

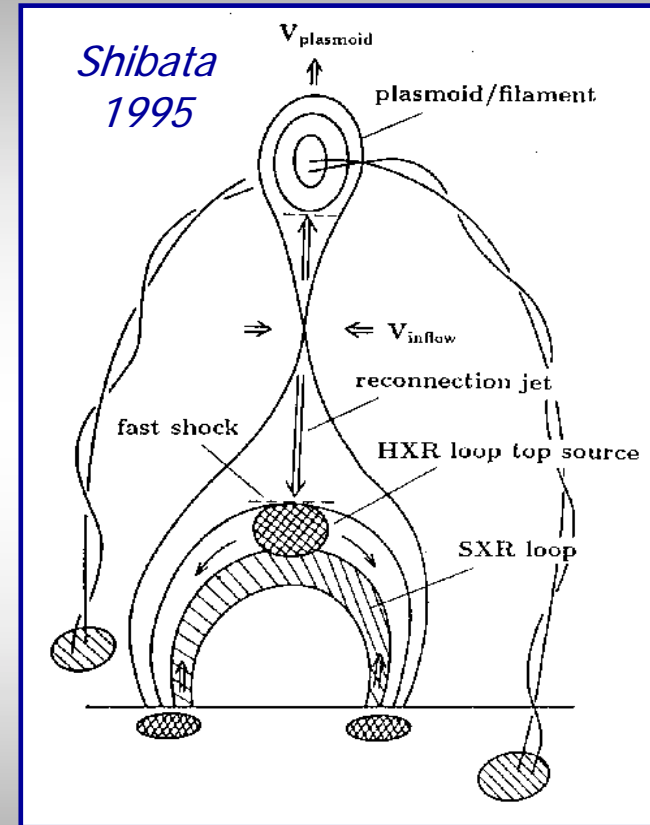
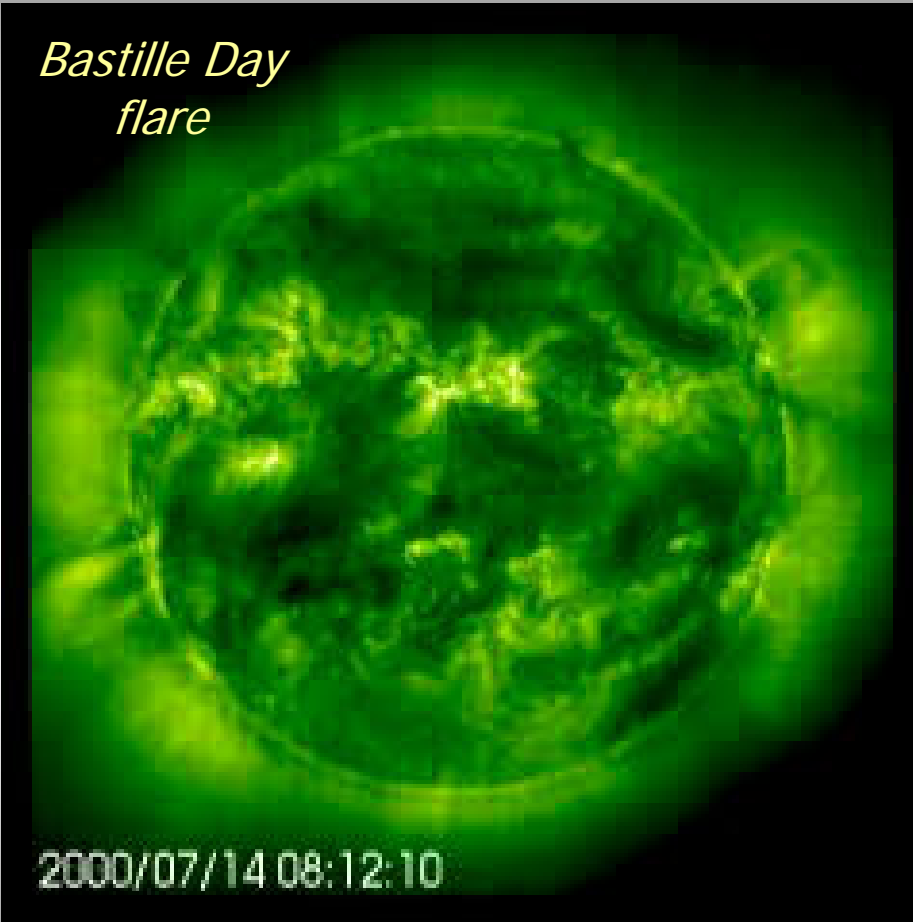
Particle acceleration at 3D magnetic reconnection sites

