

# **“Dynamical Constraints for Obscuration by Dust in a Torus”**

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# In collaboration with

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- Wolfgang Duschl & Werner Tscharnuter (Heidelberg)
- Gerd Weigelt (Bonn)
- Moshe Elitzur (Lexington, KY)

„From the circumnuclear disk in the Galactic Center to thick, obscuring tori of AGNs“

B. Vollmer, T. Beckert, W. J. Duschl 2004, A&A 413, 949


„The dynamical structure of a thick cloudy torus“

T. Beckert, W. J. Duschl (A&A accepted)

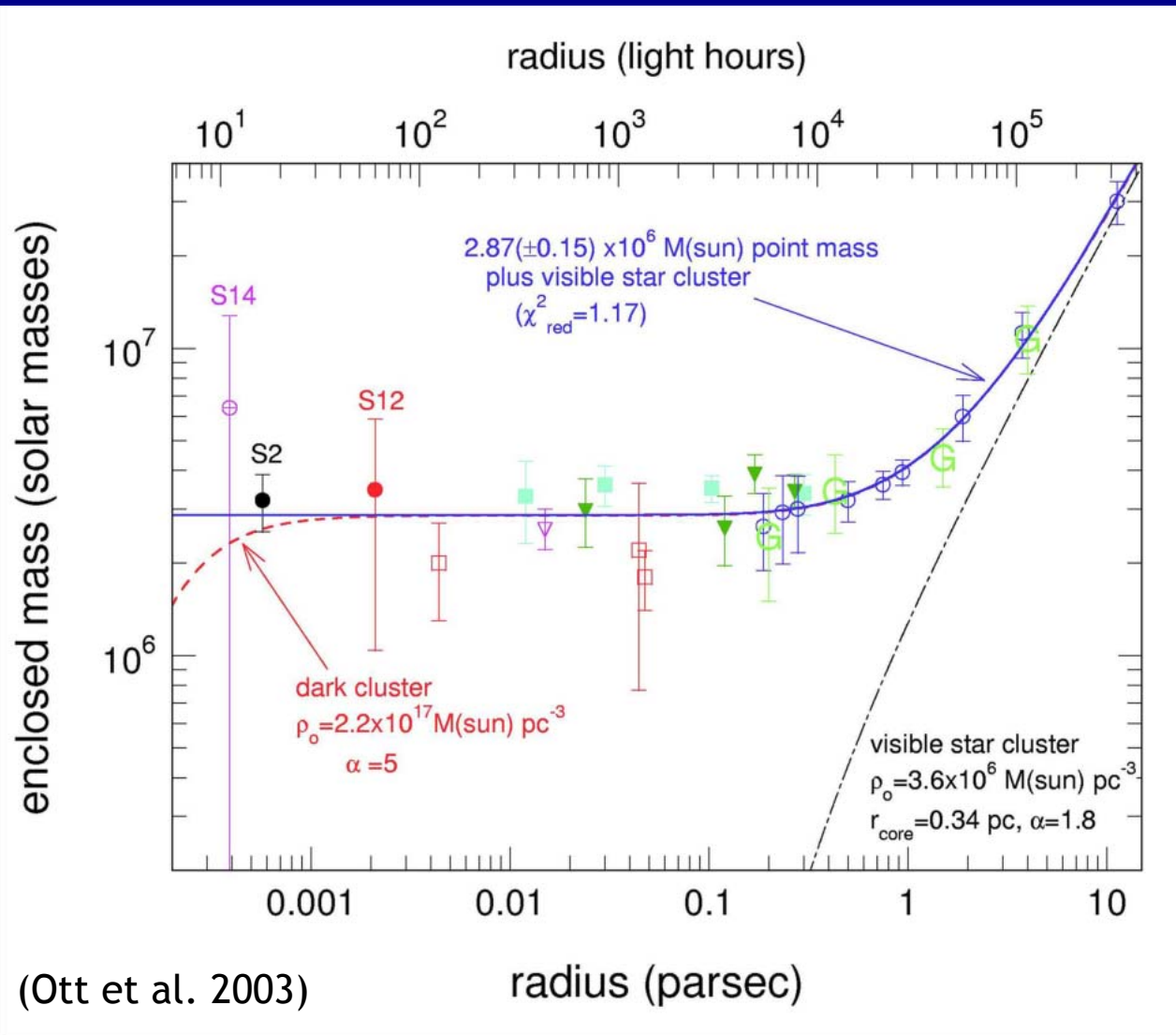
# Overview

- NIR-Observations of NGC 1068
- Viscosity & Cloud Collisions in a Torus
- Consequences for NIR imaging and Unified Schemes

# Motivation

- Unified schemes of AGN assume dusty torus
- What are the requirements to make it work ?
- Are the conclusions of Krolik & Begelman (1988) still valid ?
-  How to support and maintain the thickness of the torus ?

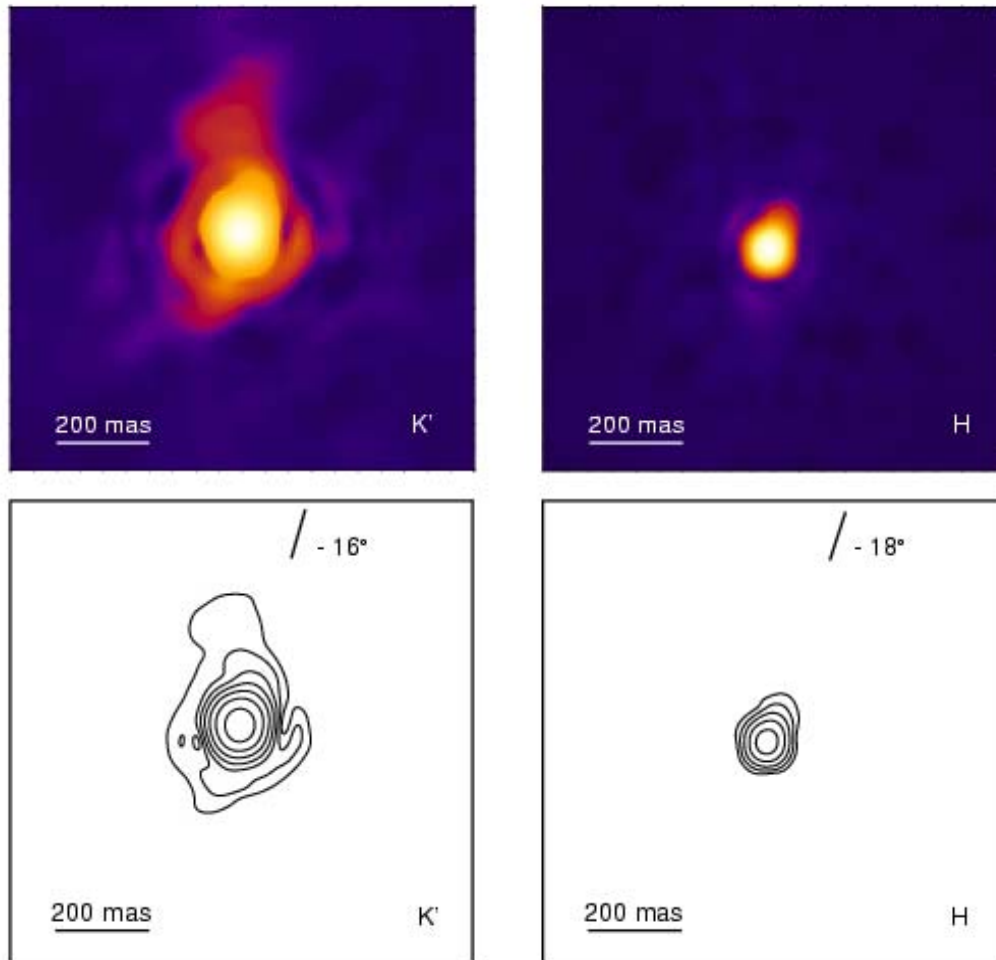
# The Mass Distribution (GC)



