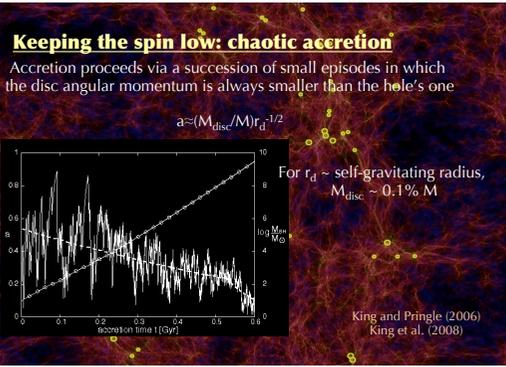
•• Spin evolves according to:

$$a_f = (r_{\text{ISCO},i}^{1/2}/3) * \chi * [4 - (3r_{\text{ISCO},i} * \chi^2 - 2)^{1/2}]$$

Where $\chi = M_f/M_i$

→ An initially non-rotating BH is spun up to $a_f=1$ if $\chi = M_f/M_i = 6^{1/2} \approx 2.45$

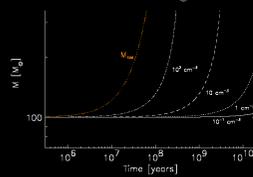
↓
Prolonged coherent accretion episodes imply high efficiency!



More problems: fueling rates

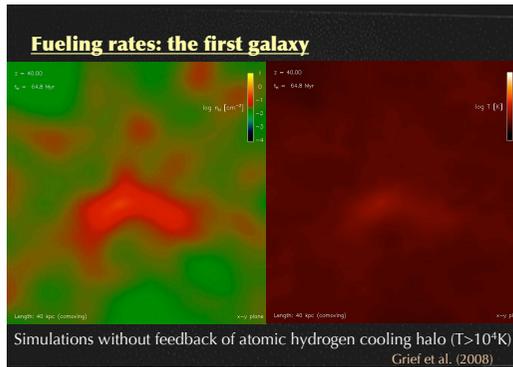
$$\frac{dM_{\text{BH}}}{dt} \Big|_{\text{Bondi}} = \frac{2\pi(1-\eta)(GM_{\text{BH}})^2 m_{\text{H}} n}{c^3} \quad \frac{M_{\text{BH}}}{M_i} = \frac{1}{1-t/t_{\text{Bondi}}}$$

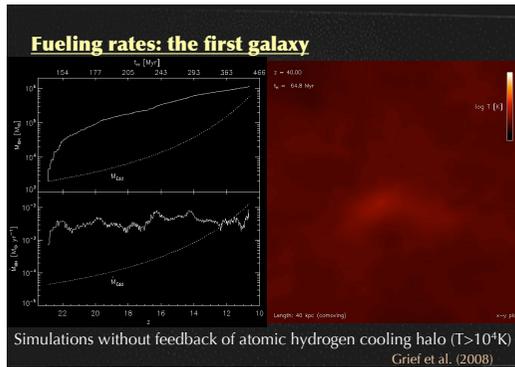
$$t_{\text{Bondi}} \simeq 3.5 \times 10^8 \left(\frac{M_i}{100 M_{\odot}} \right)^{-1} \left(\frac{0.9}{1-\eta} \right) \left(\frac{T}{200 \text{K}} \right) \left(\frac{n}{\text{cm}^{-3}} \right)^{-1} \text{ yrs}$$