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Solar flares

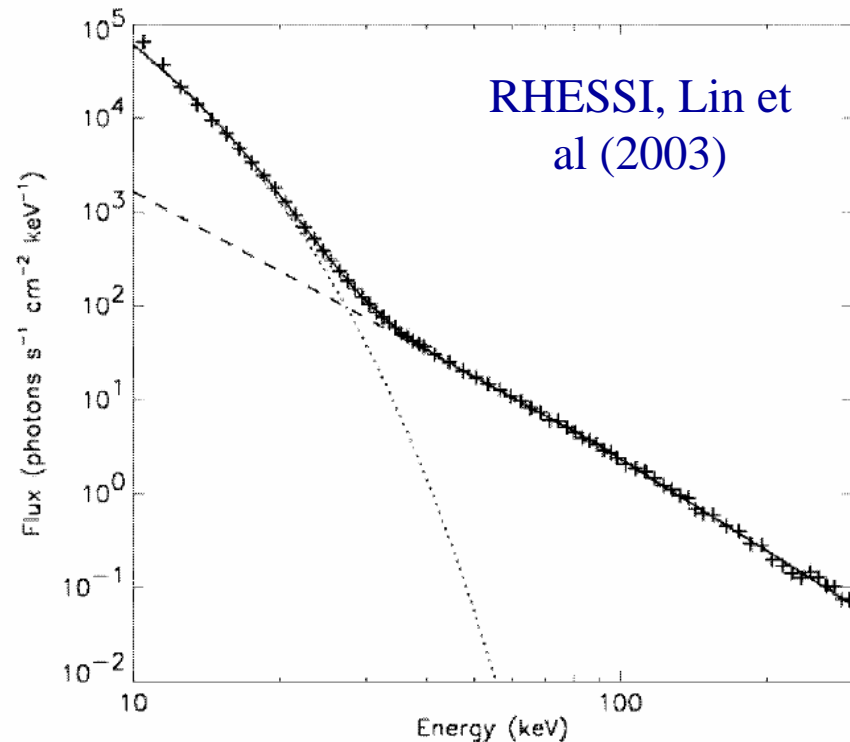
*Bastille Day
flare*



- Release of stored magnetic energy by magnetic reconnection

Particle acceleration in solar flares

- Flares produce large numbers of high energy (non thermal) protons and electrons as well as heating plasma
- Magnetic reconnection is believed to be fundamental process of energy release in solar flares
- Are direct electric fields associated with reconnection responsible for generating high energy particles?
- Similar processes may accelerate particles in other contexts – Earth's magnetosphere, extragalactic jets, pulsars etc