# NGC 1068 – the classical case for the unified model

- Distance 14.4 Mpc  $\rightarrow$  1''  $\leftrightarrow$  70 pc
  - Seyfert 2 with a hidden Sy 1 core seen in polarized lines
  - Compton thick in X-rays:  $N_{\text{H}} > 5 \cdot 10^{24} \text{ cm}^{-2}$
  - Conical (collimated) Narrow Line Region
  - Massive outflow & Weak radio jet
- 
- Masers in a almost perpendicular disk tracing (??) rotation  
Mass estimate  $\sim 1.2 \cdot 10^7 M_{\text{solar}}$
  - Luminosity:  $0.4\text{--}2 \cdot 10^{45} \text{ erg/s}$
-

# NGC 1068 — NIR speckle images



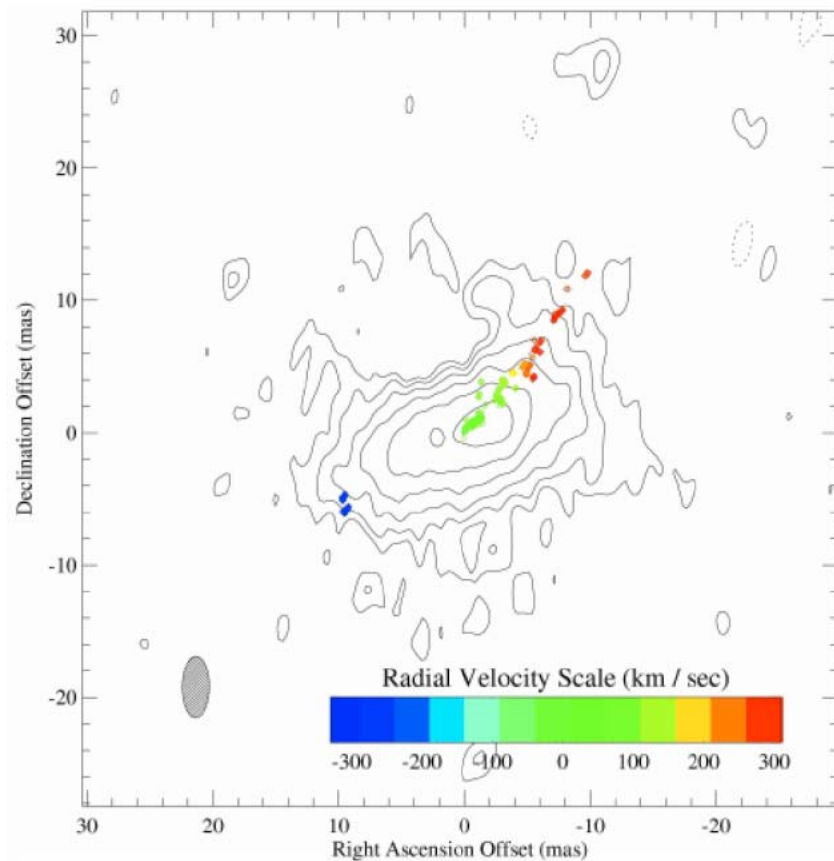
(Weigelt et al.  
2004,  
A&A accepted)