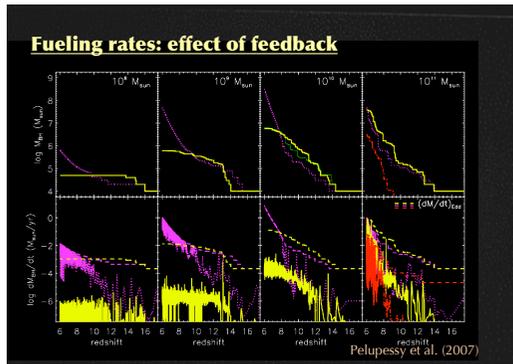


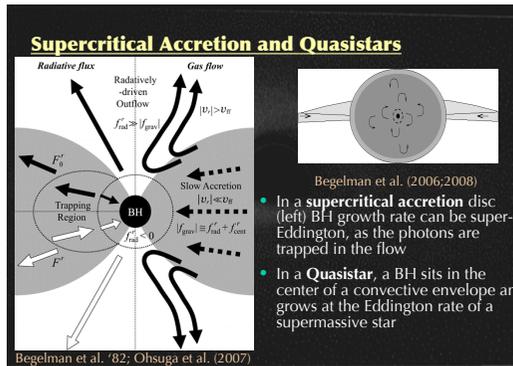
Johnson and Bromm (2007)

Simulations **without feedback**
found very low gas densities
($< 1 \text{ cm}^{-3}$ for $\sim 10^8$ yrs)
around first PopIII stars

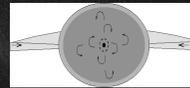
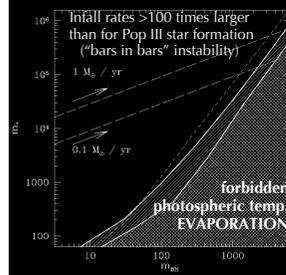








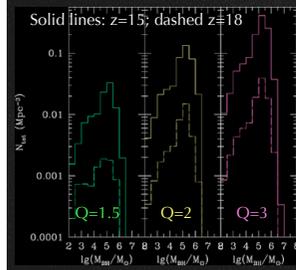
Growth and death of a Quasistar



- A PopIII **Quasistar** may form by direct collapse in the core a pristine pre-galactic halos where atomic hydrogen cooling is efficient
- In less than a million years a 10^4 BH may be born

Begelman et al. (2008)

Predictions: early (seed) BH mass functions



- Using a model for gas collapse in pre-galactic halos that counts for **gravitational stability and fragmentation**
- Q is a **stability** parameter: lower Q implies more stable discs

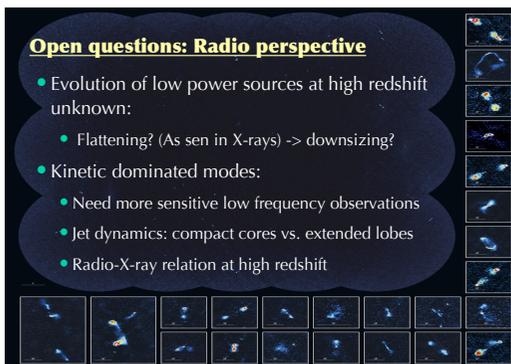
Volonteri et al. (2008)

Conclusions: SMBH evolution at $z < 4$

- SMBH grow with a **broad accretion rate distribution**
- Most of SMBH growth occurred in **radiatively efficient** episodes of accretion.
- The **anti-hierarchical trend** is clearly seen in the low- z evolution of SMBH **mass function**. Reversal at higher z ?
- Feedback from “Low-luminosity AGN” are most likely **dominated by kinetic energy**
- The efficiency with which growing black holes convert mass into mechanical energy is **0.3-0.5%** (but strongly dependent on BH mass and redshift).

Open questions: Radio perspective

- Evolution of low power sources at high redshift unknown:
 - Flattening? (As seen in X-rays) \rightarrow downsizing?
- Kinetic dominated modes:
 - Need more sensitive low frequency observations
 - Jet dynamics: compact cores vs. extended lobes
 - Radio-X-ray relation at high redshift



Open questions: the first Black holes

- What typical mass (where the peak of Mass Function)
 - PopIII star remnants vs. direct collapse
- What growth mode:
 - Standard accretion vs. chaotic accretion vs. super-critical accretion
- What is the mass of typical host halo
 - Correlation function
 - Early M-sigma?



The M87 jet
Hubble Heritage Project
<http://heritage.stsci.edu/2000/20/index.html>

Footnotes:

*Kicked BH after mergers can introduce a loss term in the continuity equation

**Merging BH alter the Mass function