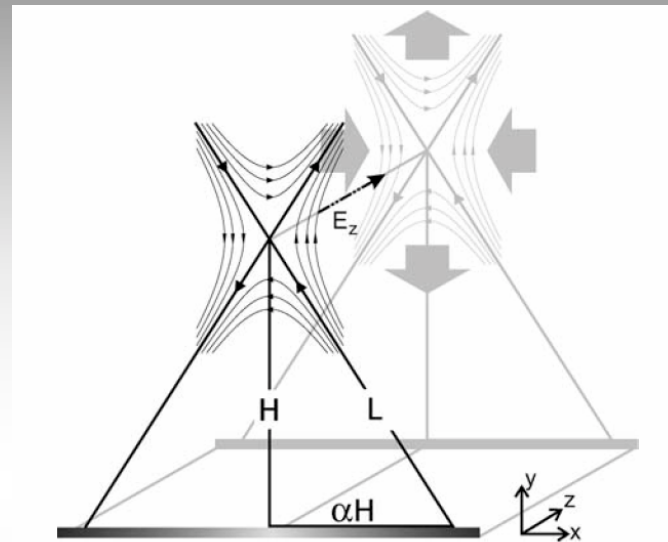
Test particle approach to modelling particle acceleration

- Take magnetic and electric field configuration representative of reconnection
- Neglect the fields generated by the test particles and (usually) neglect collisions of test particles with the background plasma – no energy losses
- Integrate equations of motion numerically
- We undertake test particle simulations in 3D fields

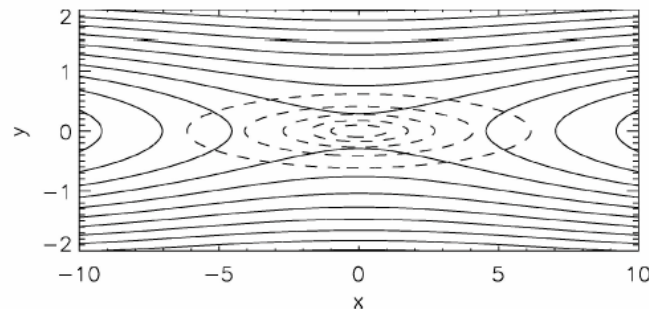
$$\frac{d\mathbf{x}}{dt} = \frac{\mathbf{p}}{m\gamma}$$
$$\frac{d\mathbf{p}}{dt} = q \left(\mathbf{E} + \frac{1}{c} \frac{\mathbf{p}}{m\gamma} \times \mathbf{B} \right)$$

2D test particle models

- X-point or current sheet configuration
- Invariance of all quantities in the 'z' direction - with or without "guide field" B_z
- Many studies e.g. Vekstein and Browning (1997), Deeg et al (1990), Browning & Vekstein (2001), Zharkova & Gordovsky (2005), Hamilton et al (2003).. etc



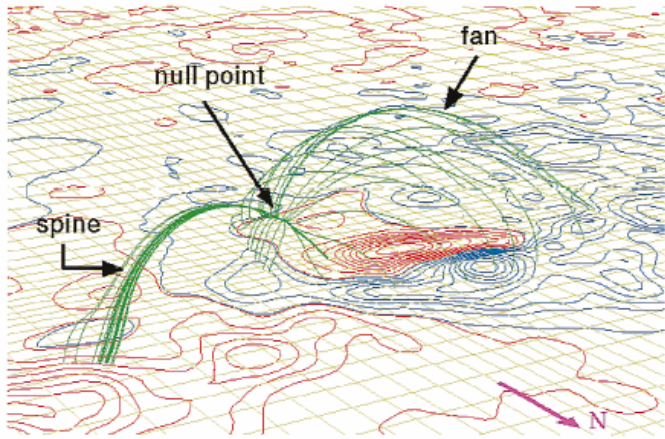
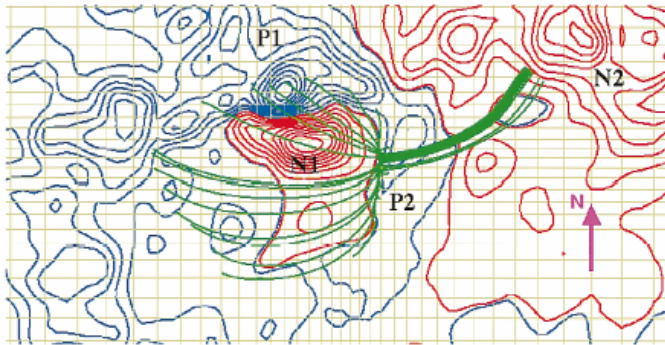
Hannah & Fletcher, 2006