Core:  
18 x 40 mas  
= 1.3 x 2.8 pc  
Flux: 350 mJy

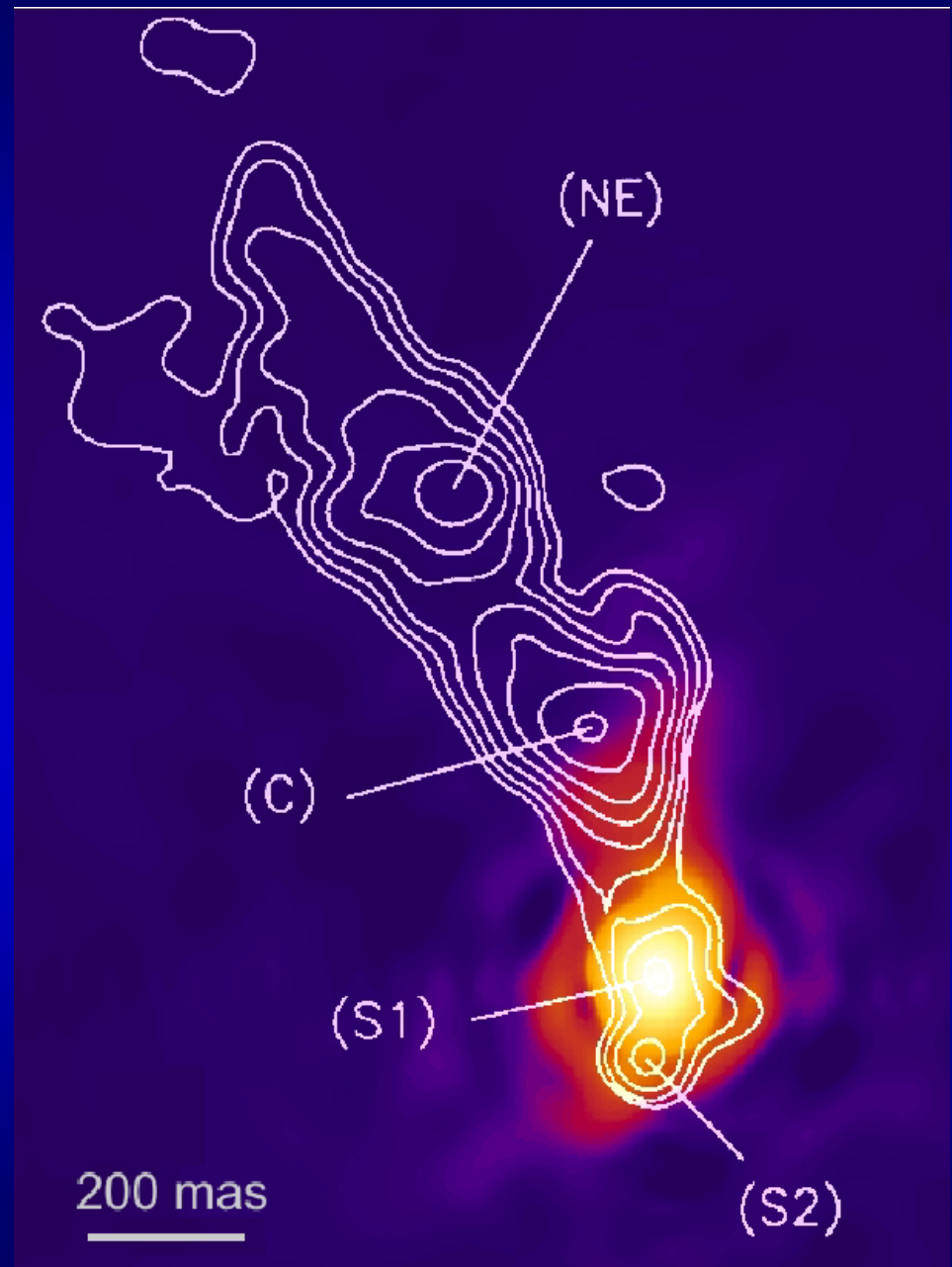
& extended  
emission

# NGC 1068

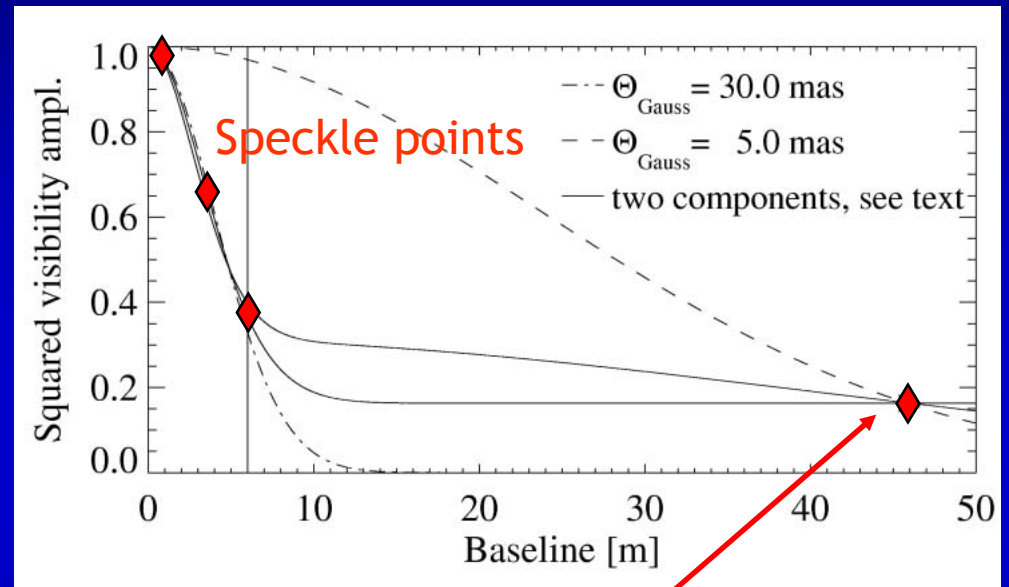
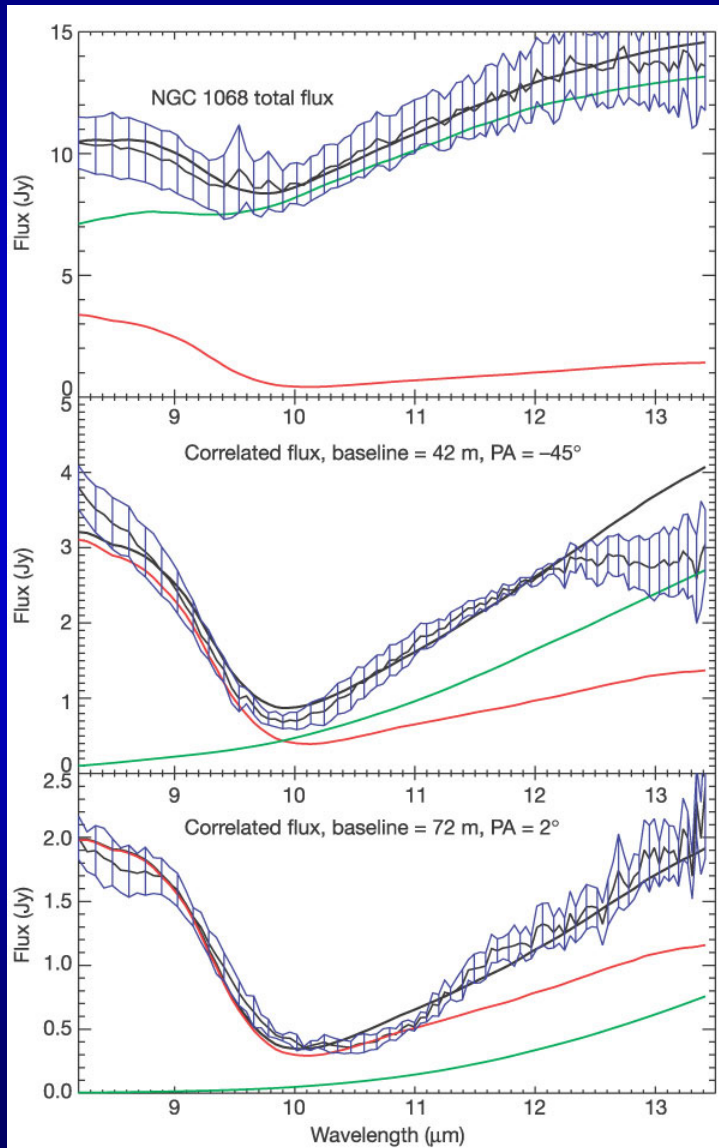
## The radio-NIR view



(Gallimore et al. 2004, preprint)



# Infrared Interferometry (VLTI)



VINCI ( $2.2 \mu\text{m}$ )  $T = T_{\text{sub}} (< 3 \text{ mas})$

$\longrightarrow < 0.2 \text{ pc} \longrightarrow$  Substructure

(Wittkowski et al. 2004)

MIDI: Two-Comp.:  $T=320 \text{ K}$  ( $30 \times 50 \text{ mas}$ )

&  $T > 800 \text{ K}$  ( $\leq 10 \text{ mas}$ )

(Jaffe et al. 2004)

