Infrared Emission from the dusty veil around AGN

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In collaboration with

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"From the circumnuclear disk in the Galactic Center to thick, obscuring tori of AGNs"
B. Vollmer, T. Beckert, W. J. Duschl 2004, A&A 413, 949

"The dynamical structure of a thick cloudy torus" T. Beckert, W. J. Duschl (A&A, in press)

Overview

- NIR-Observations of NGC 1068 (Talk by G. Weigelt)
- Viscosity & Cloud Collisions in a Torus
- Consequences for NIR imaging, IRspectra, and Unified Schemes

Motivation



- Are the conclusions of Krolik & Begelman (1988) still valid ?
- How to support and maintain the thickness of the torus ?
- Is the "torus" really a torus or a cloudy veil.

/e are used to think in (model) pictures. -> Are the pictures correct?

The Mass Distribution (GC)



NGC 1068 — the classical case for the unified model

- Distance 14.4 Mpc \rightarrow 1'' \Leftrightarrow 70 pc
- Seyfert 2 with a hidden Sy 1 core seen in polarized lines
- Compton thick in X-rays: N_H > 5 10²⁴ cm⁻²
- Conical (collimated) Narrow Line Region

•Massive outflow & Weak radio jet

•Masers in a almost perpendicular disk tracing (??) rotation Mass estimate ~ 1.2 10⁷ M_{solar}

•Luminosity: 0.4–2 1045 erg/s

200 mas

K'

NGC 1068 – NIR speckle images





Weigelt et al. 200 and his talk here

Core: 18 x 40 mas = 1.3 x 2.8 pc Flux: 350 mJy

& extended emission

Infrared Interferometry (VLTI)





 \longrightarrow < 0.2 pc \implies Substructure (Wittkowski et al. 2004)

MIDI: Two-Comp.: T=320 K (30 x 50 mas) & T > 800 K (≤ 10 mas) (Jaffe et al. 2004)

The Model: Conditions - Assumptions

- Dust can only survive in cold clouds !!!
 σ ~ 50 km/s → T~10⁵ K
- Equilibrium structure in the combined potential of Black Hole & quasi-isothermal star cluster
- Radial accretion flow due to cloud-cloud collisions
- Mass is supplied at an outer radius (ISM, starburst ring, bar driven accretion)

The Accretion Scenario

Cloud-Cloud Interactions:

effective viscosity (Goldreich & Tremaine 1978)

$$v_{\rm eff} = \frac{\tau}{1 + \tau^2} \frac{\sigma^2}{\Omega}$$

dimensionless collision frequency $\tau = \omega_{coll}/\Omega$

- Mass accretion from differential rotation & angular momentum redistribution
- Vertical scale height

$$H = \sigma / \Omega = l_{\rm coll} \tau$$

The collisional particle disk:

- Symmetries of a thin accretion disk
- Keplerian rotation & vertical hydrostatic equilibrium
- Triaxial Gaussian velocity distribution
- Viscosity follows

$$v = \sin(2\delta) \frac{\sigma_2^2 - \sigma_1^2}{2R(-\Omega')}$$

$$N = \frac{1}{2}\sin(2\delta)\Sigma R^2(\sigma_2^2 - \sigma_1^2)$$

Cloud Collisions



Goldreich & Tremaine (1978):
Only momentum along λ is lost in inelastic collisions (assumption)
Elasticity ε
Coefficient of restitution

Only $\frac{1}{2}$ (1- ϵ^2) of the kinetic energy is dissipated.

Results for a thin disk



Modifications for a thick torus

- 1. Advective terms (change of σ^2 and compression)
- No collective effects (enhanced viscosity via non-local interactions; Wisdom & Tremaine 1988)

$$\sigma >> R_{\text{Cloud}} \left| \frac{\partial \Omega}{\partial \log R} \right| \quad \Rightarrow \quad \Phi_{\text{V}} << 0.35 \tau$$

 3. Elasticities ε as low as 0.3 are possible (45 % of kinetic energy can be dissipated in collisions)
 ε determines τ

The vertical structure



Exact solution for the density in an arbitrary

external potential

with a vertical cut-off height

Example for NGC 1068

Cloud Properties

- Tidal forces limit size of the largest clouds (shear-limit)
- Largest Clouds dominate appearance
- Quasi-stable clouds hold together by self-gravity

$$R_{\text{Cloud}}(c_S, M_{\text{Cloud}}) \rightarrow l_{\text{Coll}} = \frac{1}{n \pi R_{\text{Cloud}}^2}$$

Typical cloud mass

$$M \approx 50 M_{\rm solar}$$

Obscuration implies





Radial Structure



10

10

Red: m = 30

Blue: m = 10

Black: m = 3

m : mass accretion rate in units of the Eddington rate of the black hole

Timescales

- Geometrically thick accretion flows (tori) rotate slightly sub-Keplerian (4 · 10⁵ yr at 10 pc) Vertical hydrostatic equilibrium can be achieved on an orbital timescale
- 2. The accretion timescale (viscous timescale) < 10⁶ yr at 10 pc $t_{\rm acc} \propto \tau^{-1} \dot{M}^{-1} R$

3. Collapse time is the shortest timescale involved

$$t_{\text{Collapse}} = \frac{0.2}{\tau \left(1 - \varepsilon\right)} t_{\text{Orbit}}$$

Without energy gains from accretion or other processes a $\tau \sim 1$ torus would collapse to a thin disk within an orbital timescale Energy gains from accretion is required

Appearance in NIR

Based on the formalism of Nenkova, Ivezić & Elitzur (2002)



Dynamic range: 10⁵; spatial scale in units of the sublimation radius

The IR-Spectrum

Based on the formalism of Nenkova, Ivezić & Elitzur (2002)



Sublimation radius 0.9 pc ; $L_{bol} = 8 \ 10^{45} \ erg/s$ Degeneracy Mdot - Inclination

Torus & Unified Scheme



Predictions for Evolution

• Surface density

$$\Sigma \sim \tau \, \frac{M(R)}{2R^2} \frac{c_s}{v_{\text{circ}}}$$

Thickness of Torus

$$H_{R} \sim \left(\frac{1+\tau^{2}}{\tau^{2}}\frac{R\dot{M}}{M(R)c_{s}}\right)^{1/2}$$

• Mean free path

$$\frac{l_{\rm coll}}{R} = \frac{H}{\tau R}$$

Mean number of clouds (midplane)

$$N \propto \tau \ \dot{M}^{-1/2}$$

Summary

- Dusty tori can be modelled as thick, clumpy accretion flows
- Geometrically thick tori need huge mass accretion rates



- Next question:

How do cloud collisions really look like ? How does circumnuclear starformation feed the torus ?