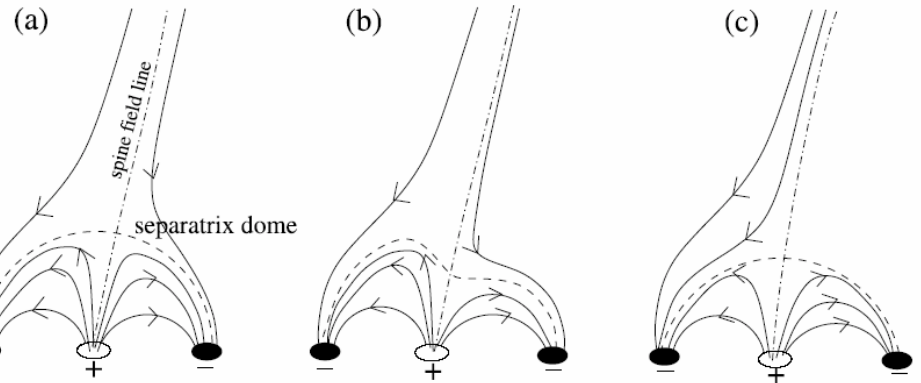
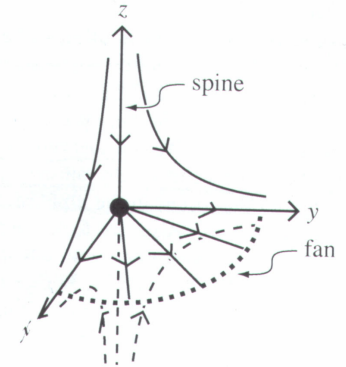
Wood & Neukirch, 2005

Figure 1. An example of the field configuration used in this paper (using $L = 10$). Note that for graphical reasons the configuration is not shown with the correct aspect ratio, so that it is actually a lot more elongated in the x -direction than shown here. The *solid lines* are projections of magnetic field lines onto the x - y plane (contours of the flux function A). The non-ideal region in which a field-aligned component of the electric field exists is outlined by the *dashed contours*.

3D null points in flares



3D null points

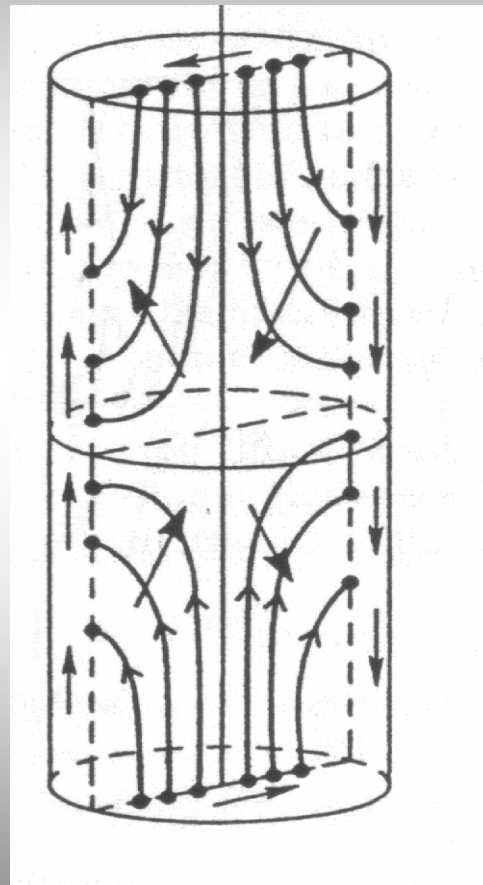
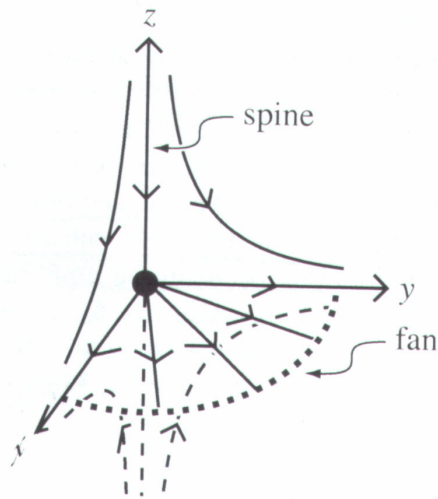