# The Model: Conditions - Assumptions

- Dust can only survive in cold clouds
- **Equilibrium structure** in the combined potential of Black Hole & quasi-isothermal star cluster

$$\dot{M} = \text{const.}$$

- **Radial accretion flow due to cloud-cloud collisions**
- Mass is supplied at an outer radius (ISM, starburst ring, bar driven accretion)

# The Accretion Scenario

- Cloud-Cloud Interactions:

**effective viscosity**

(Goldreich & Tremaine 1978)

$$\nu_{\text{eff}} = \frac{\tau}{1 + \tau^2} \frac{\sigma^2}{\Omega}$$

dimensionless collision frequency

$$\tau = \omega_{\text{coll}} / \Omega$$

- Mass accretion from differential rotation & angular momentum redistribution
- Vertical scale height

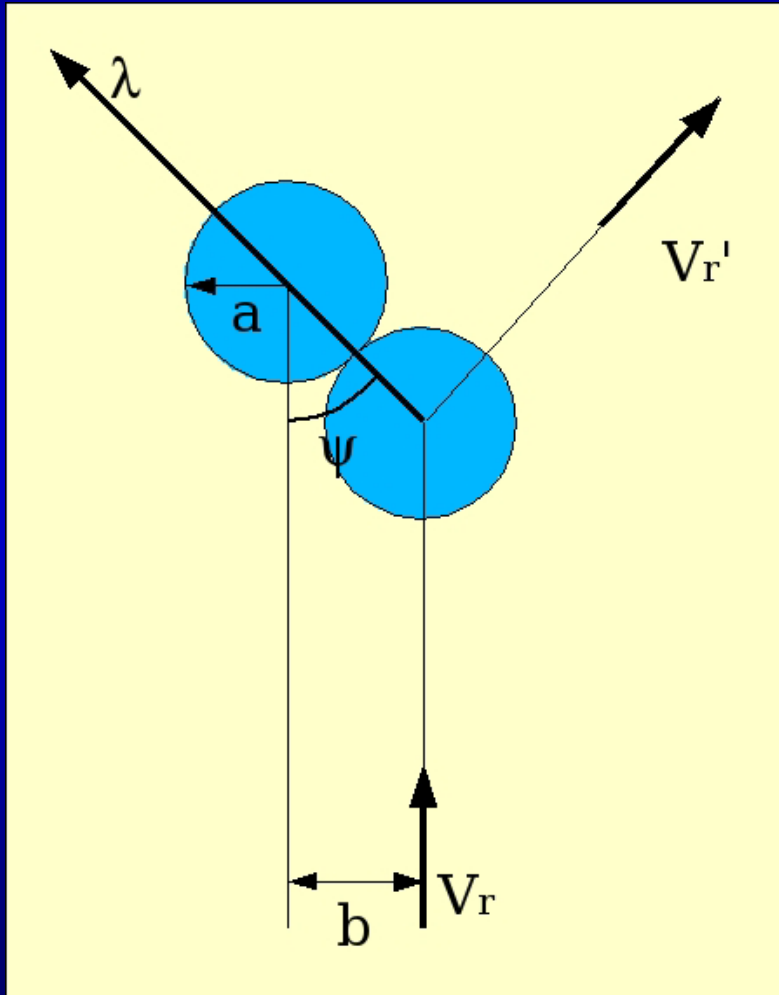
$$H = \sigma / \Omega = l_{\text{coll}} \tau$$

## The collisional particle disk:

- Symmetries of a thin accretion disk
- Keplerian rotation & vertical hydrostatic equilibrium
- Triaxial Gaussian velocity distribution
- Viscosity follows

$$v = \frac{P_{r\phi}}{n(-\Omega')} = \sin(2\delta) \frac{\sigma_2^2 - \sigma_1^2}{2R(-\Omega')}$$

# Cloud Collisions



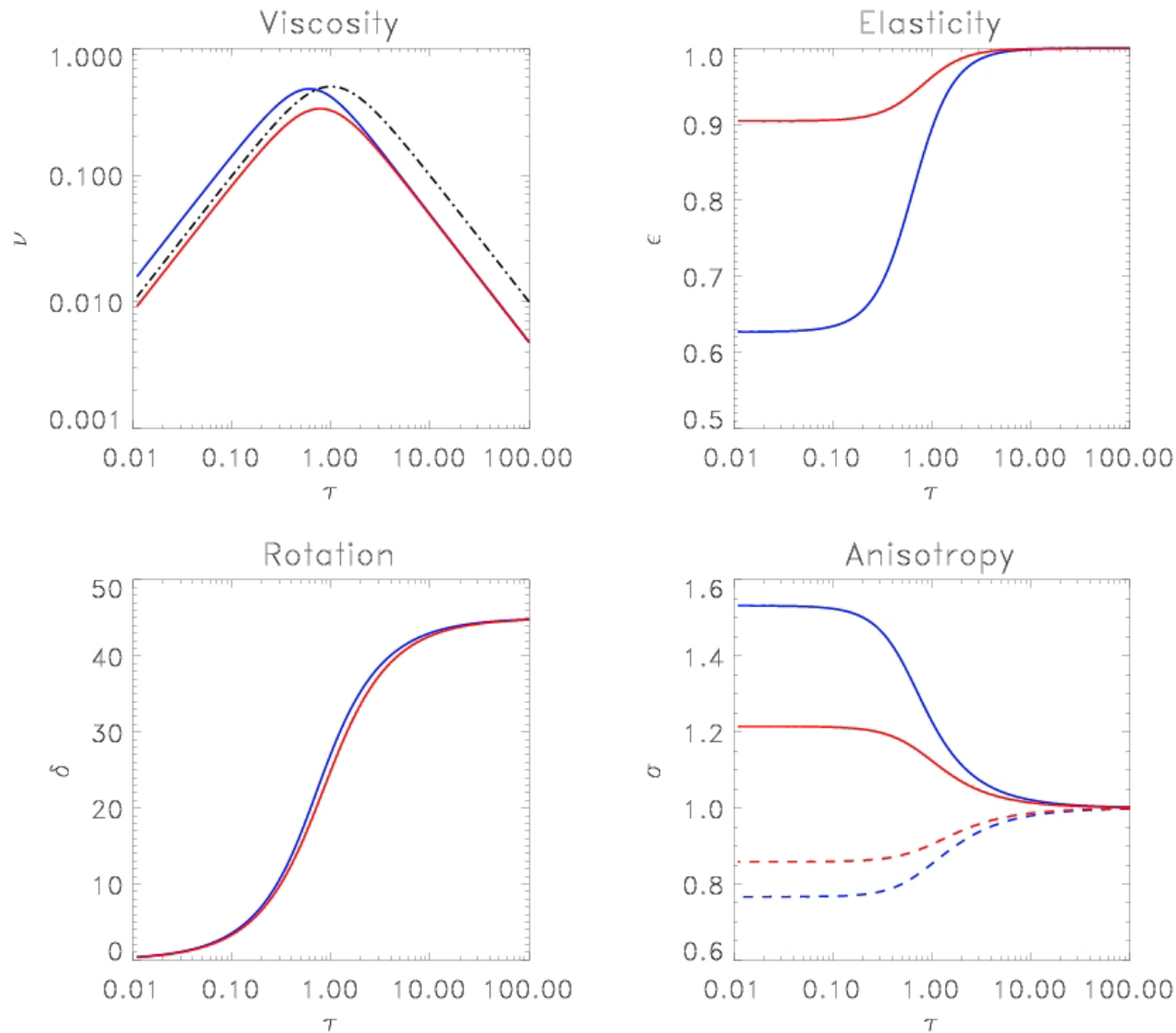
Goldreich & Tremaine (1978):  
Only momentum along  $\lambda$  is  
lost in inelastic collisions  
(assumption)

→ Elasticity  $\epsilon$

Coefficient of restitution

Only  $\frac{1}{2} (1 - \epsilon^2)$  of the kinetic  
energy is dissipated.

# Results for a thin disk



# Modifications for a thick torus

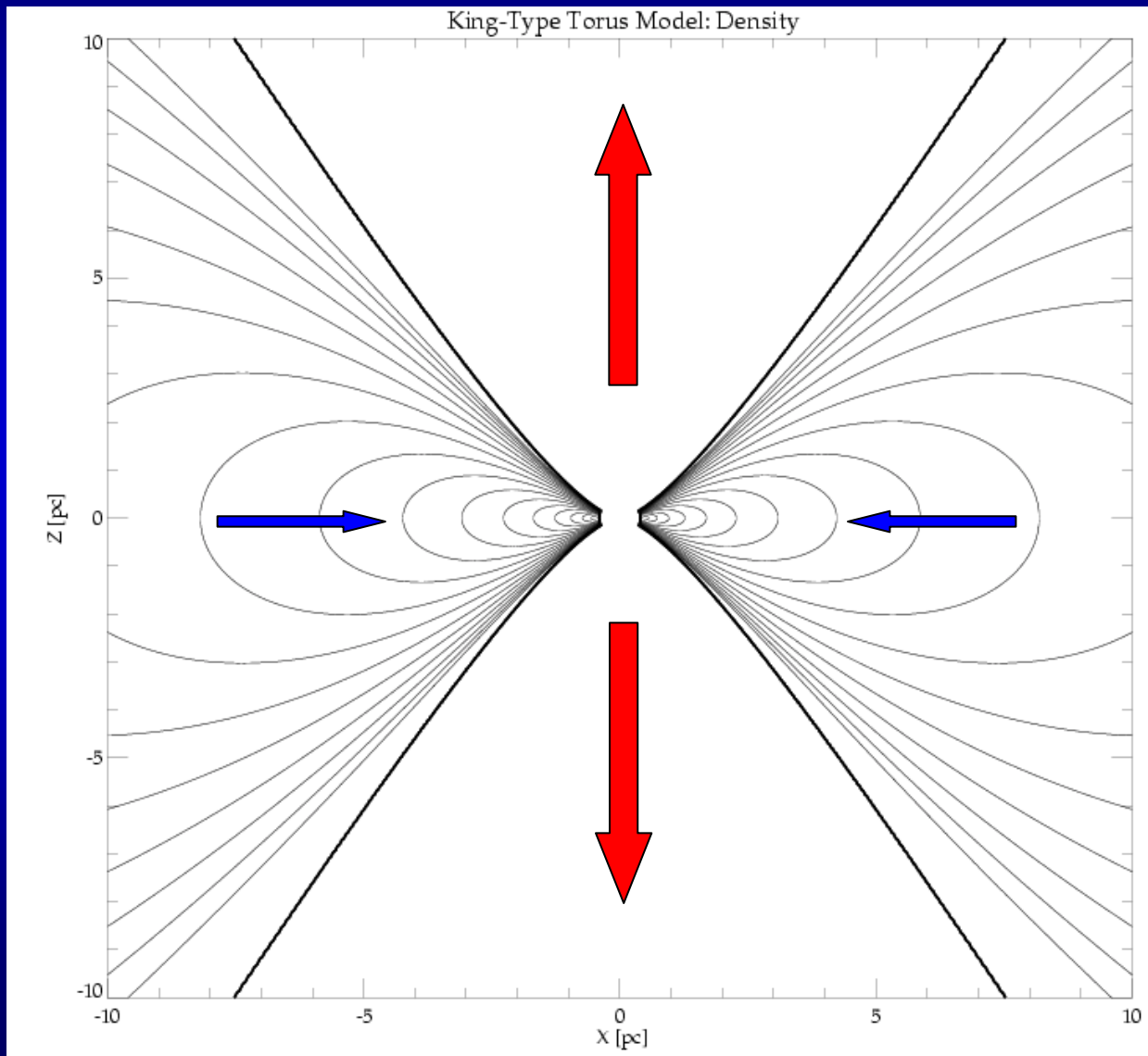
1. Advective terms  
(change of  $\sigma^2$  and compression )
2. No collective effects (enhanced viscosity via non-local interactions; Wisdom & Tremaine 1988)

$$\sigma \gg R_{\text{cloud}} \left| \frac{\partial \Omega}{\partial \log R} \right| \Rightarrow \Phi_{\text{v}} \ll 0.35 \tau$$

3. Elasticities  $\varepsilon$  as low as 0.3 are possible  
(45 % of kinetic energy can be dissipated in collisions)

**$\varepsilon$  determines  $\tau$**

# The vertical structure



**Exact solution  
for the density**

**in an arbitrary  
external potential**

**with a vertical  
cut-off height**

**Example for NGC 1068**

# Cloud Properties

- Tidal forces limit size of the largest clouds (→ shear-limit)
- Largest Clouds dominate appearance
- Quasi-stable clouds hold together by self-gravity

$$R_{\text{Cloud}}(c_s, M_{\text{Cloud}}) \rightarrow l_{\text{Coll}} = \frac{1}{n \pi R_{\text{Cloud}}^2}$$