Aulanier et al (2000)

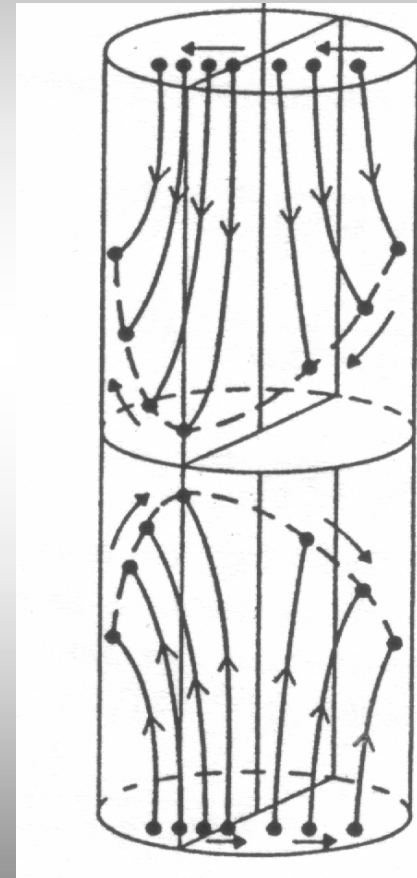
Fletcher et al (2001)

Spine and fan reconnection

Priest and Titov (1996)

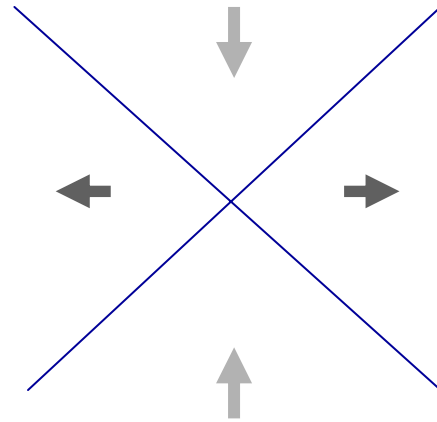
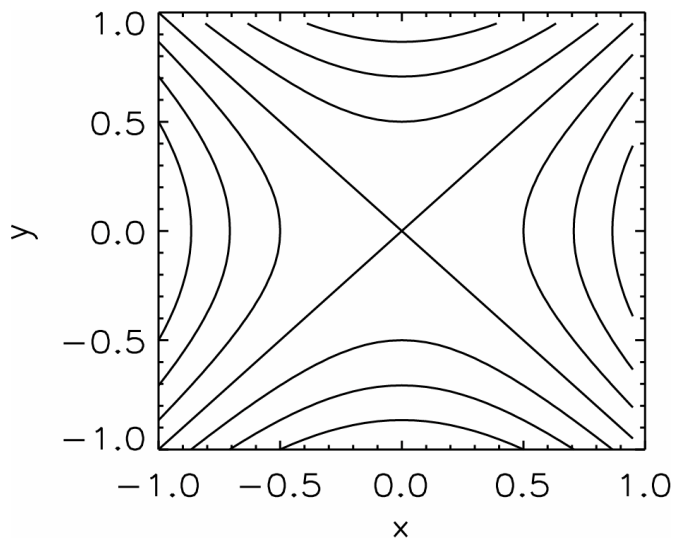