- Typical cloud mass  $M \approx 50M_{\text{solar}}$

- Obscuration implies  $\tau \approx 1$

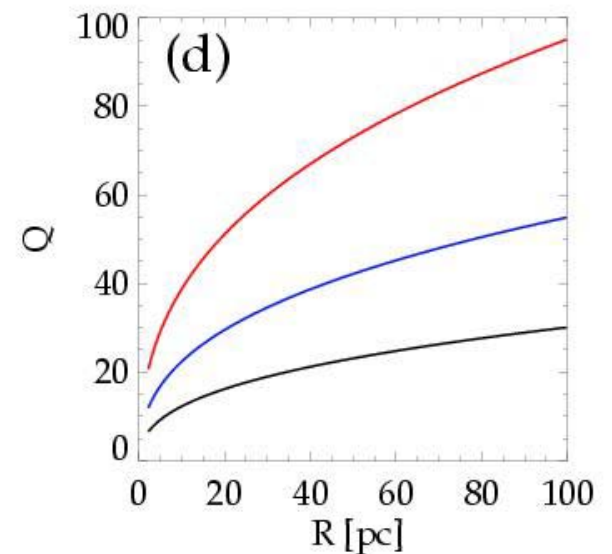
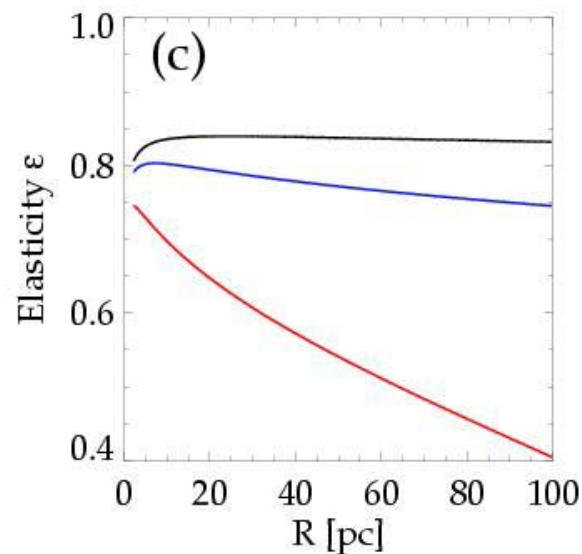
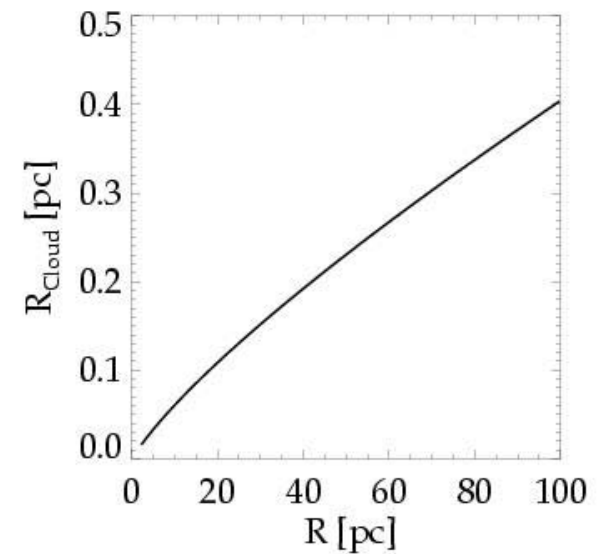
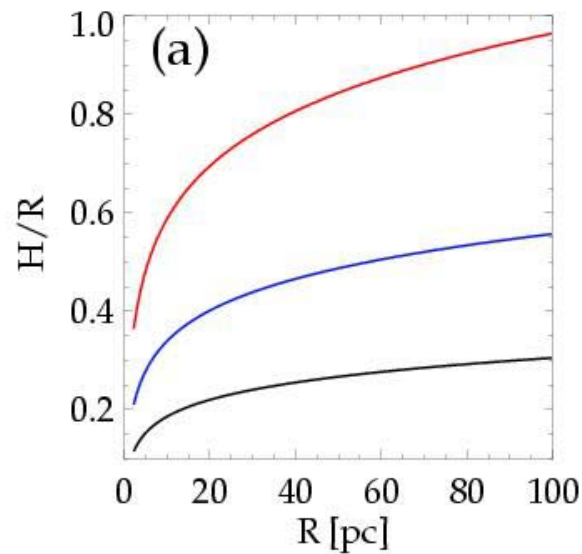
# Radial Structure

Red:  $m = 30$

Blue:  $m = 10$

Black:  $m = 3$

$m$  : mass accretion rate  
in units of the Eddington  
rate of the black hole



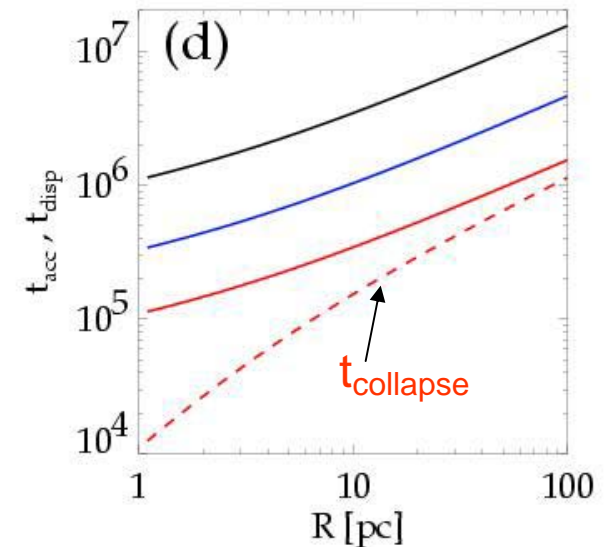
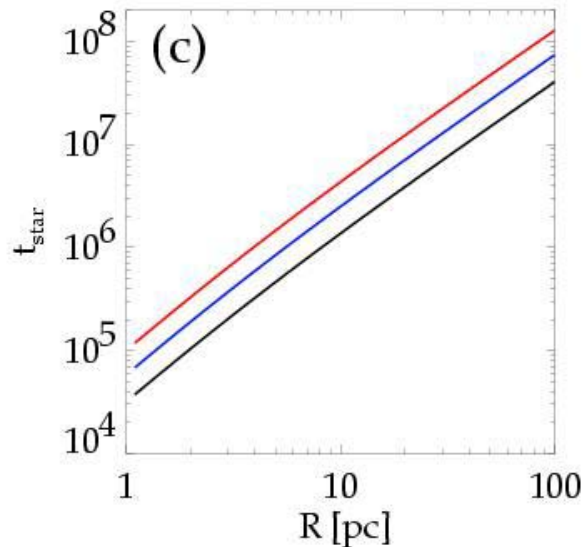
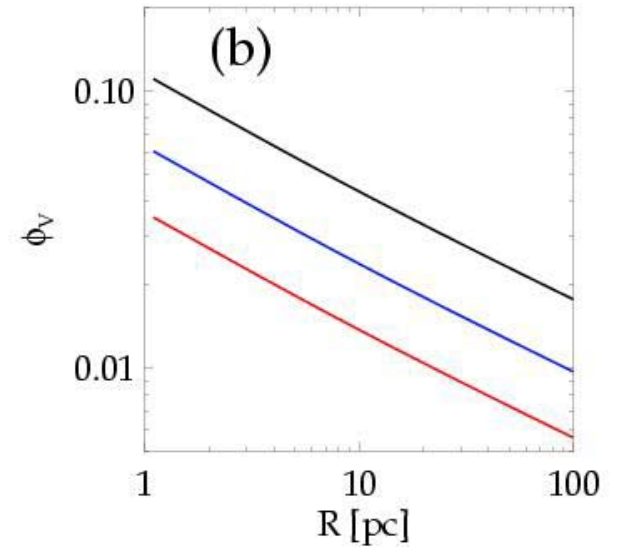
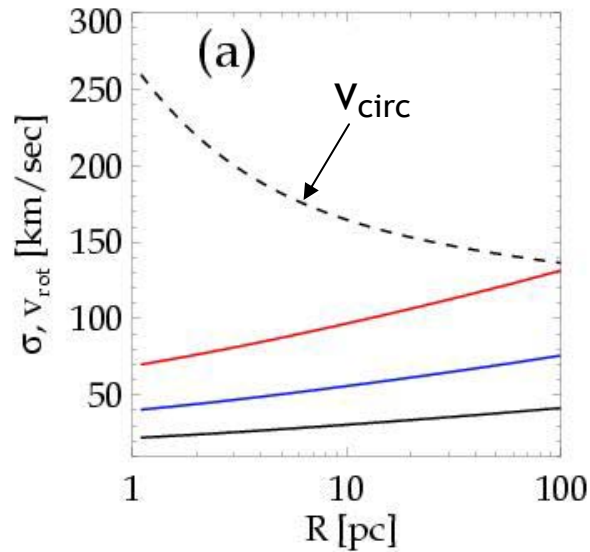
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# Timescales

1. Geometrically thick accretion flows (tori) rotate slightly sub-Keplerian ( $4 \cdot 10^5$  yr at 10 pc )  
Vertical hydrostatic equilibrium can be achieved on an orbital timescale

2. The accretion timescale (viscous timescale)

<  $10^6$  yr at 10 pc

$$t_{\text{acc}} \propto \tau^{-1} \dot{M}^{-1} R$$

3. Collapse time is the shortest timescale involved

$$t_{\text{Collapse}} = \frac{0.2}{\tau(1-\varepsilon)} t_{\text{Orbit}}$$

Without energy gains from accretion or other processes a  $\tau \sim 1$  torus would collapse to a thin disk within an orbital timescale

Energy gains from accretion is required or

4. The Circularization radius (angular momentum barrier) must lie within the inner boundary (sublimation radius)

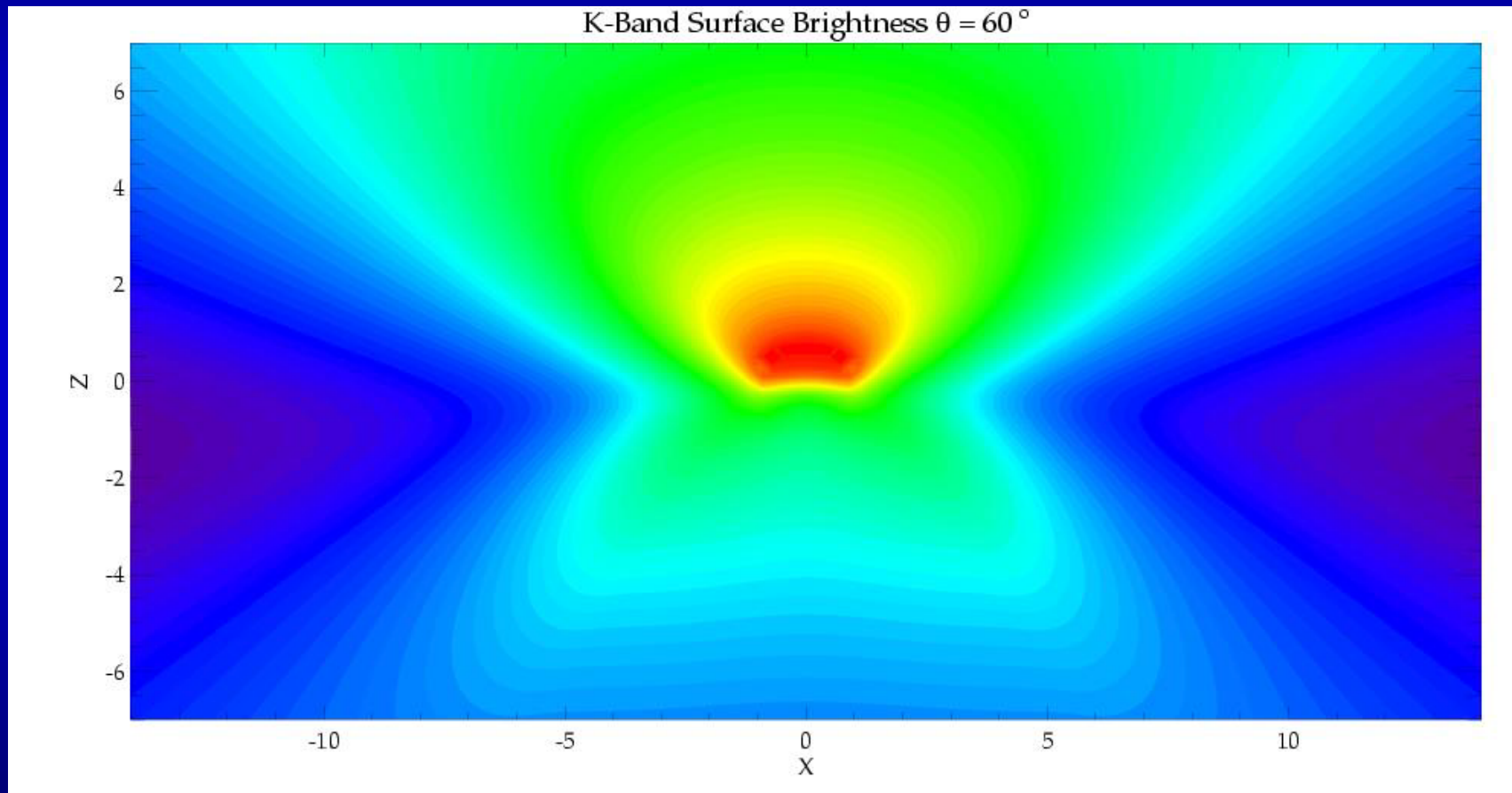
# Results

- Required accretion rate:  $\sim 10 M_{\text{solar}}/\text{yr}$
- Gas Mass 4 Mio.  $M_{\text{solar}}$  (within 80 pc)  $\rightarrow 10^4$  Clouds
- Cloud masses: 25 - 100  $M_{\text{solar}}$   
 $\rightarrow$  Clouds are optically thick individually
- Accretion lifetime  $\sim 1$  Mio. years (50 pc)
- Mean free path for Rad. Transfer