Spine

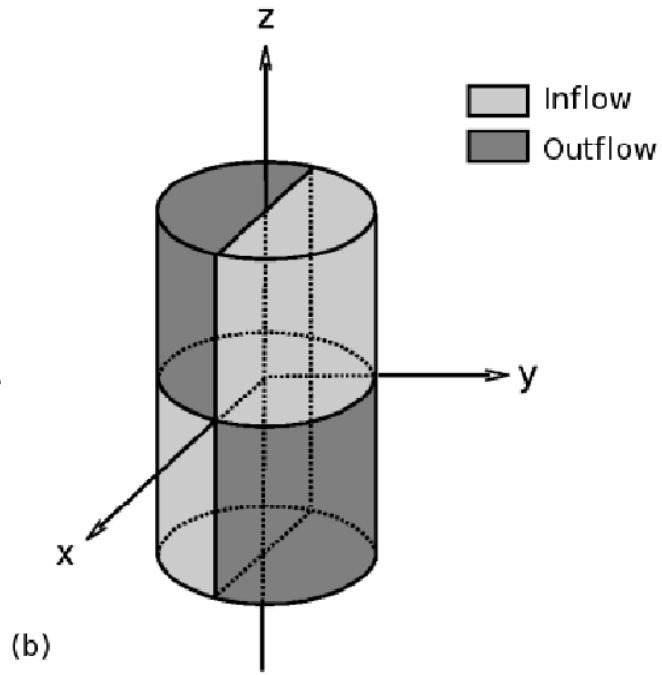
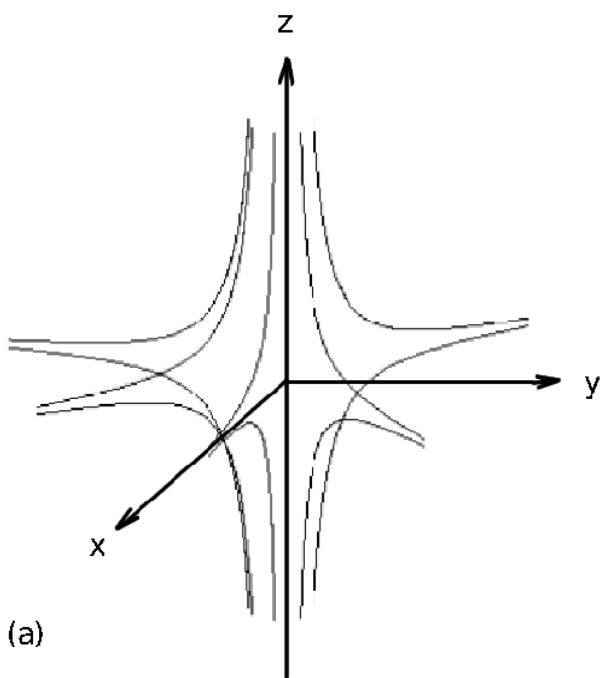


Fan

$$\mathbf{B} = B_0 \frac{R}{L} \mathbf{e}_R - 2 B_0 \frac{z}{L} \mathbf{e}_z$$



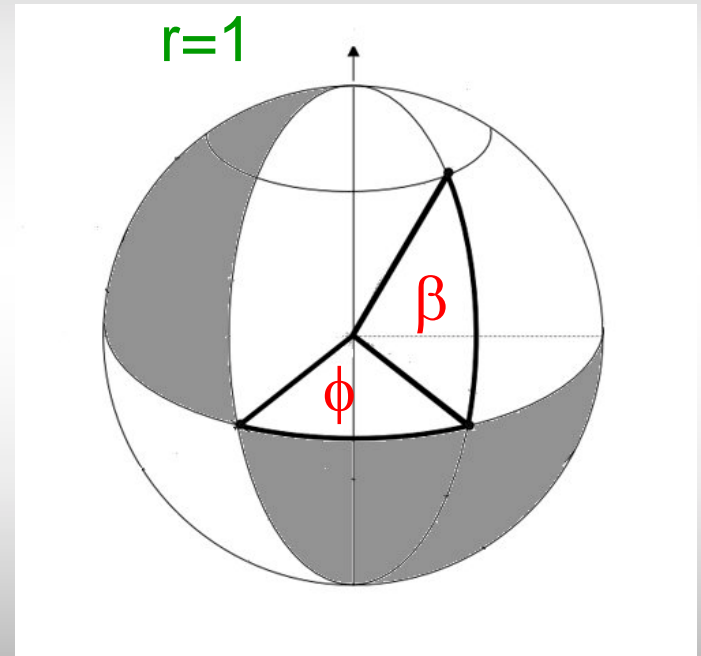
2D



3D
(spine)

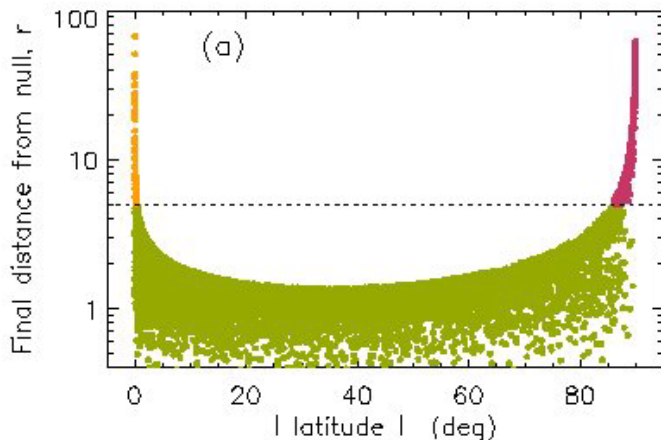
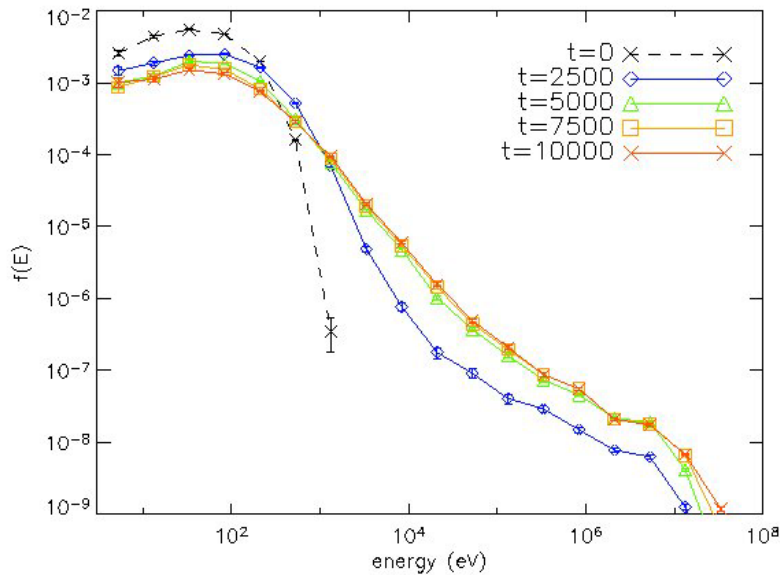
Acceleration of test particles

- Consider only outer “ideal” region - acceleration due to non-uniform $E \times B$ drift
- Particles become unmagnetised (large Larmor radii) near spine/fan
- Ignore (for now) acceleration due to E_{\parallel} in resistive region (current sheet/filament)
- Inject ensemble of particles on a sphere of normalised radius $r=1$
- 10,000 protons. Start in inflow regions, with uniform random positions. Velocities from Maxwellian ($T = 1$ million K)



□ Inflow □ Outflow

Results – see poster!



- Calculate energy spectra of particles in spine and fan reconnection
- Also analyse spatial distribution and pitch angles of particles
- Energy gain depends strongly on entry location – investigate individual particle trajectories
- In spine reconnection, a jet of high energy particles escapes along the spine – also trapped high energy particles are generated
- Acceleration is less effective in fan reconnection
- Scaling with reconnection parameters is considered