$$l_{\text{coll}} \propto R^q; \quad 0.5 \leq q \leq 1.5$$

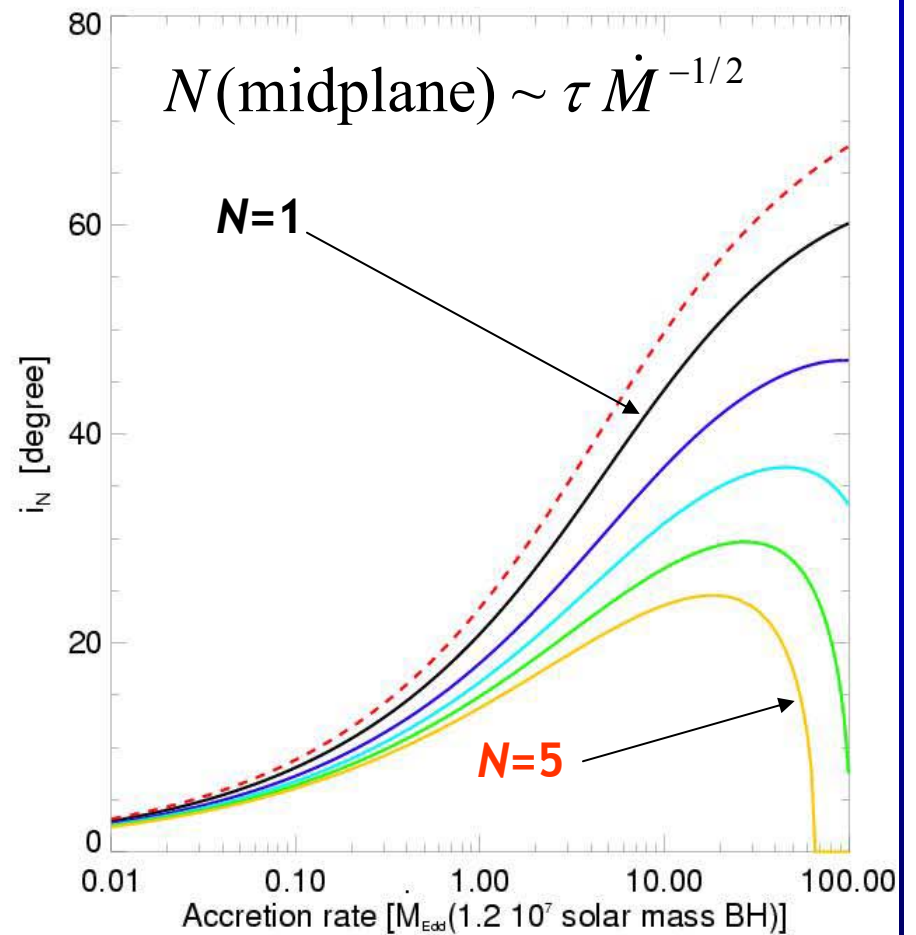
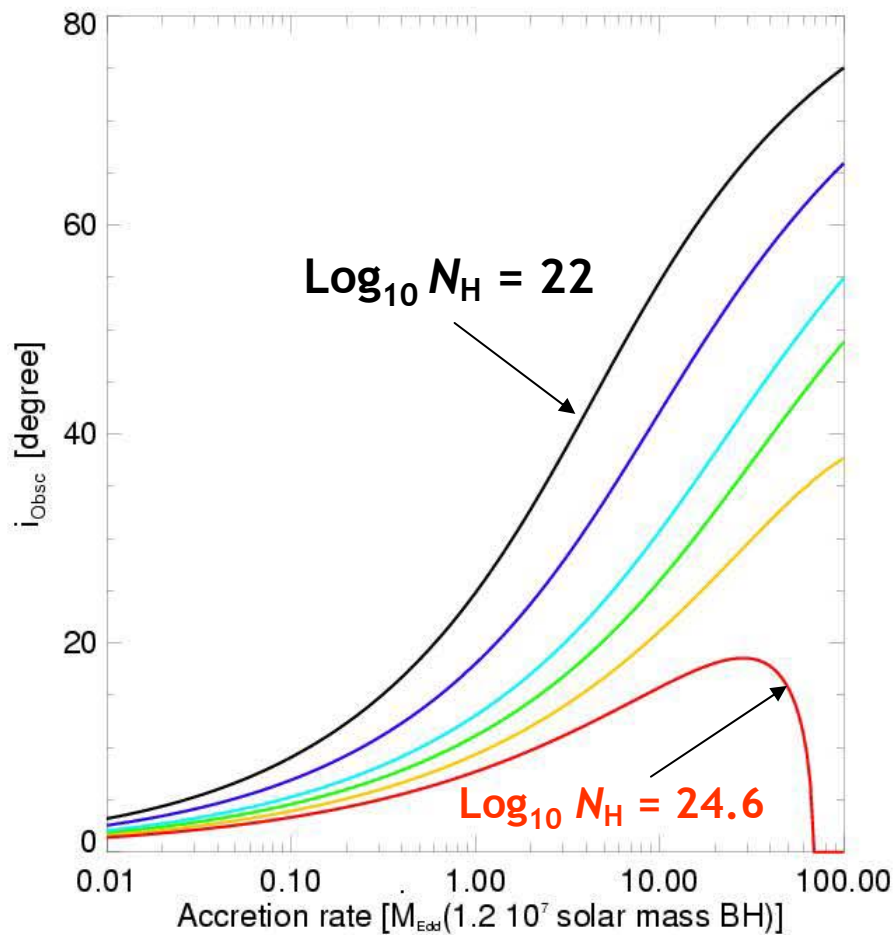
# Appearance in NIR

Based on the formalism of Nenkova, Ivezić & Elitzur (2002)



Dynamic range:  $10^5$  ; spatial scale in units of the sublimation radius

# Torus & Unified Scheme



# Predictions for Evolution

- Surface density

$$\Sigma \sim \tau \frac{M(R)}{2R^2} \frac{c_s}{v_{\text{circ}}}$$

- Thickness of Torus

$$H/R \sim \left( \frac{1 + \tau^2}{\tau^2} \frac{R \dot{M}}{M(R) c_s} \right)^{1/2}$$

- Mean free path

$$\frac{l_{\text{coll}}}{R} = \frac{H}{\tau R}$$

- Mean number of clouds (midplane)

$$N \propto \tau \dot{M}^{-1/2}$$

# Summary

- Dusty tori can be modelled as thick, clumpy accretion flows

- Geometrically thick tori need huge mass accretion rates

$$\dot{M}_{\text{Torus}} > \dot{M}_{\text{Edd,BH}}$$

- Provides a basis for radiative transfer models  
Anisotropic absorption and almost isotropic emission  $\longrightarrow$  spectral fitting

- Next question:

How do cloud collisions really look like ?

How does circumnuclear starformation feed the torus ?

